

Simulations including gas

Summary and perspectives

# Dynamics of Barred Galaxies in Triaxial Dark Matter Haloes

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# Outline of the presentation



- 2 Collisionless simulations
- Simulations including gas
- 4 Summary and perspectives

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# Introduction

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## Dark matter haloes: shapes

#### Cosmological simulations



- from the cosmological simulations of structure formation:
- dark matter haloes are generally triaxial
- major-to-minor axis ratio of as much as 2 is not uncommon

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## Dark matter haloes: shapes





Springel et al. (2008)

- from the cosmological simulations of structure formation:
- dark matter haloes are generally triaxial
- major-to-minor axis ratio of as much as 2 is not uncommon

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# Haloes and barred galaxies





 What about bar formation within an elongated potential?

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# Numerical simulations



#### N-body problem

 calculate force on each particle due to N - 1 other particles

• 
$$F_i = \sum_{i \neq j} \frac{G m_i m_j}{|r_i - r_j|^2}$$

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# Numerical simulations



#### N-body problem

 calculate force on each particle due to N - 1 other particles

• 
$$F_i = \sum_{i \neq j} \frac{G m_i m_j}{|r_i - r_j|^2}$$

 instead, approximate far-away particles by the centre-of-mass of that cell

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# Numerical simulations



#### colisionless simulations

- gyrfalcON code (Dehnen 2000)
- $N \sim 10^6$  particles
- mass resolution  $\sim 10^5 M_{\odot}$

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# Numerical simulations



#### GADGET2 (Springel 2005)

- TreePM: Tree and particle-mesh
- SPH: smoothed particle hydrodynamics

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# Part II

# **Collisionless simulations**

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# Initial conditions: haloes



halo IC: iterative method (Rodionov et al. 2010)

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# Initial conditions: density profiles

#### DISC (exponential)

$$\rho_d(R,z) = \frac{M_d}{4\pi z_0 R_d^2} \exp\left(-\frac{R}{R_d}\right) \operatorname{sech}^2\left(\frac{z}{z_0}\right)$$

#### HALO (Hernquist)

$$ho_h(r) = rac{M_h}{2\pi^{3/2}} rac{lpha}{r_c} rac{\exp{(-r^2/r_c^2)}}{r^2+\gamma^2}$$

 $R_d = 1$  is the unit of length (say, 3.5 kpc)

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# Epycicle approximation

#### Measure from the triaxial haloes:

- $\epsilon_{pot}(R)$  : ellipticity of halo potential (as a function of radius)
- $v_c(R)$  : circular velocity (as a function of radius)

# ellipticity of the orbits and ellipticity of the potential

$$\epsilon_{R} = \left[\frac{\frac{2v_{c}^{2}}{R} + \frac{dv_{c}^{2}}{dR}}{\frac{2v_{c}^{2}}{R} - \frac{dv_{c}^{2}}{dR}}\right]_{R_{0}} \epsilon_{pot}$$

# reassignment of disc orbit coordinates:

$$\begin{aligned} \mathbf{\hat{r}} & R = R_0 \left[ 1 - \frac{\epsilon_R}{2} \cos(2\Omega_0 t) \right] \\ \varphi &= \Omega_0 t + \frac{\epsilon_R + \epsilon_v}{4} \sin(2\Omega_0 t) \end{aligned}$$

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# Disc growth





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# Circularisation of the haloes



----- t=400 t=800

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# Haloes: initially triaxial



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# Discs: initially elliptical



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# Circularisation of the haloes



halo 1 halo 2 halo 3

- two phases of circularisation
- circularisation linked to bar strength

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## Elliptical versus circular discs



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### Elliptical versus circular discs



initially circular disc x initially elliptical disc

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# Bar formation causes further halo circularisation



#### compare halo shapes in the absence of bars:

- Iarge halo core
- Iess massive disc
- In the second second
- rigid disc
- artificially axisymmetric disc

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# Bar formation causes further halo circularisation





- Iarge halo core
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# Bar formation causes further halo circularisation





- large halo core
- 2 less massive disc
- In the second second
- I rigid disc
- artificially axisymmetric disc

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# Bar formation causes further halo circularisation

#### how to check this?

#### compare halo shapes in the **absence of bars**:

- large halo core
  - Iess massive disc
- In the test of test
- Irigid disc
- artificially axisymmetric disc

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## 1. Large halo core



halo too susceptible to circularisation

• halo completely circularised by disc growth alone

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### 2. Less masive discs



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# 3. Hot discs



- very weak bars
- haloes remain triaxial

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# 4. Rigid discs



no bars at all (analytic disc potential)haloes remain triaxial

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## 5. Artificially axisymmetric discs





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# Vertical flattening



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## Peanut-shaped bulges



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# Halo kinematics: disc-like halo particles



- some halo particles rotate (in a layer within |z| < 2 kpc)</li>
- rotation is less important in triaxial models





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# Halo kinematics: velocity anisotropy



 haloes remain anisotropic (even after circularisation)





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- a fraction of the disc mass is in the form of gas
- 15 models: 3 halo shapes × 5 gas fractions



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# Gas fraction



10

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# Disk, stars, gas: face-on (t=6 Gyr)



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## Disk, stars, gas: face-on (t=10 Gyr)



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## Disk, stars, gas: edge-on





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# Halo circularisation



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# Different stellar ages: bar strength

#### disc components:

 stars, young stars, youngest stars and gas



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## Different stellar ages: bar strength



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## Different stellar ages: angular momentum



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#### Disc vertical structure



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# Halo kinematics: disc-like halo particles



 peak tangential velocities correlated to bar strength



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# Part IV

# Summary and perspectives

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# Summary

#### circular discs are not adequate IC for triaxial haloes

- a halo is circularised in two phases:
  - o during disc growth
  - during bar formation

in the absence of a bar the halo may remain triaxial

- presence of gas inhibts strong bars more importantly than halo triaxiality
- rotation of disc-like halo particles is more important in the spherical case and is correlated do bar strength
- Itriaxial haloes retaint the anisotropy of their velocity dispersions even after being circularised

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Summary and perspectives

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Summary and perspectives

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Summary and perspectives

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Summary and perspectives

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Summary and perspectives

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## Perspectives

- Use models for statistical study of the orbits.
- I How do gas properties (SF, feedback, etc) affect the evolution?
- So far we have only considered isolated galaxies. It would be interesting to study interactions with such models.
- Similar work, but in a cosmologically-motivated setting.

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the end