Lecture 8

stellar atmospheres prof. Marcos Diaz



treasure map: Williams, R. ("The Analysis of Emission lines", 2009) Gray: pg 232 H&M: pg 144, 181, 186 Bohn-Vitense: pg 117,167 Rutten: pg 83

The line absorption and emission

I. the classical radiation driven damped oscillator:

 $m(\ddot{\mathbf{d}} + \omega_0^2) = e\mathbf{E}_0(\omega) - m\gamma \,\dot{\mathbf{d}}$

with the classical damping constant:

 $\gamma \equiv \frac{2e^2\omega_0^2}{3mc^3}$ solved for charged particle acceleration provide:

$$\ddot{\mathbf{d}} = \dot{\mathbf{v}} = \frac{e\,\omega^2}{m} \Re e \left[\frac{\mathbf{E}_0(\omega)e^{i\omega t}}{(\omega^2 - \omega_0^2) + i\,\gamma\,\omega} \right]$$

with Poynting flux:

$$P(t) = \frac{2e^2\ddot{d}}{3c^3}$$

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EM radiation orthogonal E and H fields with amplitude E_0 and frequency $\omega = 2\pi v$ yields a randomly directed Poynting vector power:

$$\left\langle P(\omega) \right\rangle_{\mathrm{T}} = \left(\frac{e^4 \omega^4}{3m^2 c^3} \right) \frac{E_0^2}{(\omega^2 - \omega_0^2)^2 + \gamma^2 \omega^2}.$$

with
$$I_0 \equiv E_0^2 / 8\pi$$
 and $\langle P(\omega) \rangle_{T,4\pi} = \oint \sigma(\omega) I_0 d\Omega$ (6)

(Poynting flux)

$$\sigma(\omega) = \frac{(8\pi e^4 \omega^4 / 3m^2 c^4)}{[(\omega^2 - \omega_0^2)^2 + \gamma^2 \omega^2]}$$

for a free electron $\omega_0 = 0$:

$$\sigma(\omega) \to \sigma_{\rm T} \equiv \frac{8\pi e^4}{3m_e^2 c^4} = 6.65 \times 10^{-25} \ {\rm cm}^2$$
$$\sigma_{class} = \int_0^\infty \sigma(\omega) d\omega = \frac{\pi e^2}{mc} \to \frac{\pi e^2}{mc} f \quad (\sigma_{\rm quantum})$$

f = fraction of quantum oscillators allowed to participate

II. Equilibrium LTE gf $N_{u}A_{ul} + N_{u}B_{ul}I_{v} + C_{ul} = N_{l}B_{lu}I_{v} + C_{lu}$ (rate equation)

With detailed balance of collisional transitions:

$$I_{v} = \frac{A_{ul}}{B_{lu}(N_{l}/N_{u}) - B_{ul}} \qquad \text{LTE} \rightarrow \qquad I_{v} = \frac{A_{ul}}{(g_{l}/g_{u})B_{lu}e^{hv/kT} - B_{ul}}$$

Wien (hv >> kT)

$$I_{v} \propto v^{3} e^{-\frac{hv}{kT}} \rightarrow B_{ul} = B_{lu} g_{l} / g_{u}$$

Rayleigh-Jeans ($hv << kT$) $\rightarrow I_{v} \approx \frac{A_{ul}}{(g_{l}/g_{u})B_{lu} - B_{ul} + (g_{l}/g_{u})B_{lu}hv/kT}$

$$I_{v} \approx \frac{2kT v^{2}}{c^{2}} \rightarrow A_{ul} = \frac{2hv^{3}}{c^{2}}B_{ul}$$

$$\sigma_{tot} = \frac{\pi e^2}{mc} f_{lu} = B_{lu} h \nu = \frac{mc^3}{2\pi e^2 \nu^2} \frac{g_u}{g_l} A_{ul}$$

$$g_l f_{lu} = g_u f_{ul} = g f$$

with $g_{l,u} = 2J_{l,u} + 1$

III. Bound state quantum numbers & selection rules

$$-\frac{\hbar^2}{2m}\nabla^2\psi + V(x, y, z)\psi = E\psi \quad \text{with} \quad \frac{1}{m} = \frac{1}{m_N} + \frac{1}{m_e} = \frac{m_N + m_e}{m_N m_e}.$$

$$-\frac{\hbar^2}{2m}\frac{1}{r^2} \left[\frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r}\right) + \frac{1}{\sin\theta} \frac{\partial}{\partial \theta} \left(\sin\theta \frac{\partial \psi}{\partial \theta}\right) + \frac{1}{\sin^2\theta} \frac{\partial^2 \psi}{\partial \phi^2}\right] - \frac{Ze^2}{r}\psi = E\psi.$$
with $\psi(r, \theta, \phi) = R(r)\Theta(\theta)\Phi(\phi)$

$$\frac{d^2\Phi}{d\phi^2} + m^2\Phi = 0,$$

$$\frac{1}{\sin\theta}\frac{d}{d\theta}\left(\sin\theta \frac{d\Theta}{d\theta}\right) + \left[l(l+1) - \frac{m^2}{\sin^2\theta}\right]\Theta = 0, \quad \text{Legendre } (l,m)$$

$$\frac{1}{r^2}\frac{d}{dr}\left(r^2\frac{dR}{dr}\right) - \frac{l(l+1)}{r^2}R + \frac{2m}{\hbar^2}\left(E + \frac{Ze^2}{r}\right)R = 0.$$
Laguerre (n,l)

Solution for single electron systems:

$$\psi_{n,l,m}(r,\theta,\phi) = A_{n,l,m}L_{n,l}(r) P_{l,m}(\cos\theta) e^{im\phi}$$

For many electron systems the energy levels are given by the quantum numbers of each electron:

- n_i : orbit of electron *i*
- l_i : orbital angular momentum of electron *i*; with $l_i < n_i (l_i = 0 \text{ or } "s"; l_i = 1 \text{ or } "p"; ...)$
- *L*: orbital angular momentum the state: vector sum of l_i (L = 0 or "S"; L = 1 or "P"; ...)
- S: spin angular momentum of the state: vector sum of s_i ($s_i = 1/2$)
- J: total angular momentum of the state: vector sum of L + S

Spectroscopic Notation:

e.g. for a 12 electron atom:
$$(n_i, l_i)^{2S+1}L_J$$

$$[1s^2 2s^2 2p^6] 3s^{2-1}S_0$$

for 2 different Helium states:

$$1s^{2} {}^{1}S_{0}; 1s2p {}^{3}P_{2}$$

Solution for single electron systems:

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For many electron systems the energy levels are given by the quantum numbers of each electron:

- orbit of electron *i* n_i :
- l_i : L: orbital angular momentum of electron *i*; with $l_i \le n_i (l_i = 0 \text{ or } "s"; l_i = 1 \text{ or } "p"; ...)$
- orbital angular momentum of the state: vector sum of l_i (L=0 or "S"; L=1 or "P"; ...)
- spin angular momentum of the state: vector sum of s_i ($s_i = 1/2$) *S*:
- total angular momentum of the state: vector sum of L + SJ:

Spectroscopic Notation:

e.g. for a 12 electron atom:
$$(n_i, l_i) \stackrel{2S+1}{\longrightarrow} L_J$$

$$[1s^2 2s^2 2p^6] \ 3s^{2-1}S_0$$

for 2 different Helium states:

$$1s^{2} {}^{1}S_{0}; 1s^{2}p {}^{3}P_{2}$$



from R. Williams 2017



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Selection rules

$$\Delta S = (S - S) = 0*$$

$$\Delta L = (L_l - L_l) = 0, \pm 1, \pm 2(E2)$$

$$\Delta J = (J_l - J_l) = 0, \pm 1, \pm 2(E2)$$

i. for atom **electric dipole** E1(permitted) transitions:

→ change in parity of the wave function that describes the state → $\Delta l_i = \mp l$; $\Delta n = any$

ii. for atom electric quadrupole E2 / mag dipole M1(forbidden "[X]") transitions:
 → no change in parity of the wave function that describes the state

→ $\Delta l_i = 0$; $\Delta n = 0$ (same eletronic config) OR one electron change n (M1)

* LS coupling violation $\Delta S \neq 0 \rightarrow$ intercombination (or semi-forbidden "[X") transitions.

Atomic Energy Levels: O III (1s²2s²2p²)



Main link for Atomic data:

https://www.nist.gov/pml/atomic-spectra-database

Grotrian Diagrams:

https://ned.ipac.caltech.edu/level5/Ewald/Grotrian/frames.html