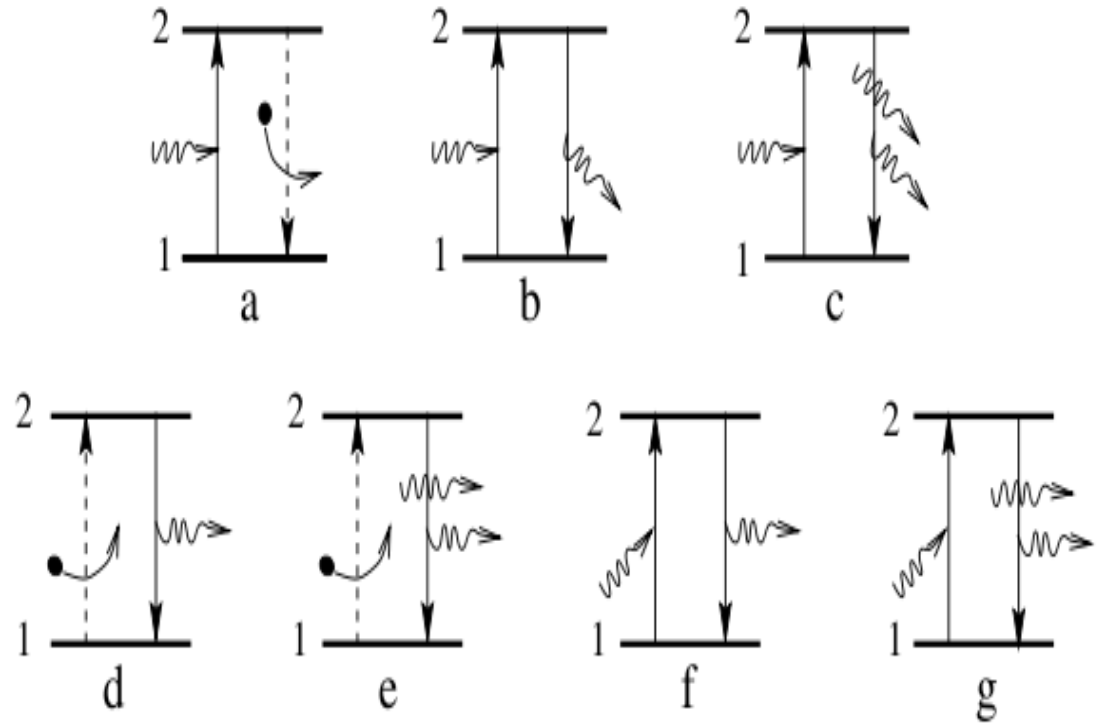


# Atmosferas Estelares

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IAG-USP 2023



Contributions to beam intensity from two-level bound-bound sequences.

from Rutten 2004

## Lecture 2

### **treasure map:**

H&M: pg 334-361

Rutten: pg 4

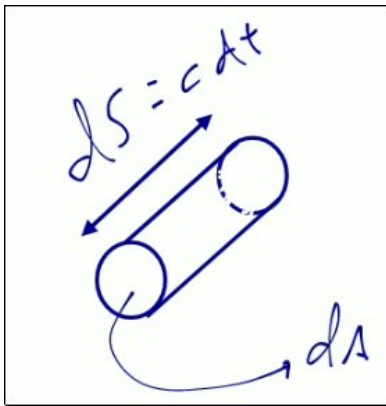
Vitense: pg 39

Gray: pg 127

Keep it small. keep it *simple*:

local linear approximation to beam contributions:

$$k(\mathbf{r}, \mathbf{n}, \nu, t) I(\mathbf{r}, \mathbf{n}, \nu, t) + \eta(\mathbf{r}, \mathbf{n}, \nu, t)$$



$$dN_{total\ ph} = \frac{1}{h\nu} I(\mathbf{r}, \mathbf{n}, \nu, t) dA d\omega d\nu dt$$

i. “ $k_\nu$ ” the linear absorption coefficient

$$P = \frac{\Delta a}{\Delta A}; \quad P = \frac{da}{dA} \quad (1) \quad \text{interaction probability per photon}$$

$$da = \sigma N_{particles} dA dS$$

is the total unsaturated cross-section, in (1):

per unit length  $S$ :  $P = \sigma N_{part} = \chi$  with mean free path:  $1/\chi$

$$dN_{ph} = dN_{tot\ ph} P$$

$$dN_{ph} = \frac{\chi}{h\nu} I(\mathbf{r}, \mathbf{n}, \nu, t) dA d\omega d\nu dt$$

$$dE = \chi I(\mathbf{r}, \mathbf{n}, \nu, t) dA d\omega d\nu dt$$

with  $\chi = \text{sct.} + \text{true abs.} - \text{induced emission}$

ii.  $\eta_\nu$  (the local emission coefficient)

$$dE_{ph} = \eta(\mathbf{r}, \mathbf{n}, \nu, t) dA d\omega d\nu dt$$

with  $\eta =$  all spontaneous emission

With (i) and (ii), along  $dS$  the radiative energy change by:

$$\begin{aligned} dE_{ph} &= [I(\mathbf{r} + d\mathbf{S}, \mathbf{n}, \nu, t + dt) - I(\mathbf{r}, \mathbf{n}, \nu, t)] dA d\omega d\nu dt \\ &= \eta(\mathbf{r}, \mathbf{n}, \nu, t) - \chi(\mathbf{r}, \mathbf{n}, \nu, t) I(\mathbf{r}, \mathbf{n}, \nu, t) \end{aligned}$$

$$diff|_{r,t} = \frac{\partial I}{\partial t} dt + \frac{\partial I}{\partial S} dS = \left( \frac{\partial I}{\partial S} + \frac{1}{c} \frac{\partial I}{\partial t} \right) dS$$

(with  $dS = c dt$ )

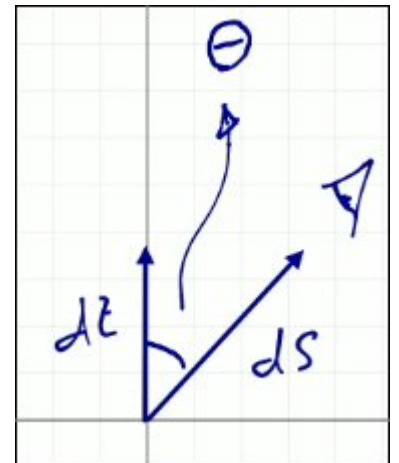
with  $\left(\frac{\partial I}{\partial S}\right) dS = \mathbf{n} \cdot \nabla I :$

$$\left(\frac{1}{c} \frac{\partial}{\partial t} + \mathbf{n} \cdot \nabla\right) I(\mathbf{r}, \mathbf{n}, \nu, t) = \eta(\mathbf{r}, \mathbf{n}, \nu, t) - \chi(\mathbf{r}, \mathbf{n}, \nu, t) I(\mathbf{r}, \mathbf{n}, \nu, t)$$

time-dependent transfer equation

steady state:  $\frac{\partial I}{\partial t} = 0$

plane-parallel:  $\frac{dz}{dS} = \cos(\theta) = \mu$



→  $\mu \frac{dI(\nu, \mu, z)}{dz} = \eta(\nu, \mu, z) - \chi(\nu, \mu, z) I(\nu, \mu, z)$

*optical depth*      $d\tau_\nu \equiv -x_\nu dS = -x_\nu \frac{dz}{\mu}$      “lagrangian coordinate”  
for radiation

$$x_\nu \frac{dI(\nu, \mu, z)}{d\tau_\nu} = -\eta(\nu, \mu, z) + x(\nu, \mu, z) I(\nu, \mu, z)$$

$$\frac{dI_\nu}{d\tau_\nu} = I_\nu - S_\nu \quad \text{with} \quad S_\nu \equiv \frac{\eta_\nu}{x_\nu}$$

with  $\tau_{\nu\perp}$  in the normal direction:      $\mu \frac{dI_\nu}{d\tau_{\nu\perp}} = I_\nu - S_\nu$

*optical depth*  $d\tau_\nu \equiv -x_\nu dS = -x_\nu \frac{dz}{\mu}$  “lagrangian coordinate”  
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with  $\tau_{\nu\perp}$  in the normal direction:  $\mu \frac{dI_\nu}{d\tau_{\nu\perp}} = I_\nu - S_\nu$

$$dN_{ph} = \frac{\eta_\nu}{h\nu} d\nu d\omega dt dA ds \quad (\text{number of photons emitted at } dV = dA ds)$$

$$dN_{ph} = \frac{\eta_\nu}{x_\nu} \frac{1}{h\nu} d\nu d\omega dt dA x_\nu ds = S_\nu \frac{4\pi}{h\nu} d\nu dt dA d\tau_\nu \quad (S_\nu \text{ isotropic})$$

$S_\nu$  is the energy added within  $\Delta\tau_\nu = 1$ ;  $J_\nu$  is the energy removed within  $\Delta\tau_\nu = 1$



# The Formal Solution

$$\frac{dI_v}{d\tau_v} = I_v - S_v$$

*the standard integrating factor solution of a first-order differential equation:*

$$\frac{dI_v}{d\tau_v} + p(\tau)I_v = q(\tau) \quad @ \text{ an arbitrary depth } \tau :$$

$$I(\tau) = e^{-\int_0^\tau p(\tau')d\tau'} \left[ \int_0^\tau e^{\int_0^{\tau'} p(\tau'')d\tau''} q(\tau')d\tau' + c \right]$$

## The Formal Solution (cont.)

with:  $p(\tau) = -1; \quad q(\tau) = -S(\tau)$

$$I(\tau) = -e^{\tau} \left[ \int_0^{\tau} e^{-\tau'} S(\tau') d\tau' + c \right]$$

then relate the intensities at 2 arbitrary depths  $\tau_1$  and  $\tau_2$  as follows:

$$I(\tau_2) e^{-\tau_2} - I(\tau_1) e^{-\tau_1}$$

assuming p.p. geometry, using  $\tau_{1,2} (\mu=1) \equiv \tau_{1\perp,2\perp}$ :

$$\tau_1 = \tau_{1\perp} / \mu \quad \text{and} \quad \tau_2 = \tau_{2\perp} / \mu$$

## The Formal Solution (cont.)

$$\tau_{\perp} = \bar{\tau} \quad (\text{change of notation to avoid double indices})$$

$$I(\bar{\tau}_1, \mu) = I(\bar{\tau}_2) e^{\frac{\bar{\tau}_1 - \bar{\tau}_2}{\mu}} + \frac{1}{\mu} \int_{\bar{\tau}_1}^{\bar{\tau}_2} S(\tau') e^{\frac{\bar{\tau}_1 - \tau'}{\mu}} d\tau'$$

The formal solution above is particularly useful when the intensity is known at a given depth and a recipe for the depth dependence of the source function can be found.