Day 3. Cosmic chemical evolution

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History of the Universe





Hubble Space Telescope • Wide Field Planetary Camera 2

PRC97-34a • ST Scl OPO • October 21, 1997 • B. Whitmore (ST Scl) and NASA







M 82 (NGC 3034)

FOCAS (B, V, Hlpha)

Subaru Telescope, National Astronomical Observatory of Japan

March 24, 2000

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3 types of galaxy models





One-zone (monolithic) models

Instantaneous mixing approximation
SFR/inflow/outflow with analytic formula
Average evolution of a galaxy (or a shell of galaxy)

Semi-analytic (hierarchical) models

Mass assembly history based on λCDM scenario
Global properties of galaxies in a large scale



Chemo-dynamical simulations

Inhomogeneous chemical enrichment
Internal structures kinematics of stars/gas, spatial distribution of elements

Other types of models

- Stochastic models (Argast+02; Ishimaru & Prantzos; Cescutti+; Wehmeyer+)
- Chemodynamical models without hydro (e.g., Minchev & Chiappini)

Basic equations of Chemodynamical Simulations

CK 2004 CK, Springel, White 2007 Taylor & CK 2014

N-body methods

Gravity for DM and stars ★ direct summation – O(N²) GRAPE, GPU

$$m_{i}\frac{d^{2}x_{i}}{dt^{2}} = \sum_{j \neq i}Gm_{i}m_{j}\frac{x_{j} - x_{i}}{|x_{j} - x_{i}|^{3}}$$

***** Tree method – O(N logN)<sup>-
$$k$$</sup> distant particles are bundled up

* Particle-mesh method – O(N logN)

$$\nabla^2 \phi = 4\pi G \rho \rightarrow -k^2 \phi_k = 4\pi G \rho_k \rightarrow \phi(\mathbf{r})$$





SPH method

onling karnal 11/

SPH formulation (Navarro & White 1993)

Smoothed Particle Hydrodynamics for gas (Gingold &

Monaghan 1977; Lucy 1977; Monaghan 1992 for a review)

★ equations to integrate:

$$\frac{D\rho}{Dt} + \rho \nabla \cdot \mathbf{v} = 0$$

$$\frac{D}{Dt} + \rho \nabla \cdot \mathbf{v} = 0$$

$$\frac{D}{Dt} = -\frac{1}{\rho} \nabla P - \nabla \Phi$$

$$\frac{Du}{Dt} = \frac{P}{\rho^2} \frac{D\rho}{Dt} + \frac{\nabla \cdot (\kappa \nabla T)}{\rho} + \frac{\Gamma - \Lambda}{\rho}$$

$$\nabla^2 \Phi = 4\pi G\rho$$

$$P = (\gamma - 1)\rho u$$

$$P_i = \sum m_j W(\mathbf{r}_i - \mathbf{r}_j; h)$$

$$\frac{D\mathbf{v}_i}{Dt} = -\sum m_j \left(\frac{P_i}{\rho_i^2} + \frac{P_j}{\rho_j^2} + \Pi_{ij}\right) \nabla_i W(\mathbf{r}_i - \mathbf{r}_j; h) - (\nabla \Phi)_i$$

$$\frac{Du_i}{Dt} = -\sum m_j \left(\frac{P_i}{\rho_i^2} + \frac{1}{2} \Pi_{ij}\right) (\mathbf{v}_i - \mathbf{v}_j) \nabla_i W(\mathbf{r}_i - \mathbf{r}_j; h) + \frac{\Gamma - \Lambda}{\rho}$$

Radiative Cooling



- * Calculated by MAPPINGS III (Sutherland & Dopita 1993).
- * Molecule cooling (T<10⁴K) is not included.
- * UV background heating is subtracted (Kats, Weinberg & Hernquist 96).

Star Formation

★ Katz (1992)

(i) (∇ ⋅ v)_i < 0, $t_{cool} = \frac{\rho u}{\Lambda}$,
★ converging flow
★ rapid cooling
★ Jeans unstable
(ii) $t_{cool} < t_{dyn}$, $t_{dyn} = \frac{1}{\sqrt{4\pi G\rho}}$,
(iii) $t_{dyn} < t_{sound}$, $t_{sound} = \frac{h_i}{c_s}$,

* Density criteria (Stinson+06; Governato+10, Nature) $\rho > \rho_{crit}$

* Star formation timescale $t_{sf} = \frac{1}{c} t_{dyn}$, c=0.02-0.1

★ Stochastic treatment for the timestep ∆t
 ★ Form stars if random number [0,1] is smaller than

$$p \leq \frac{m}{m_{g,0}/N_*} (1 - \exp[-\frac{\Delta t}{t_{sf}}]) \quad m_{*,0} = m_{g,0}/N_* \text{ with } N_* = 2.$$

Chemical Enrichment 1/2



Chemical Enrichment 2/2



 $\phi(m) = m^{-x} \frac{1-x}{m_u^{1-x} - m_\ell^{1-x}} \qquad \phi_{\rm d}(m) = m^{-x} \times \frac{-x}{m_{{\rm d},u}^{-x} - m_{{\rm d},\ell}^{-x}}$

Supernova Feedback 1/2

* Energy by stellar winds for given mass

$$e_{e,SW} = \begin{cases} 0.2 \times 10^{51} \begin{pmatrix} Z \\ Z_{\odot} \end{pmatrix}^{0.8} & (m_{2,u} < m \le m_{u}) \\ (0.2 \times 10^{51} & (m_{2,\ell} < m \le m_{2,u}) \end{pmatrix} & (erg), \\ (17) \\ \text{ * Energy per core-collapse supernova} \\ e_{e,II} = \int_{m_{1}}^{m_{2,u}} p_{e,m} \frac{1}{m} \phi(m) dm \\ p_{e,m} = \begin{cases} 1 & (m_{2,\ell} < m \le 20M_{sun}) \\ (1 - \varepsilon_{HN}) + [10,10,20,30]\varepsilon_{HN} & (20,25,30,40M_{sun}) \end{cases} \times 10^{51} (erg) \\ \varepsilon_{HN} = [0.5, 0.5, 0.4, 0.01, 0.01] & \text{for } Z \neq [0, 0.001, 0.004, 0.02, 0.05] \\ \text{ * Energy per Type Ia supernova} \end{cases}$$

$$e_{\rm e,Ia} = 1.3 \times 10^{51} (m_{1d,\ell} < m \le m_{1d,u}) (\rm erg)$$
 (19)

Supernova Feedback

★ Total energy, mass, metal ejection from a star particle j at time t to a neighbor gas particle i

$$\begin{split} E_{e,j}(t,Z_{j}) &= m_{*}^{0} [e_{e,SW}(Z_{j}) R_{SW}(t) + e_{e,II}(t,Z_{j}) + e_{e,Ia} R_{Ia}(t,Z_{j})] = \sum_{i=1}^{N_{FB}} E_{e,i}(t,Z_{j}) W(\mathbf{r}_{i} - \mathbf{r}_{j};h_{j}) \\ E_{m,j}(t,Z_{j}) &= m_{*}^{0} [e_{m,SW}(t,Z_{j}) + e_{m,II}(t,Z_{j}) + e_{m,Ia}(t,Z_{j})] = \sum_{i=1}^{N_{FB}} E_{m,i}(t,Z_{j}) W(\mathbf{r}_{i} - \mathbf{r}_{j};h_{j}) \\ E_{Z,j}(t,Z_{j}) &= m_{*}^{0} [e_{Z,SW}(t,Z_{j}) + e_{Z,II}(t,Z_{j}) + e_{Z,Ia}(t,Z_{j})] = \sum_{i=1}^{N_{FB}} E_{Z,i}(t,Z_{j}) W(\mathbf{r}_{i} - \mathbf{r}_{j};h_{j}) \end{split}$$

Internal energy (and mass, metal) that gas particle i get from star particle j

$$\frac{du_i}{dt} = \sum_j \int_{t-dt_j}^t dt \frac{E_{e,j}(t, Z_j) W(\mathbf{r}_i - \mathbf{r}_j; h_j)}{\sum_{k=1}^{N_{\text{FB}}} W(\mathbf{r}_k - \mathbf{r}_j; h_j)}$$
100% thermal

Other feedback methods

* artificial cooling off during heating (Governato+ 10)

* kinetic feedback (Navarro & White 93, Springel & Hernquist 03)

$$v_{\rm wind} = \sqrt{\frac{2\varepsilon_{\rm kin}E_{\rm SN}}{m}}$$

* kinetic feedback (Navarro & White 93, Springel & Hernquist 03)

$$v_{\text{wind}} = 484 \text{ km/s for } p \le \frac{m}{m_{g,0} / N_*} (1 - \eta \exp[-\frac{\Delta t}{t_{\text{sf}}}]), \eta = 2$$

* momentum driven-winds (Dave & Oppenheimer 06) $v_{\text{wind}} = 3\sigma \sqrt{f_L - 1} + 2\sigma, \ \eta = \frac{\sigma_0}{\sigma} = 2, \ \sigma_0 = 300 \text{ km/s}, \ f_I = 2 \text{ or}$ $f_L = f_{L,\odot} \times 10^{-0.0029*(\log Z + 9)^{2.5} + 0.417694}.$

* stochastic treatment of feedback (Dalla Vecchia & Schaye 12) * $\Delta \varepsilon$ (~1.82 10¹⁴ erg/s) is given if random number [0,1] is smaller than $p = f_{\text{th}} \frac{\epsilon_{\text{SNII}}}{\Delta \epsilon} \frac{m_*}{\sum_{i=1}^{N_{\text{ngb}}} m_i} , f_{\text{th}}=1, N_{\text{ngb}}=48$

Density-temperature diagram



CK, Springel, & White 2007

Isolated Disk

* Fixed NFW halo $M_{tot} = 10^{10}$ /h M_{\odot} CK, Springel, & White 2007 * 10% gas, N=160000 (m_{gas}=6.3×10³ /h Mu), λ =0.1 Face-on

t = 0.00 Gyr





~50% gas, ~25% metals are ejected by HNe. Resolution independent.

Other Simulation Codes

Cosmological

- ★ Springel 00 (Gadget SPH), Di ★ Steinmetz & Muller 94 (SPH) Matteo (AGN), Scannapieco (la), Schaye
- ★ Frenk, Okamoto, ... (Gadget)
- ★ Dave, Oppenheimer (Gadget)
- ★ Mosconi, Tissera, et al. 01 (AP3MSPH, Ia)
- ★ Teyssier 02, Devriendt (RAMSES - AMR)
- ★ Wadsley 04 (GASOLINE –SPH), Governato, Brook
- ★ Martinez-Serrano 08 (SPH)
- ★ Springel 10 (AREPO)

Galactic

- ★ Raiteri, Villata & Navarro 96 (Ia)
- ★ Carraro, Lia & Chiosi 98 (Ia)
- ★ Hensler 90, Recchi
- ★ Burkert, Naab 99 (N-body)
- ★ Hernquist, Hopkins 05 (AGN)
- * Mori 97
- * Nakasato 03
- Gibson, Kawata 04 (Ia), Brook *
- Tornatore et al. 04 (la) *
- ★ Elmegreen
- * . . .

* ...

AGN feedback



Active Galactic Nuclei (AGN)

* 1950s: radio sources discovered, optical counterparts?

* 1960: 3C48, blue star with broad emission lines

★ 1963: 3C273, redshift $z=0.158 \rightarrow \text{distant object}$



★ very luminous, 10¹² L_{sun},~ 4×10⁴⁵ erg/s, energyproduction mechanism?

The Unified Model of AGN

Seyfert 1 with broad & narrow emission lines



Seyfert 2 with narrow emission lines

AGN Feedback?

AGN not only suppress SF but also enhance SF!



Figure 2. Evolution of density in simulation E. The width and height of each panel are 1 kpc. (A color version, animations (A, B, C, D, E), and the complete figure set (five images) of this figure are available in the online journal.)

Wagner & Bicknell 2011; and many

Co-evolution of SMBH & galaxies

Quasar space density Schmidt+ 95



Stellar + AGN winds



ESO optical/submm/Xray

 $M_{BH} \sim 10^{6} M_{\odot}$ SFR(ring)~0.05M_☉/yr

AGN feedback in Semi-Analytic Models

* Croton et al. 2006, Semi-analytic Model

- ★ quasar mode: BH growth by mergers
- ★ radio mode: heating by hot gas accretion on BH



Standard AGN model in hydro-simulations

* Co-evolution of BH-galaxy is <u>assumed</u>.

- ★ BH Formation (Booth & Schaye 09; also Springel, Di Matteo+ 05, Sijacki+07)
 - ★ Regularly run a friends-of-friends group finder on all of the dark matter particles during the simulation, get DM halo mass.
 - ★ If the halo is >M_{halo,min}=100m_{DM} (~10¹⁰M_☉) and does not already contain a BH, then its most gravitationally bound baryonic particle is converted into a collisionless BH particle with $M_{seed} = 0.001 m_g$ (~10⁵M_☉).
- ***** Growth
- ★ Feedback

The first BHs?

- * The first chemical enrichment is driven by ~10-50 M_{\odot} stars (faint SNe), leaving ~10 M_{\odot} BHs.
 - ★ No signature of ~140-300M_☉ stars (PISNe) observed in any chemical abundances.
- ★ If the accretion is suppressed, <100M_☉ stars form (Hosokawa+ 11). Otherwise, >300M_☉ stars <u>may</u> form (Ohkubo+ 09), which collapse to 100-1000M_☉ BHs.
 - ★ Direct collapse of gas can form ~10⁵ M_☉ BHs (e.g., Madau & Rees 01), but rare.
- What is the role of these "first" BHs in galaxy formation?
- * Can they grow into super-massive BHs at present?
- * Are they enough for the down-sizing?

New AGN model

★ BH Formation - not merging products, but originated from primordial SF

 $ho_{
m g} >
ho_{
m c}, \ Z=0$ then $m M_{
m seed}$ (~1000 $m M_{\odot}$)

Growth - accretion & mergers (same as in previous works; Springel, Di Matteo+ 05, Booth & Schaye 09)

$$egin{aligned} \dot{M}_{
m BH} &= (1-\epsilon_{
m r}) imes \min(\dot{M}_{
m acc},\dot{M}_{
m Edd}) \ \dot{M}_{
m acc} &= lpha rac{4\pi G^2 M_{
m BH}^2
ho}{(c_{
m s}^2+v^2)^{3/2}}, \ \dot{M}_{
m Edd} &= rac{4\pi G M_{
m BH} m_{
m p}}{\epsilon_{
m r} \sigma_{
m T} c} \end{aligned}$$

***** Feedback – thermal

$$E_{\rm FB} = \epsilon_{\rm r} \epsilon_{\rm f} \dot{M}_{\rm acc} c^2 \Delta t$$

Cosmological Simulations

CK, Springel, White 2007 Taylor & CK 2014, 2015abcd





Cosmological Simulation t = 0.39 Gyr, z = 11.72

Star



Constraining parameters

from cosmic SFRs, Magorrian relation, and galaxy size-mass relation $\rightarrow \alpha = 1$, $\epsilon_f = 0.25$, $\rho = 0.1$, $M_{seed} = 1000 M_{\odot}$



Taylor & CK 2014, MNRAS, 442, 2751

Observational Constraints



Simulation Outcome

Not used for parameter determination



Other Big Simulations

- Illustris w AREPO (106.5Mpc)
 - ★ Size-Mass relation?



Schaye et al. 14





AGN-driven outflows

Taylor & CK 2015b, MNRAS, 452, L59



★ AGN-driven winds remove the remaining gas (3% of baryon) and ~2% of total produced metals.

Metallicity of galaxies



Star Forming Galaxies

- Emission line ratios [OII],Hβ,
 [OIII],Hα,[NII],[SII]
- Photoionization models MAPPINGS (Sutherland & Dopita 199), Cloudy (Ferland et al. 1998)
- Stellar population synthesis
 models Starburst99 (Leitherer et al. 1999)
- Degeneracy between Z and ionization parameter q
- Solar Abundance = 8.69 (Allende Prieto et al. 2001), 8.66 (Asplund et al. 2004)

al. 2004)



Fig. 2.—Robust best-fit *M-Z* relations calculated using the different metallicity calibrations listed in Table 1, except the T_e method. The top panel shows the rms scatter in metallicity about the best-fit relation for each calibration in 0.1 dex bins of stellar mass. The *y*-axis offset, shape, and scatter of the *M-Z* relation differ substantially, depending on which metallicity calibration is used.

Early-type Galaxies (ETG)

- ★ Integrated Absorption Lines → Lick indices Mg₂, Mg_b, Fe5270, Fe5335, Hβ, Hγ
- Population synthesis
 models Worthey 1994,
 Thomas, Maraston & Bender
 2003
- Broadening of velocity dispersion
- Contamination of emission lines
- * Age-Metallicity Degeneracy

Metallicity [M/H]



[X/Fe]-[Fe/H] relations of ETGs



Abundances at high-redshifts



The origin of MZR



Gas in galaxies, identified by FOF Wind Gas, is not in galaxy, but has been in some galaxies $M_*/M_b=0.10$ $M_{g,gal}/M_b=0.10$ $M_w/M_b=0.20$

CK+ 07

The origin of MZR



- * Metals are effectively ejected from (present) dwarfs.
- * The origin of the mass-metallicity relation --- Massdependent Galactic Winds.

CK+ 07

Mass Metallicity Relations (MZR) Taylor & CK 2015a, MNRAS, 448, 1835





Taylor & CK 2015, MNRAS, 448, 1835

Other models



Figure 9. Comparison between observed and predicted $[Z/H]_0$ -mass and $[\alpha/Fe]_0$ -mass relationships. Open black circles represent central values of our sample galaxies. The black solid line is the weighted least-squares linear fit to the data points. In the left plots, we plot the predictions from Arrigoni et al. (2009) (green shading). In the right plots, the models of Calura et al. (2009) are shown (blue shading).

- ★ Matteucci 1994, GCE
 - ★ Inverse wind?
- * Arrigoni et al. 10, SAM
 - ★ flat IMF (x=1.1)
 - ★ yields: Woosley & Weaver 95
 - SNIa: bimodal DTD (inconsistent with obs. in the solar neighborhood)
- ★ Calura & Menci 09, SAM
 - ★ SFR-dependent IMF
 - ★ yields: Woosley & Weaver 95
 - ★ SNIa: Matteucci & Greggio 86
- ★ Yates et al. 13, SAM
 - ★ Chabrier (03) IMF
 - ★ yields: Portinari+ 98
 - ★ SNIa: power-law DTD

Cosmic evolution



Cosmic evolution

*Time evolution of scaling relations (Size/metallicity/BH mass – galaxy mass)

→ JWST (>2018)

***Internal structures** (2D map of gas, stars, and elements) within galaxies

 \rightarrow IFU (Integral field unit)

Time evolution of M-Z relation



Time evolution of $\mathbf{M}_{\text{BH}}\text{-}\sigma$ relation

Taylor & CK 2016, ArXiv1608.06685



Time evolution of M_{BH} - σ relation

Taylor & CK 2016, ArXiv1608.06685



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Metallicity gradients

GRADIENTS OF ABSORPTION-LINE STRENGTHS IN ELLIPTICAL GALAXIES

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For example, if the Worthey (1994) calibration is adopted, one finds that the typical metallicity gradient of elliptical galaxies is $\Delta \log Z/\Delta \log r \simeq -0.3$. The radial gradient of metallic line strength can be naturally explained by the dissipative collapse picture. However, the measured gradients are less steep than those predicted by numerical simulations of the collapse model. For example, Larson's hydrodynamical simulations gave $\Delta \log Z/\Delta \log r \sim -0.35$ (Larson 1974a) and -1.0 (Larson 1975), and Carlberg's N-body simulations gave $\Delta \log Z/\Delta \log r \sim -0.5$ (Carlberg 1984). This discrepancy could be interpreted if mergers flatten an original gradient: indeed numerical simulations showed that the gradient in a disk galaxy should be halved after three successive mergers of galaxies with similar sizes (White 1980). However, both simulations of the dissipative collapse and successive mergers leave room for the improvement because some essential physical processes, such as star formation, thermal feedback from supernovae, and metal enrichment, were not taken into account.



Metallicity Gradients vs Mass



Measured in MAX(2re,10kpc), V-luminosity/SFR weighted

Future

* Galaxies consist of stars. The formation and evolutionary histories will be revealed from their chemical evolution ("g"alactic archaeology).

James Webb Space Telescope (JWST) http://www.jwst.nasa.gov

Now

* Meantime, let's update our knowledge of stellar physics and the Milky Way, from elemental abundances and kinematics of stars with precision spectroscopy.

