**Evolution and** Nucleosynthesis during the AGB phase

> **Dr. Amanda Karakas School of Physics & Astronomy** Monash University, Australia

> > The Helix Nebula - NGC 7293 O HUBBLESITE.org



I'm giving three lectures, which will be broken down into the following components:

- 1. Introduction some basics
- 2. Nucleosynthesis prior to the asymptotic giant branch (AGB) phase
- 3. The evolution and nucleosynthesis of AGB stars
- 4. The slow neutron capture process

#### **Asymptotic Giant Branch stars**



- The asymptotic giant branch is the last nuclear burning phase for stars with mass < 8-10Msun</li>
- AGB stars are cool (~3000 K) evolved giants, spectral types M, S, C
- It is during the AGB where the products of nucleosynthesis reach the stellar surface
- Many AGB stars are observed to be losing mass in dense outflows of material
- $\rightarrow$  Enriching the interstellar medium
- ➔ Progenitors of planetary nebulae
- ➔ Reviews by Herwig (2005, ARAA) and Karakas & Lattanzio (2014, PASA)

H-exhausted core

## **Mixing and mass loss**

- Convective mixing (dredge-up) mixes the products of nucleosynthesis from the (hot) interior to the surface.
- Mass loss removes the enriched envelope, expelling the products into the interstellar medium.
- → When does most of the mass loss occur? When does the most nucleosynthesis occur?

For low and intermediate-mass stars, that is during the asymptotic giant branch (AGB)

## Where mixing takes place



## **Products of nucleosynthesis**

Low and intermediate-mass stars go through central hydrogen and helium burning

During the AGB, they have shells burning H and He

- 1. First dredge-up: Products of (partial) H burning
- 2. Second dredge-up: Products of H burning
- 3. Third dredge-up: Products of H, He-burning and neutroncapture nucleosynthesis
- 4. Hot bottom burning: Products of H-burning
- 5. Extra mixing processes: Products of H-burning
- $\rightarrow$  We we will now discuss the AGB phase of evolution



### **He-shell instabilities**

- The He-shell thins as the star ascends the AGB and becomes thermally unstable
- He-burning in a thin shell leads to a thermal runaway, similar to the core He-flash
- Why?
- Not caused by electron degeneracy, although the shell is partially degenerate
- Caused by the shell being thin
- Contracting shell → hotter → ε ∝ T<sup>40</sup> → but shell can't expand enough to cool → thermal runaway
- Luminosities can reach > 10<sup>8</sup> solar luminosities

## **He-shell burning in AGB stars**

- Up to  $\sim 10^8$  Lsun can be generated by a thermal pulse
- Energy goes into expanding the star
- He-shell becomes unstable to convection → mixes products of Heburning throughout shell



#### 2Msun, Z = 0.014 model star:

#### Intershell convection during thermal pulses

- The enormous amount of energy drives a convective region in the intershell
- Extends over almost the whole intershell
- Homogenises abundances within this region
- The mass of the pocket ~ few 10<sup>-2</sup> Msun, depending on the stellar mass
- The duration of convection is ~few hundred years
- Composition: result of partial He-burning: ~70%
   <sup>4</sup>He, ~25 <sup>12</sup>C and ~5% other stuff (<sup>22</sup>Ne, <sup>16</sup>O etc)

Convection zones = green, radiative = pink



Results for a 1.9Msun, Z = 0.008 model Model number proxy for time

## The thermal pulse cycle

Deep convective envelope

He-rich intershell

He-burning shell

Carbon-Oxygen core



#### Interpulse phase

## **The AGB Evolution Cycle**

- 1. On phase: He-shell burns brightly, producing up to  $10^8 L_{sun}$ , drives a convection zone in the He-rich intershell and lasts for ~ 100 years
- 2. Power-down: He-shell dies down, energy released by flash drives expansion which extinguishes the H-shell
- 3. Third dredge-up: convective envelope moves inward into regions mixed by flash-driven convection. Mixes partially He-burnt material to surface.
- Interpulse: star contracts and H-shell is re-ignited, provides most of the surface luminosity for the next ~10<sup>5</sup> years

Pulse (He-burning)  $\rightarrow$  TDU (mixing)  $\rightarrow$  Interpulse

Few ~10<sup>2</sup> yrs  $\rightarrow$  ~10<sup>2</sup> years  $\rightarrow$  ~10<sup>5</sup> yrs

#### Let's look at a thermal pulse again

Extent of convective pocket is 1.7 x 10<sup>-2</sup> Msun About half gets mixed into envelope



 $22^{nd}$  thermal pulse for the 3Msun, Z = 0.02 model

### The importance of the third dredge-up

- The third dredge-up determines how much He-shell material is mixed from the core to envelope
- Mass loss determines the number of thermal pulses
- So the combination (depth of dredge-up and mass loss rate) determine the role that AGB stars play in the evolution and origin of elements in the Universe!!

## Third dredge-up

- Badly named, can re-occur after each thermal pulse
- Inward movement of convective envelope, reaches into the He-shell
- Right-hand panel shows the evolution of the core in a low-mass AGB model
- Six (third)-dredge-up events are visible. Each one will mix He-shell material to the surface



#### **Non-energetic reactions**

- He-burning occurs in the *ashes* of H-burning
- The composition is typically 98% <sup>4</sup>He,  $\sim$ 2% <sup>14</sup>N
- Remember that the CNO cycle produces mostly <sup>14</sup>N, which can capture alpha particles to produce secondary nuclei, depending on T:
  - ${}^{14}$ N(α, γ)<sup>18</sup>F(β<sup>+</sup>ν)<sup>18</sup>O(α, γ)<sup>22</sup>Ne
  - <sup>22</sup>Ne +  $\alpha \rightarrow$  <sup>25,26</sup>Mg (+n or  $\gamma$ ) when T > 300 million K
- These reactions produce little energy but are important for nucleosynthesis
- Example, the <sup>22</sup>Ne(α,n)<sup>25</sup>Mg (Q = -0.478MeV) reaction releases *free* neutrons that can be used to produce heavy elements i.e., <sup>56</sup>Fe(n,γ)<sup>57</sup>Fe(n,γ)...

#### **Fluorine production**

- It's complicated! (e.g., Lugaro et al. 2004)
- The reaction chain:  ${}^{18}O(p, \alpha){}^{15}N(\alpha, \gamma){}^{19}F(\alpha, p){}^{22}Ne$
- Fluorine production takes place in the He-intershell: This is a region rich in <sup>4</sup>He, <sup>12</sup>C
- There are almost no protons or <sup>15</sup>N
- These are created by other reactions including:
  - $^{14}N(\alpha,\,\gamma)^{18}F(\beta^{+})^{18}O$  main reaction to produce  $^{18}O$
  - ${}^{13}C(\alpha, n){}^{16}O$  produces free neutrons (also for the s-process)
  - ${}^{14}N(n, p){}^{14}C$  produces free protons
  - <sup>18</sup>F( $\alpha$ , p)<sup>21</sup>Ne new, alternative proton production
  - ${}^{14}C(\alpha, \gamma){}^{18}O$  alternative reaction
  - $^{18}O(\alpha,\,\gamma)^{22}Ne$  main  $^{18}O$  destruction reaction
  - <sup>15</sup>N(p,  $\alpha$ )<sup>12</sup>C destroys <sup>15</sup>N

From a nucleosynthesis point of view:

- The triple alpha and <sup>12</sup>C(α, γ)<sup>16</sup>O reactions convert <sup>4</sup>He into <sup>12</sup>C and <sup>16</sup>O
- Secondary reactions can produce <sup>18</sup>O, <sup>19</sup>F, <sup>22</sup>Ne, <sup>25</sup>Mg, <sup>26</sup>Mg
- Final composition depends on temperatures, densities, and the duration of burning
- Secondary reactions can produce free neutrons (e.g., <sup>13</sup>C(a,n)<sup>16</sup>O, <sup>22</sup>Ne(a,n)<sup>25</sup>Mg) which drives the s-process

#### **Products of He-shell nucleosynthesis**

3Msun, Z = 0.014:

Surface abundance of carbon (left) and fluorine (right) during the AGB

 $\rightarrow$  We can make a carbon-rich star, which has C/O > 1



#### Mass range of carbon stars?

• From Karakas (2014) for [Fe/H] = -0.3, 0.0, +0.3



### **Third dredge-up uncertainties**

- It is important to know if the models are providing an accurate description of mixing in real AGB stars
- Because the third dredge-up determines how much Heshell material is mixed from the core to envelope
- Do current models predict enough TDU?
- Or too much?
- → Do the model predict the right mass and luminosity ranges for carbon stars?

### **Carbon star luminosity functions**

- Distances to the Magellanic Clouds are known
- Can derive accurate C-star luminosity functions
- These indicate that (most) stellar models do not predict enough dredge-up at low enough masses
- And it is deeper at these lowest masses than current models predict
- Can "force" the TDU in lowmass models...



Stancliffe, Izzard, & Tout (2005)

#### **Uncertainties: The amount of third dredge up**



#### 1.25Msun, Z = 0.01:

- Forcing dredge-up by extending the base of the envelope by N scaleheights
- e.g., Karakas et al. (2010); Frost & Lattanzio (1996)



Figure 4.17: Highlighting the strong effect of including overshoot on the AGB. The time evolution of the mass of the H-exhausted core is plotted. Both stars, having a mass of 2 M<sub> $\odot$ </sub> and a metallicity of [Fe/H]=-5.45, started with the same initial conditions except for the inclusion of overshoot in one (lower curve,  $f_{OS} = 0.01$ ). An enormous difference in core mass evolution is clearly seen. The model with no overshoot (upper curve) has virtually no 3dup whilst the model with overshoot initially has  $\lambda_{3dup} > 1$ . As the core mass is the primary factor in AGB evolution, the vastly different core masses represent a very large uncertainty in AGB evolution.

#### 2Msun, [Fe/H]= -5.45:

Diffusive mixing + Herwig's scheme for extending the envelope using exponentially decaying overshoot From Simon Campbell

## Hot bottom burning

Occurs in stars over about 4.5Msun for Z = 0.014

Along with thermal pulses and the third dredge-up, these stars also have:

- Second dredge-up: Biggest  $\Delta Y$  (up to 0.1)
- Hot bottom burning: Proton-capture nucleosynthesis at base of envelope (products: N, Na, AI)



#### Hot bottom burning and third dredge up

Example: 6Msun, Z = 0.02

Third dredge-up (TDU) and HBB act together

CN cycle is acting close to equilibrium for ~20 thermal pulses



 $^{12}C/^{13}C \sim 3$  is the equilibrium ratio

The C/O ratio never exceeds 1

#### Hot bottom burning and third dredge up

Looking at the surface abundances of Ne to Al as a function of metallicity:

- 6Msun, Z = 0.02 has a peak temperature of ~80 million K
- 6Msun, Z = 0.004 has a peak of ~95 million K



### Lithium production

- The first thing to happen is that <sup>7</sup>Li is produced via the Cameron-Fowler Beryllium Transport Mechanism
- This is basically pp chains plus convection!
- The idea is that lithium is made by  ${}^{3}\text{He}(\alpha, \gamma){}^{7}\text{Be}$
- and then to use convection to move the <sup>7</sup>Be away from the hot region before it can complete the ppll or pplll chains:

<sup>3</sup>He (
$$\alpha, \gamma$$
) <sup>7</sup>Be ( $\beta, \psi$ ) <sup>7</sup>Li  
<sup>3</sup>He ( $\alpha, \gamma$ ) <sup>7</sup>Be ( $\beta, \psi$ ) <sup>7</sup>Li  
<sup>5</sup>Be ( $p, \gamma$ ) <sup>8</sup>B ( $\beta^* \vee$ ) Be( $\alpha$ ) <sup>4</sup>He = PPIII  
BAD!  
BAD!

Cameron-Fowler mechanism

## **Lithium production**

Lithium is produced by the Cameron-Fowler mechanism: <sup>7</sup>Be is transported by convection, where it captures an electron to produce <sup>7</sup>Li



 $\log \epsilon(\text{Li})_{\text{max}} = \log_{10}(\text{Li/H}) + 12 = 4.5$ 

#### **Uncertainties caused by convection**



for three convective prescriptions

Surface CNO abundances as a function of total mass

From Ventura & D'Antona (2005)

## **Other mixing phenomena**

- What is the impact of non-convective extra mixing processes on AGB evolution and nucleosynthesis?
- Examples include: rotation, thermohaline or double diffusive mixing, mixing induced by internal gravity waves, magnetic fields...
- Effect on the stellar yields?

I won't have time to discuss these here **Reading:** Karakas & Lattanzio (2014, PASA review, arXiv:1405.0062)

## **Summary of nucleosynthesis**

- **C/O > 1**: ~1.5 to  $4.5M_{sun}$  for Z = 0.014 (solar)
  - Inward movement of convection mixes the products of He-shell nucleosynthesis to the envelope (<sup>12</sup>C,<sup>19</sup>F, *s*-process)
- **C/O < 1**: Above ~4.5M<sub>sun</sub> for Z = 0.014
  - Hydrogen burning reactions at base of convective envelope (e.g., <sup>7</sup>Li, <sup>13</sup>C, <sup>14</sup>N, <sup>23</sup>Na, <sup>26,27</sup>Al, s-process?)

#### **References: (focused on nucleosynthesis results)**

Busso, Gallino & Wasserburg (1999), Forestini & Charbonnel (1997), Straniero et al. (1997), Mowlavi (1999), Herwig (2000, 2005), Stancliffe & Jeffery (2007), Campbell & Lattanzio (2008), Suda & Fujimoto (2010), Cristallo et al. (2011, 2015), Wiess & Ferguson (2009), Marigo et al. (2013), Ventura et al. (2013), Cruz et al. (2013)

## Super-AGB stars: 8-10 Msun stars

- The first models of stars in the range 8 to 10Msun were by Nomoto (1984), Garcia-Berro & Iben (1994), Ritossa et al. (1996), and Gutierrez et al. (1996)
- The paper by Garcia-Berro & Iben (1994) gave the name "super-AGB" for stars that ignite carbon and then experience thermal pulses
- These calculations are difficult, and no one really worked on them for a long time after, until Gil-Pons et al. (2001, 2002) and then Siess (2006)

## **Off-centre carbon ignition**

- Stars between ~8 to 10Msun go through degenerate carbon ignition
- Before ascending the thermallypulsing AGB with O-Ne cores
- Q: What fraction explode as supernovae or leave massive white dwarfs?
- E.g., Poelarends et al. (2008), Gil-Pons et al. (2013), Jones et al. (2014)
- The brightest AGB stars in young populations, with Mbol ~ -7.6, brighter than the traditional AGB limit (Mbol ~ -7.1)



7.5Msun, Z= 10<sup>-4</sup> model by Siess (2007)

### Carbon ignition: 9Msun, Z = 0.02

- Maximum temperature peaks at ~950 x  $10^6$  K.
- Duration of carbon flashes and central burning ~30,000 years (model from Karakas et al. 2012)
- Carbon burning occurs during early AGB, while second dredge-up is occurring (e..g., Gil-Pons et al. 2005, Siess 2006)
- Dredge-up is deep, can eat into the He-burning shell



#### **Super-AGB stars**

A 9Msun, Z = 0.02 model has a core mass of ~1.18Msun. Too low to become an electron capture supernovae (from Karakas et al. 2012) It will produce an O-Ne white dwarf



#### **Nucleosynthesis in super-AGB stars**

 $7M_{sun}$ , Z = 0.002 (1/100<sup>th</sup> solar). Peak temperature ~ 140 x 10<sup>6</sup> K. This is about as extreme as it gets in an AGB star!



Recent models: Siess (2010) Pumo et al. (2008), Doherty et al. (2010) Karakas et al. (2012) Herwig et al. (2012) Ventura et al. (2012) Gil-Pons et at. (2013) Takahashi et al. (2013) Doherty et al. (2014a,b) Fishlock et al. (2014) Doherty et al. (2015) Shingles et al. (2015) Woolsey & Heger (2016) Jones et al. (2016)

## Final fate of Super-AGB stars?

# The final fate of super-AGB stars is uncertain

- → Will they mostly produce massive ONe white dwarfs
- → What fraction will explode as electron capture supernova?
- → What are their nucleosynthesis products? H burning? He-shell burning? The rapid neutron capture process?
- → What happens when they are in a binary system? Will more explode?
- → How do they affect the enrichment of the galaxy?

# Lots of questions! Very exciting stuff



From Doherty et al. (2015)

#### **Globular cluster abundances**

#### In a typical cluster:

- The abundances of C-N, O-Na and Mg-Al are anticorrelated (Gratton et al. 2009, 2012)
- Sum C+N+O ~ constant (within a factor of ~2)
- No variation of alpha, s or r-process elements from star-to-star *within* a cluster...
- Does this imply the composition has been exposed to hydrogen burning (CNO, NeNa, MgAI)
- For an alternative hypothesis see Bastian et al. (2015)
   In an atypical cluster: ~10%
- NGC 1851,  $\omega$  Cen, M22, NGC 2419
- Show variations in C+N+O, s-process, r-process (rare) and iron-peak elements (e.g, Marino et al. 2012)

### **The O-Na anti-correlation**

- Why is there a correlation between O and Na in some globular cluster stars?
- Seen in all globular clusters (e.g., Carretta et al. 2009)
- Now we think it is probably pollution when the stars we see now formed → But from what?



#### Field stars versus GC stars

[Na/Fe]

- Evolved fields stars of the same metallicity as globular cluster stars show correlations between C and N
- This is caused by CN processed material being mixed into the envelope by the first dredge-up and extra mixing
- But field stars do not show correlations between O, Na (e.g., Gratton et al. 2000)
- But we also see C-N variations on the MS in GCs (e.g., figure from Cannon et al. 1998; also Briley et al. 2004)



Figure 6. The 47 Tuc colour-magnitude diagram, using the same data as Fig. 2 but with the symbols of Fig. 4 to distinguish between the CN-strong and CN-weak stars.

### Pollution by what type of stars?

- 1. Deep mixing can explain the Li, C-N trends with luminosity in some GCs (e.g., Lind et al. 2009)
- 2. Self-pollution by AGB stars experiencing hot bottom burning (e.g., Ventura et al. 2009)
- 3. Self-pollution by slow winds from rapidly rotating massive stars (Decressin et al. 2007)
- 4. Binary massive stars (De Mink et al. 2009)
- 5. Very massive stars (~10,000Msun; Denissenkov & Hartwick 2014)

AGB stars have been favoured because their slow winds can be retained by the cluster, and they produce no metals

## **GC** chemical evolution



- AGB models with third dredge-up cannot match helium enrichments along with O-Na, Mg-Al composition of GCs (e.g., NGC 6752 shown above; Fenner et al. 2004, Karakas et al. 2006)
- But see recent chemical evolution models from D'Ercole et al. (2010, 2016) using AGB models from Ventura et al. (2013)