



Black Hole

Magnetospheres

Alexander (Sasha) Tchekhovskoy

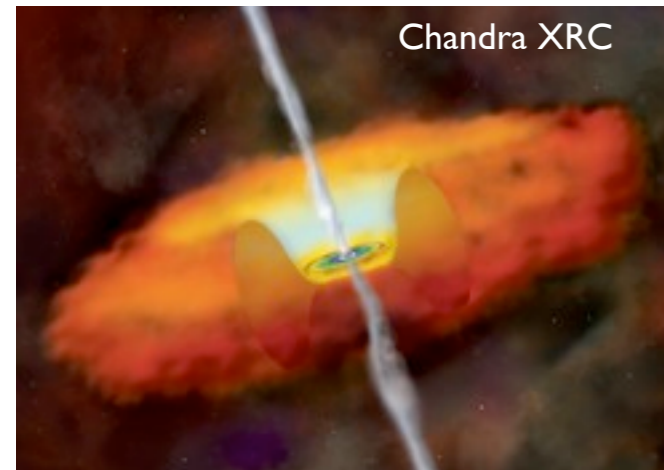
TAC Fellow
UC Berkeley

→ Northwestern University

Black Holes Power Many Transients

Supermassive

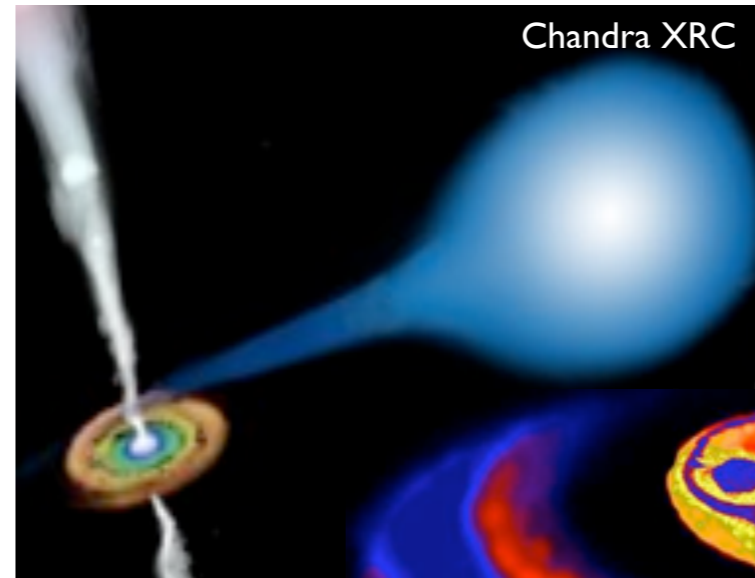
$$M \sim 10^6 - 10^9 M_{\odot}$$



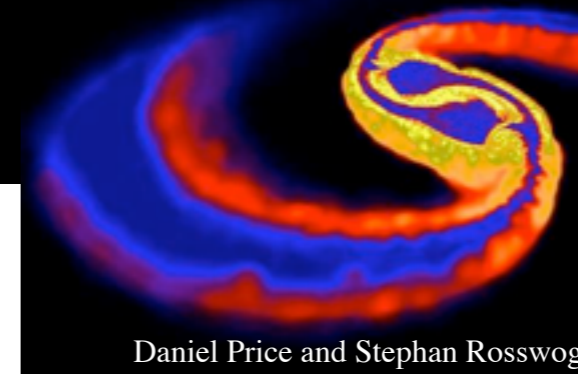
Quasars/AGN

Stellar-mass

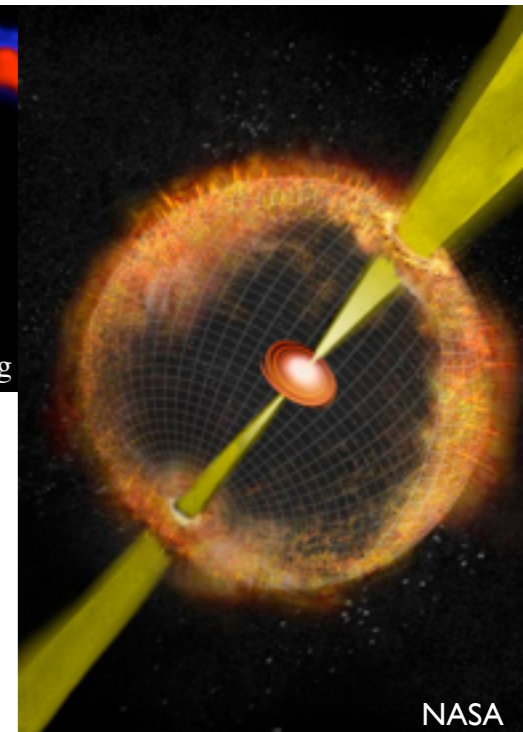
$$M \sim \text{few} - 10 M_{\odot}$$



Black
Hole
Binaries



Gamma-ray
bursts



Black hole or
Neutron star

Black Holes Power Many Transients

Supermassive

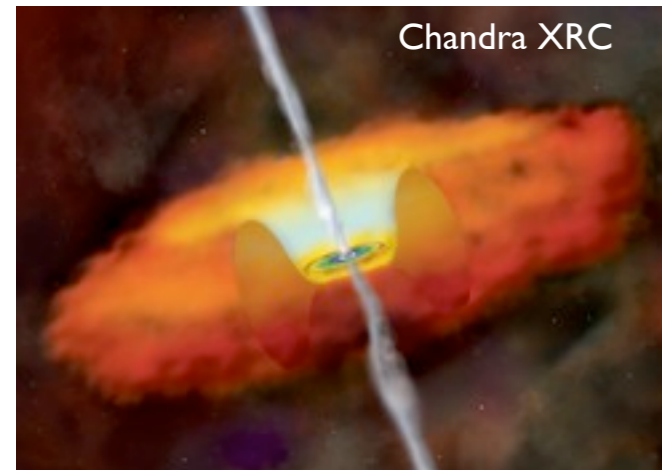
$$M \sim 10^6 - 10^{10} M_{\odot}$$

Intermediate

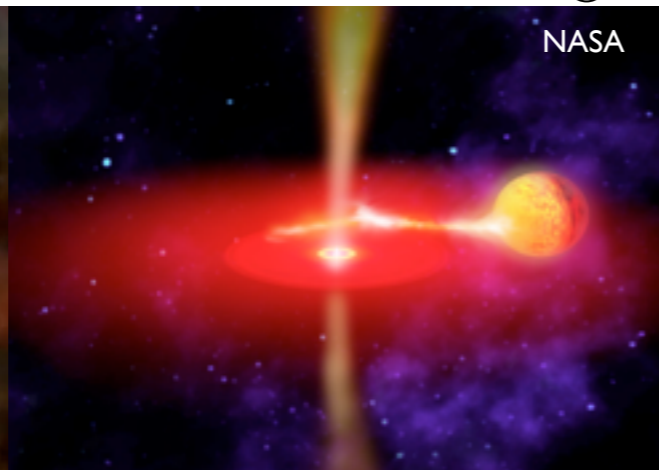
$$M \sim 10^2 - 10^5 M_{\odot}$$

Stellar-mass

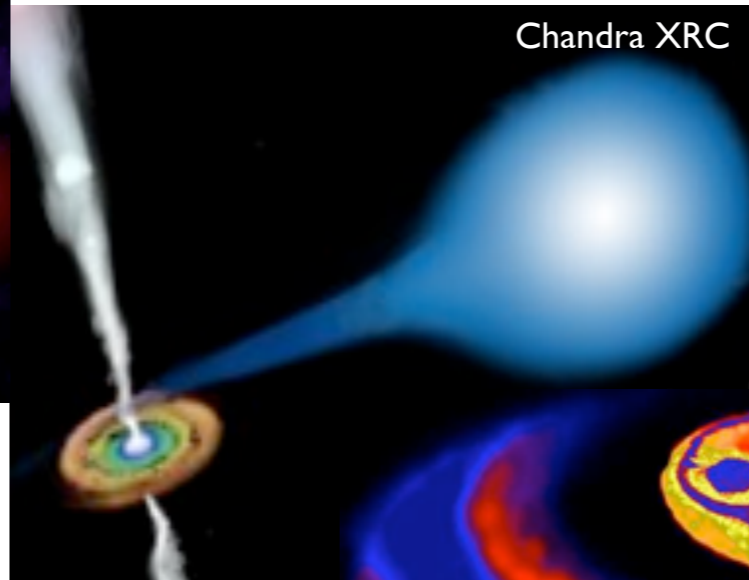
$$M \sim \text{few} - 10 M_{\odot}$$



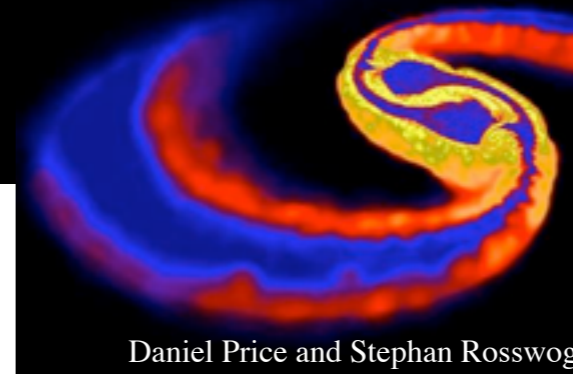
Quasars/AGN



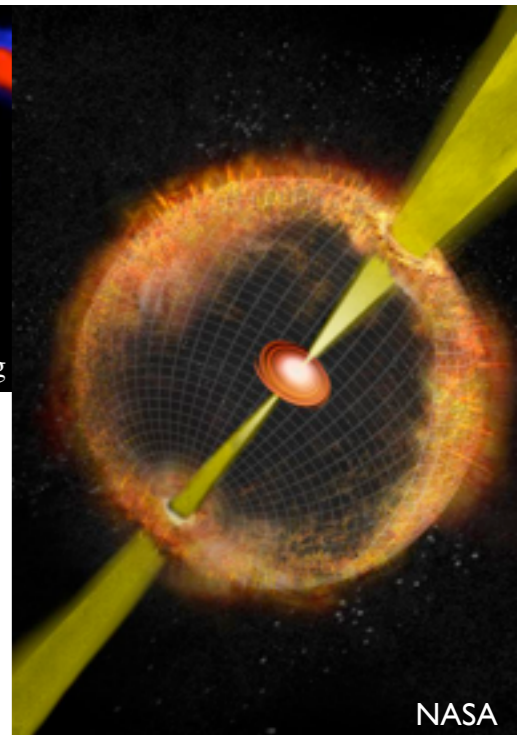
Intermediate-mass
black holes/ultra-
luminous X-ray
sources?



Black
Hole
Binaries



Gamma-ray
bursts



Black hole or
Neutron star

Black Holes Power Many Transients

Supermassive

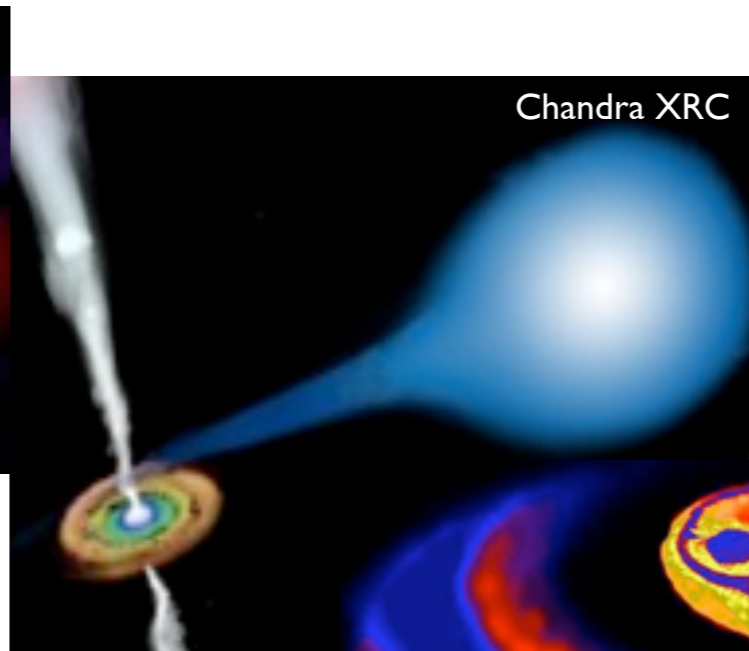
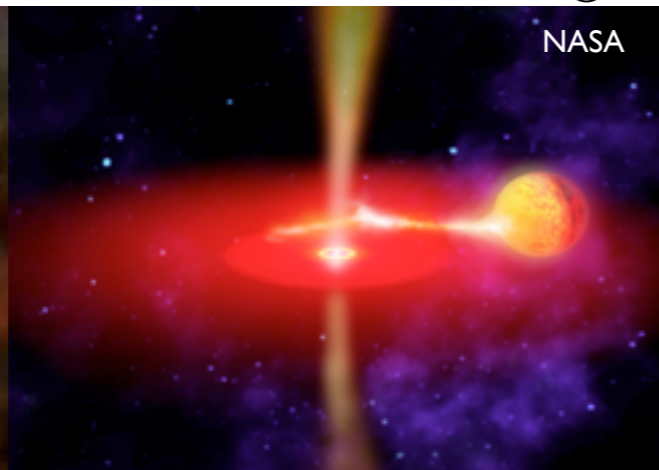
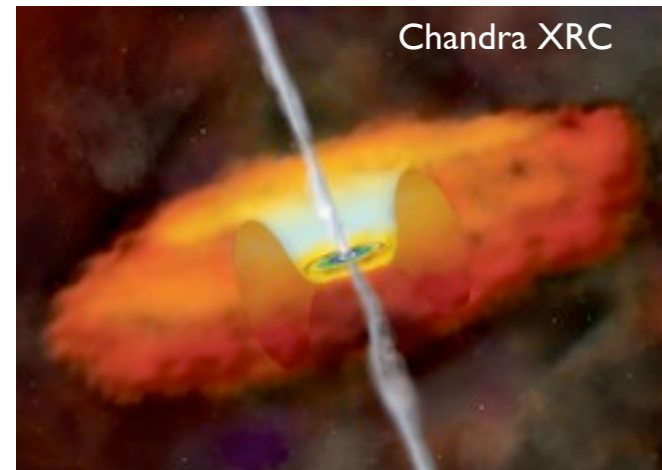
$$M \sim 10^6 - 10^{10} M_{\odot}$$

Intermediate

$$M \sim 10^2 - 10^5 M_{\odot}$$

Stellar-mass

$$M \sim \text{few} - 10 M_{\odot}$$

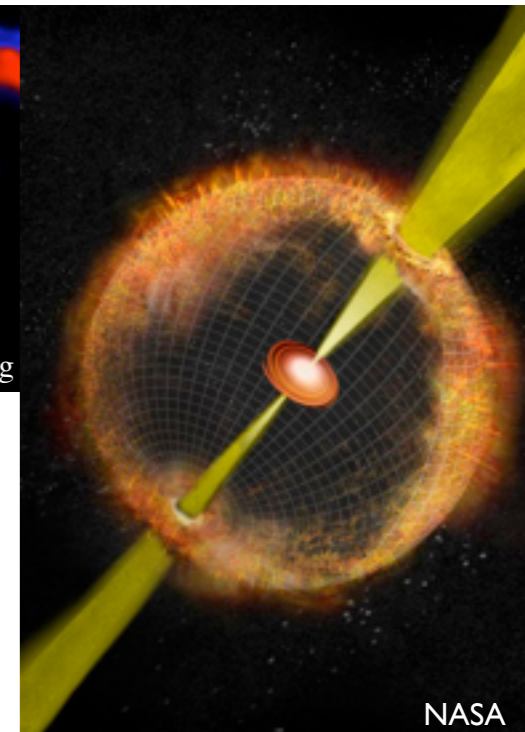
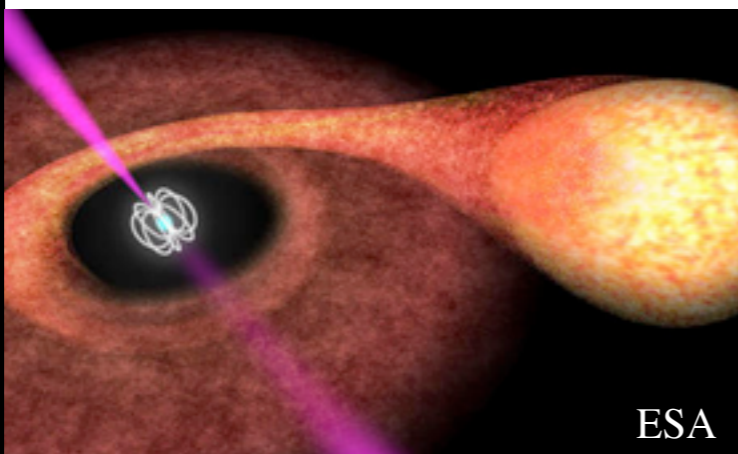
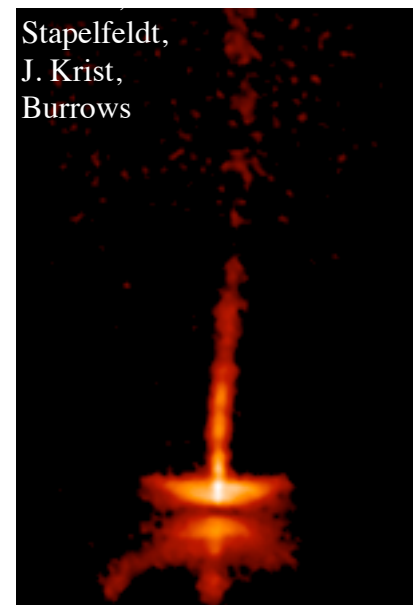


Quasars/AGN

Intermediate-mass
black holes/ultra-
luminous X-ray
sources?

Black
Hole
Binaries

Gamma-ray
bursts



Stars

Neutron Stars, White
Dwarfs; $M \sim M_{\odot}$

Black hole or
Neutron star

(AT 2015)

Black Holes Power Many Transients

Supermassive

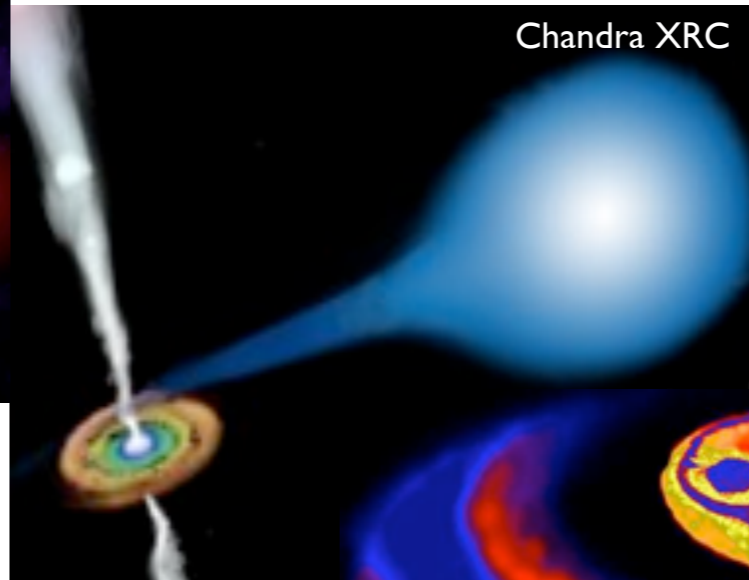
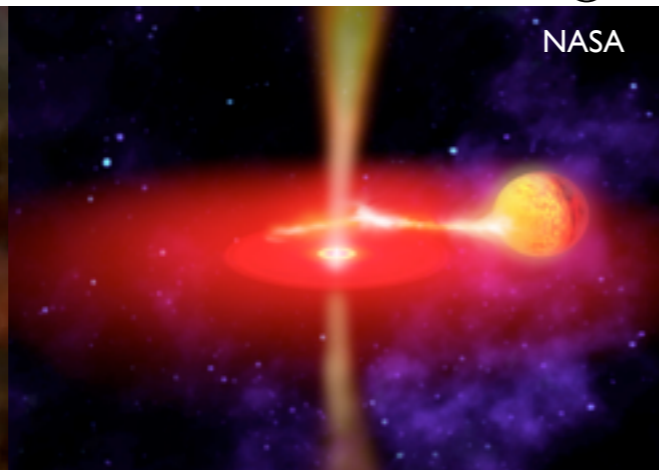
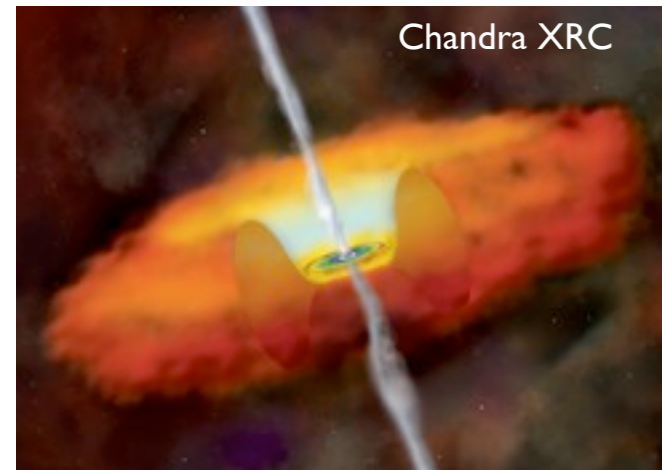
$$M \sim 10^6 - 10^{10} M_{\odot}$$

Intermediate

$$M \sim 10^2 - 10^5 M_{\odot}$$

Stellar-mass

$$M \sim \text{few} - 10 M_{\odot}$$

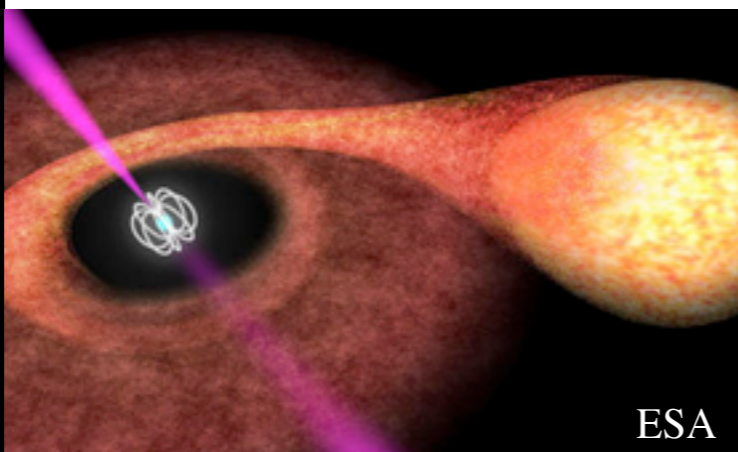
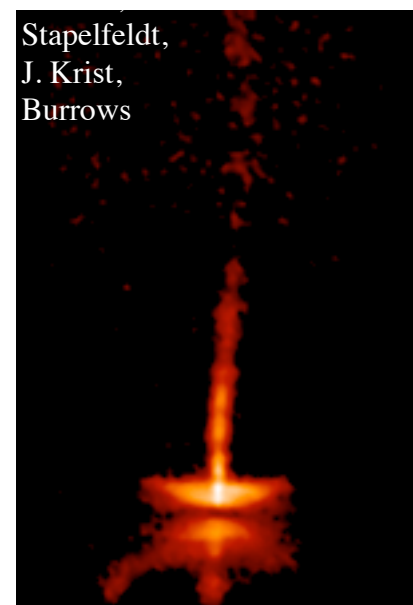


Quasars/AGN

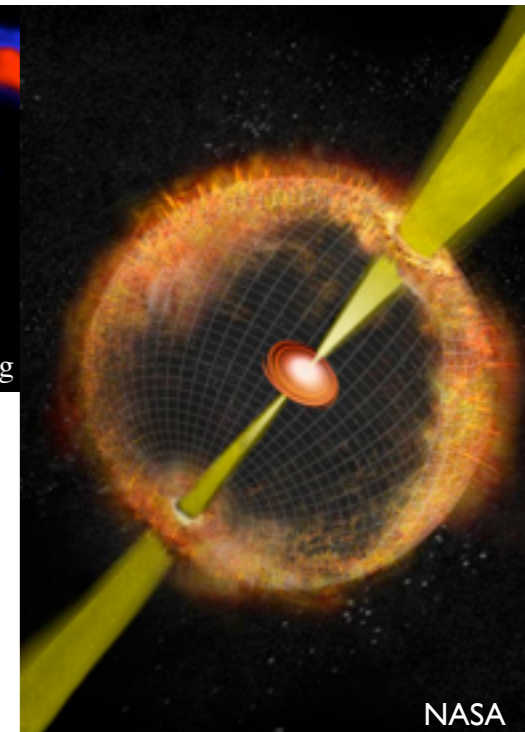
Intermediate-mass
black holes/ultra-
luminous X-ray
sources?

Black
Hole
Binaries

Gamma-ray
bursts



GWs from
binary black
holes



Black hole or
Neutron star

(AT 2015)

Stars

Neutron Stars, White
Dwarfs; $M \sim M_{\odot}$

Black Holes Power Many Transients

Supermassive

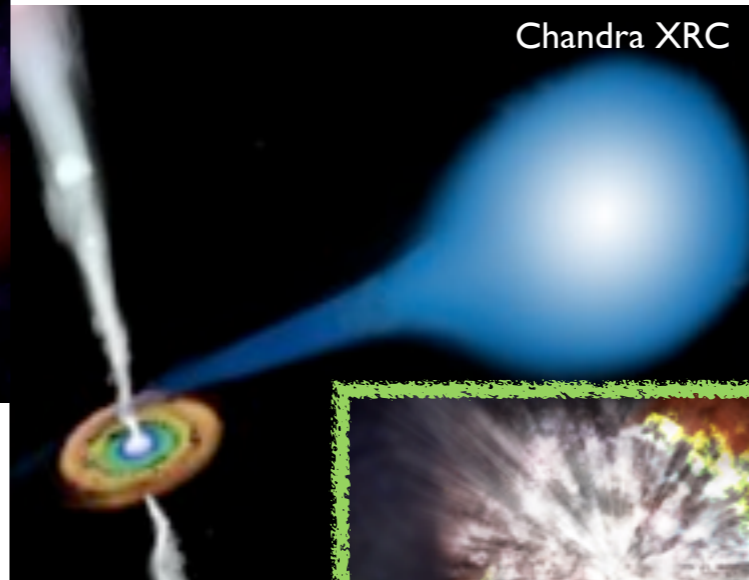
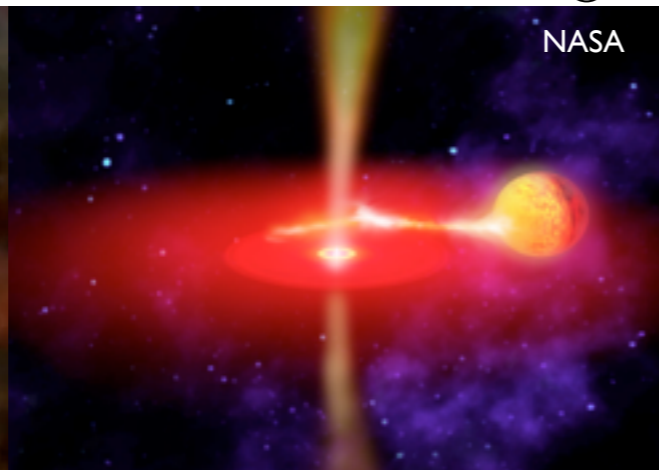
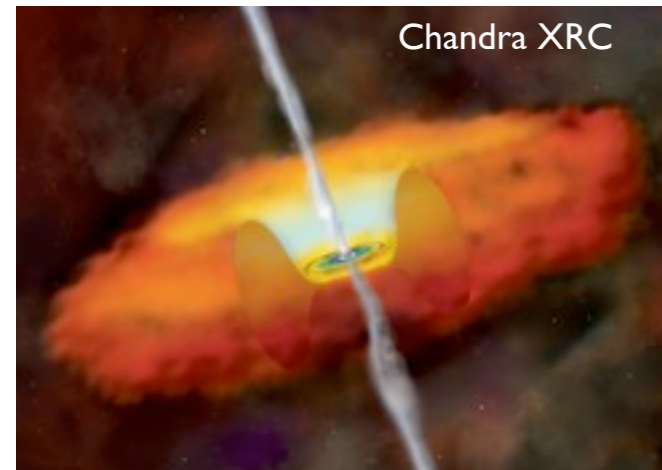
$$M \sim 10^6 - 10^{10} M_{\odot}$$

Intermediate

$$M \sim 10^2 - 10^5 M_{\odot}$$

Stellar-mass

$$M \sim \text{few} - 10 M_{\odot}$$



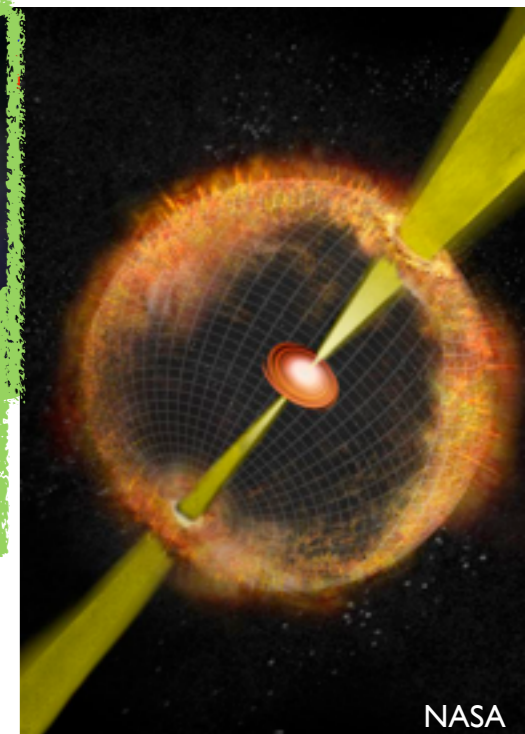
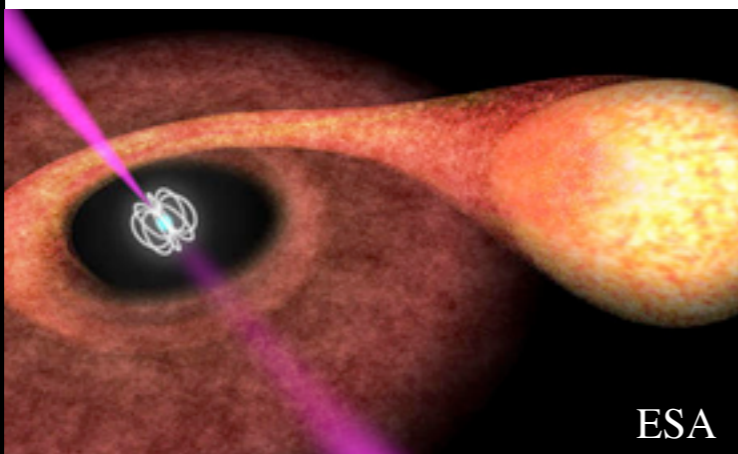
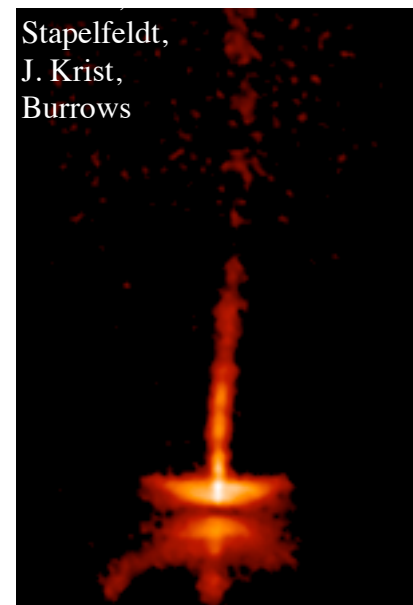
Quasars/AGN

Intermediate-mass
black holes/ultra-
luminous X-ray
sources?

Black
Hole
Binaries

factories of
heavy elements,
"kilonovae"

bursts



Black hole or
Neutron star

(AT 2015)

Stars

Neutron Stars, White
Dwarfs; $M \sim M_{\odot}$

Black Holes Power Many Transients

Supermassive

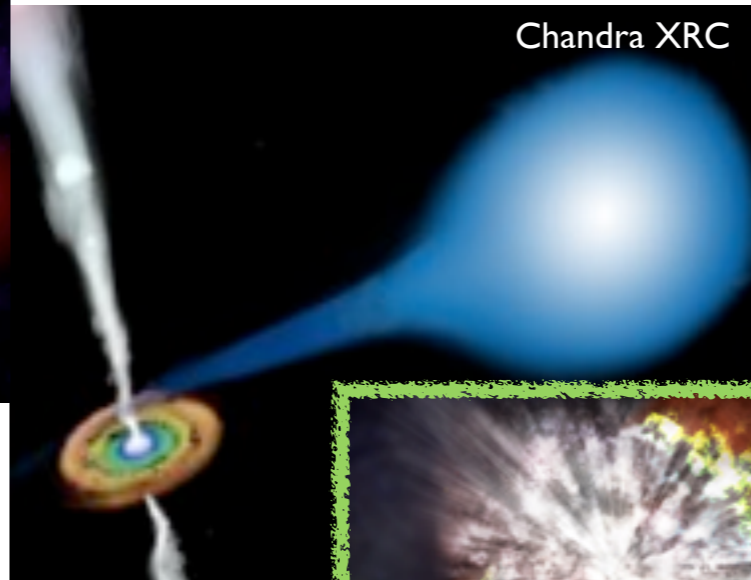
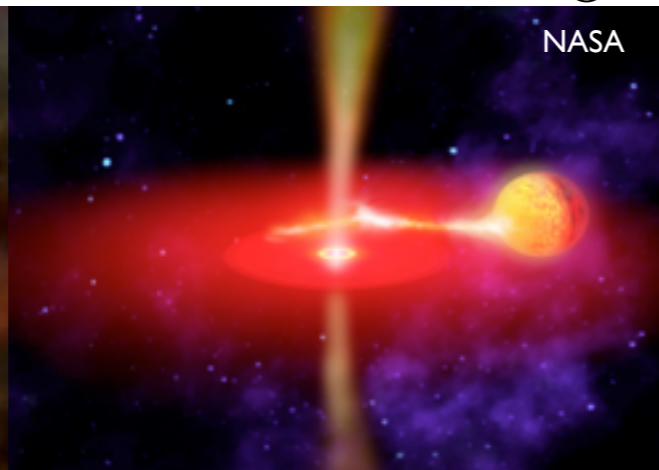
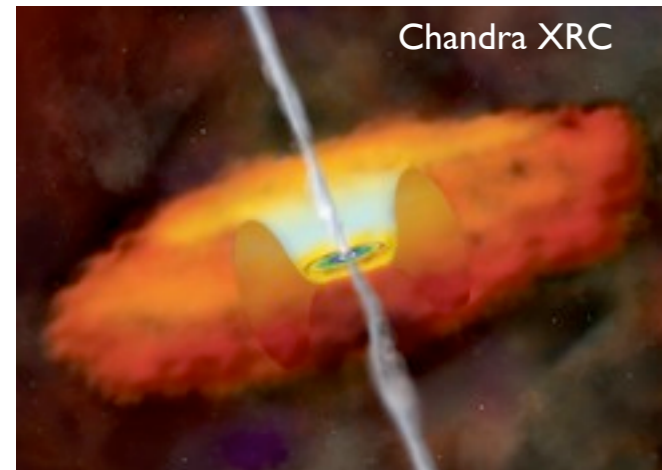
$$M \sim 10^6 - 10^{10} M_{\odot}$$

Intermediate

$$M \sim 10^2 - 10^5 M_{\odot}$$

Stellar-mass

$$M \sim \text{few} - 10 M_{\odot}$$



Quasars/AGN

Intermediate-mass
black holes/ultra-
luminous X-ray
sources?

Black
Hole
Binaries

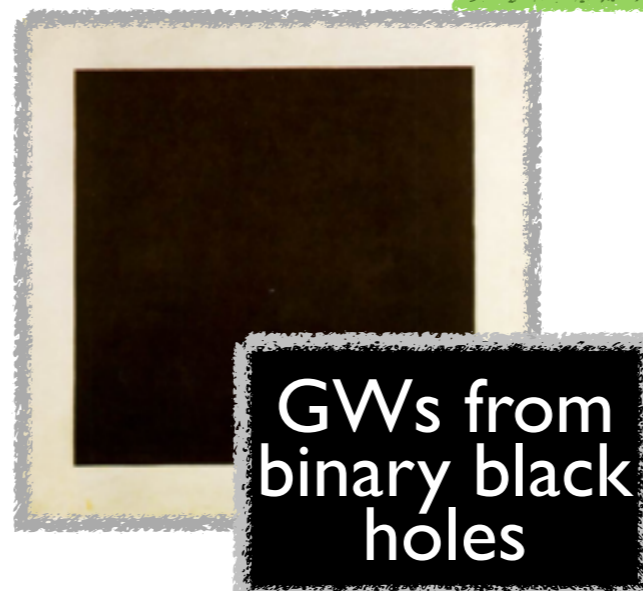
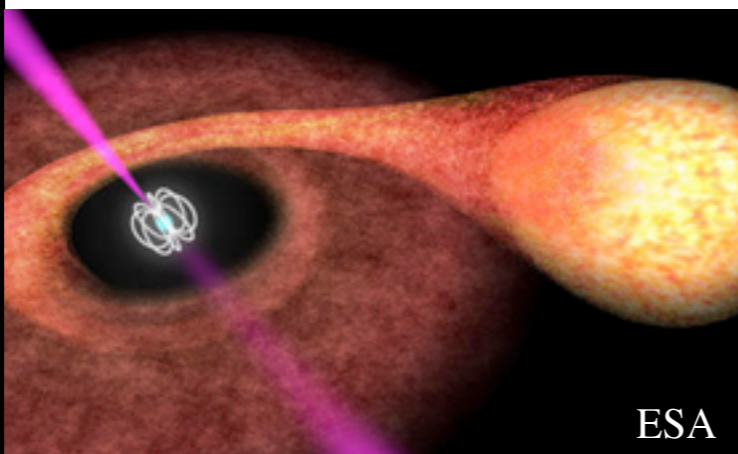
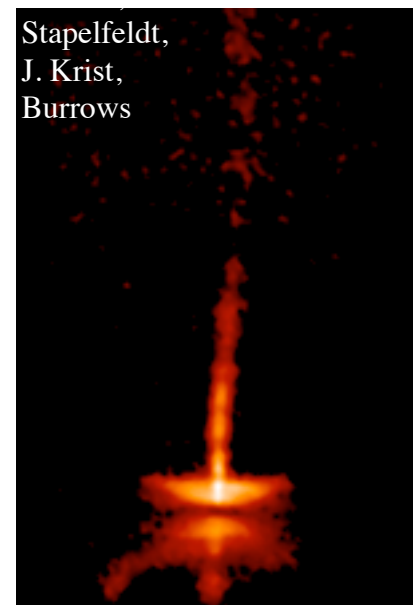
factories of
heavy elements,
“kilonovae”

bursts

explode stars
magnetically,
“superlumi-
nous SNe”

Neutron star

(AT 2015)



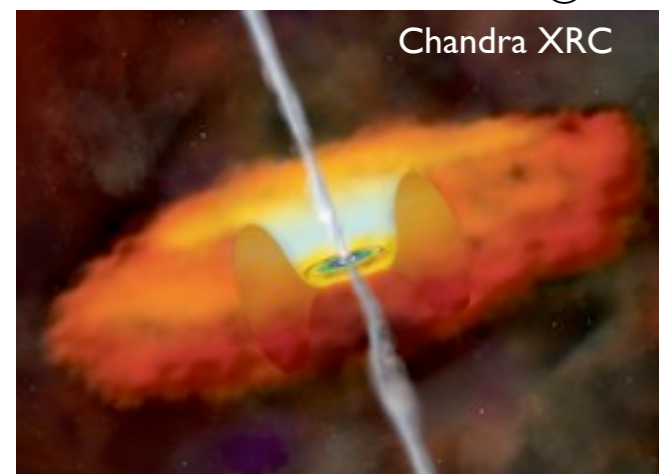
Stars

Neutron Stars, White
Dwarfs; $M \sim M_{\odot}$

Black Holes Power Many Transients

Supermassive Intermediate ————— Stellar-mass

$M \sim 10^6 - 10^{10} M_{\odot}$ $M \sim 10^2 - 10^5 M_{\odot}$ $M \sim \text{few} - 10 M_{\odot}$

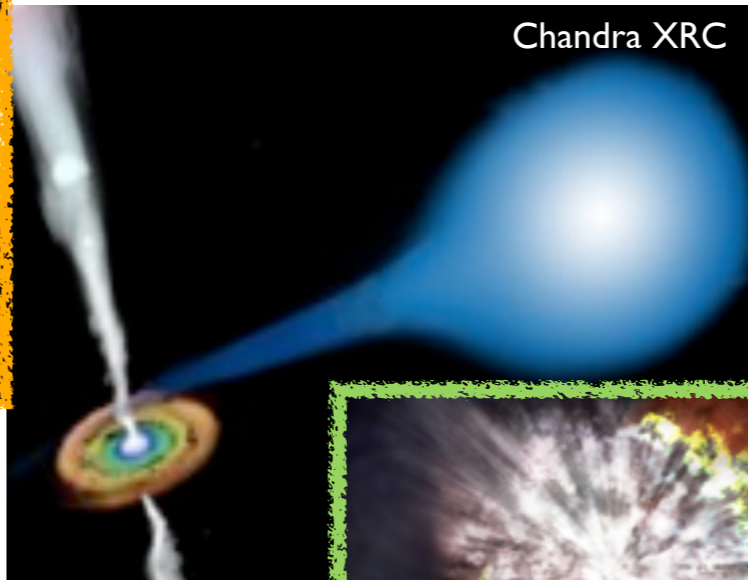


Quasars/AGN

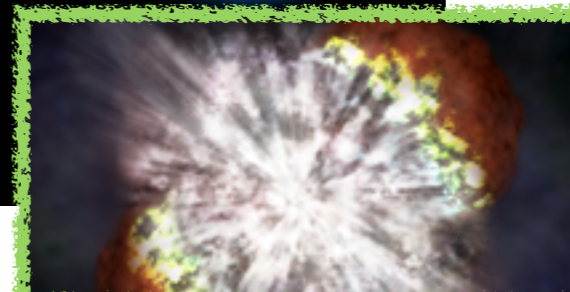


Intermediate-mass
black holes/ultra-
luminous X-ray
sources?

turn out to be
neutron stars

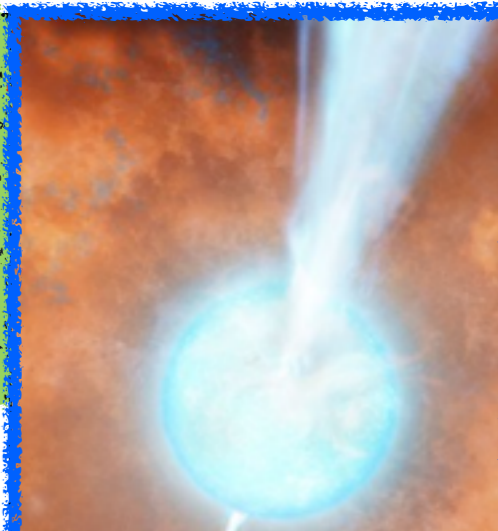


Black
Hole
Binaries



factories of
heavy elements,
“kilonovae”

bursts



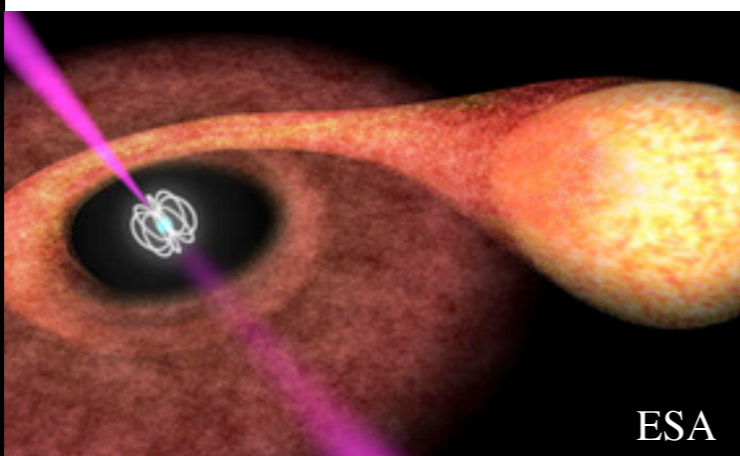
explode stars
magnetically,
“superlumi-
nous SNe”

Neutron star

(AT 2015)



Stars



Neutron Stars, White
Dwarfs; $M \sim M_{\odot}$



GWs from
binary black
holes

Black Holes Power Many Transients

Supermassive Intermediate ————— Stellar-mass

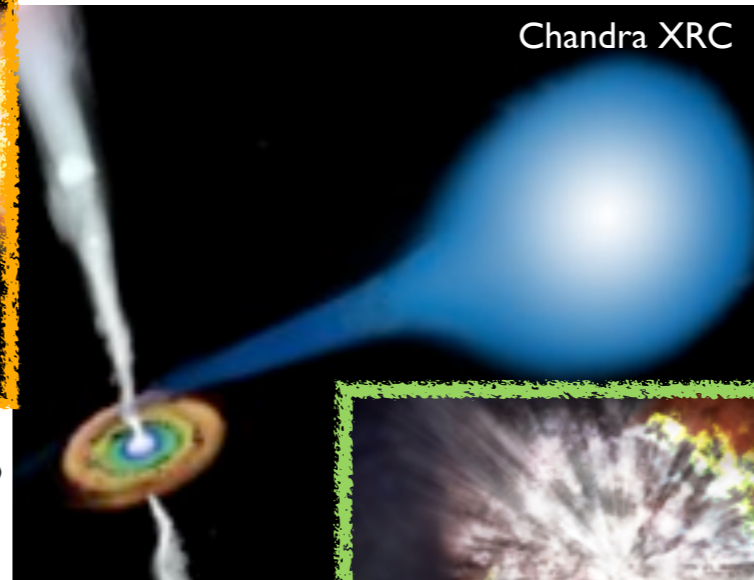
$M \sim 10^6 - 10^{10} M_{\odot}$ $M \sim 10^2 - 10^5 M_{\odot}$ $M \sim \text{few} - 10 M_{\odot}$



Quasars/AGN



Intermediate-mass
black holes/ultra-
luminous X-ray
sources?



Black
Hole
Binaries



bursts

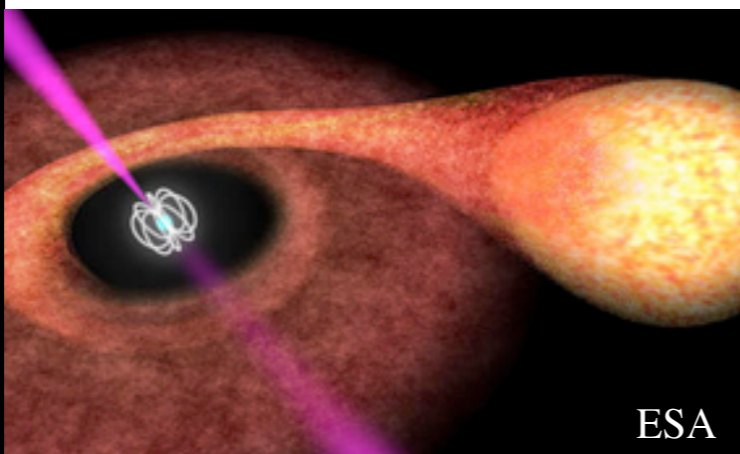


Neutron star

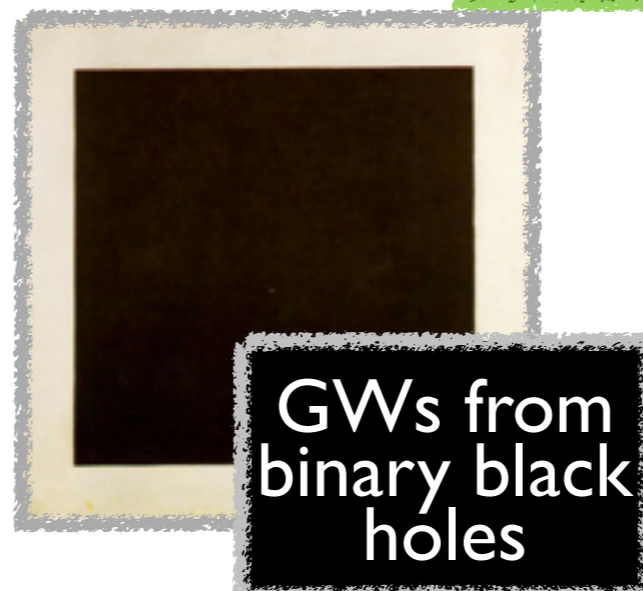
(AT 2015)



Stars

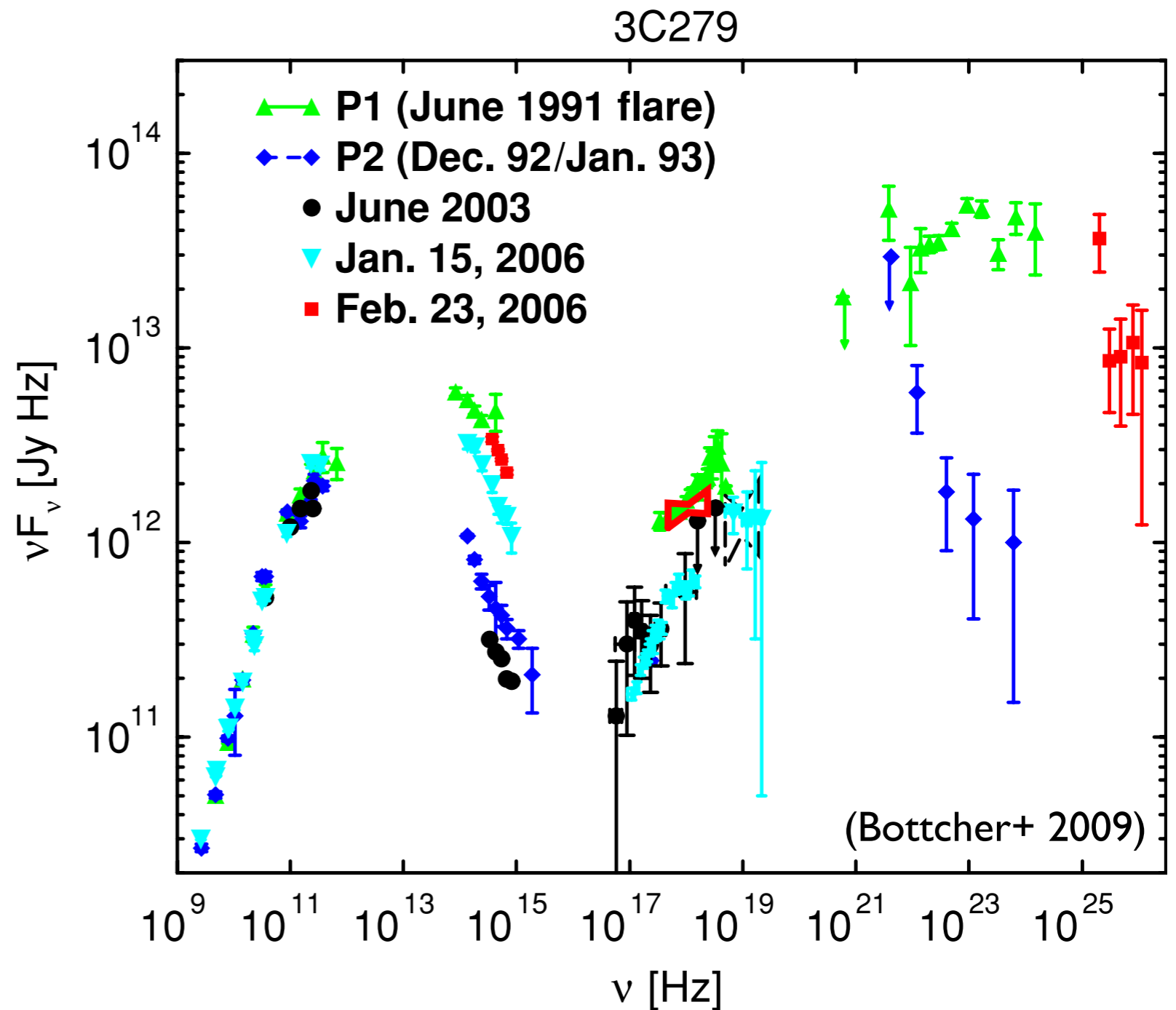
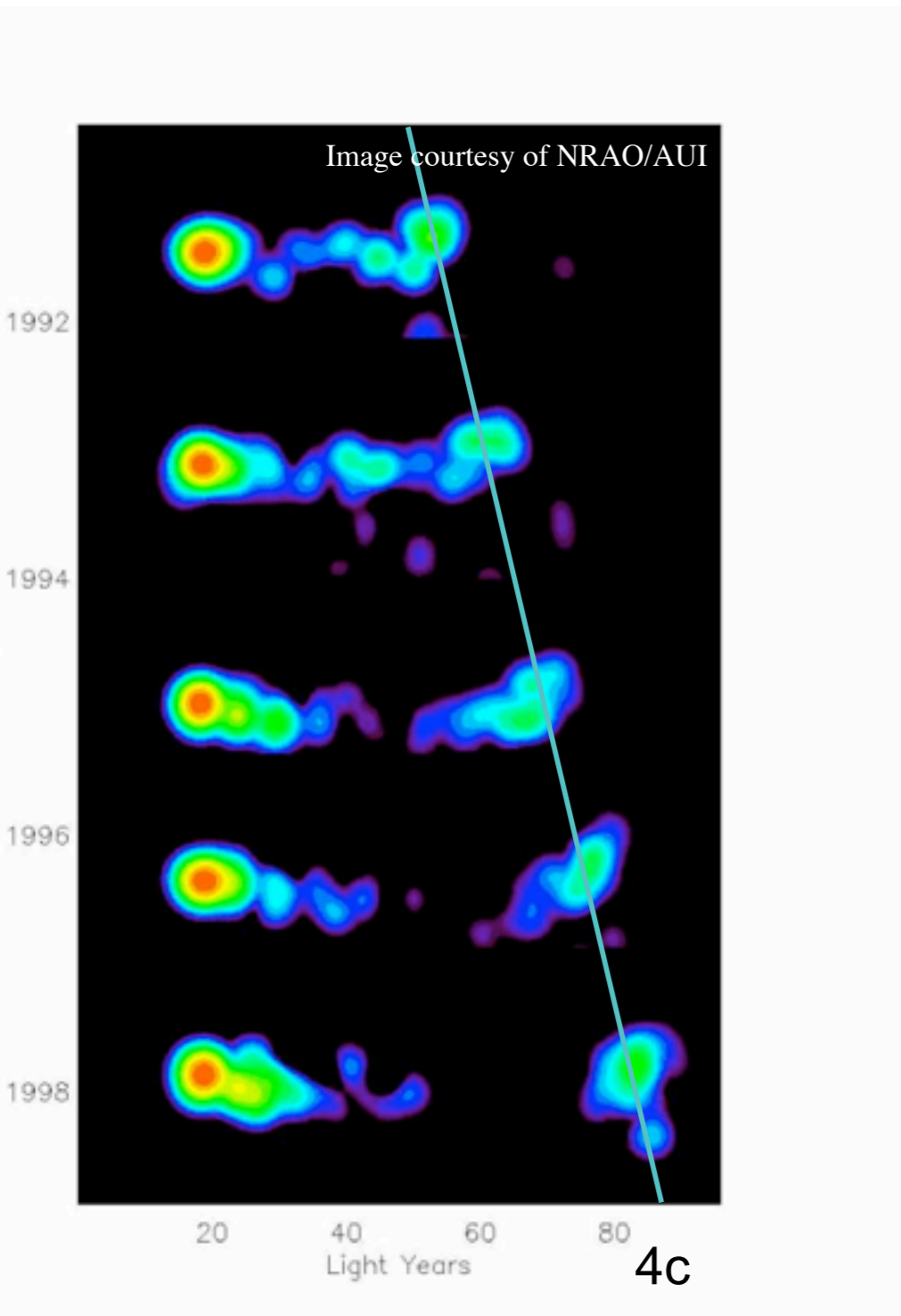


Neutron Stars, White
Dwarfs; $M \sim M_{\odot}$



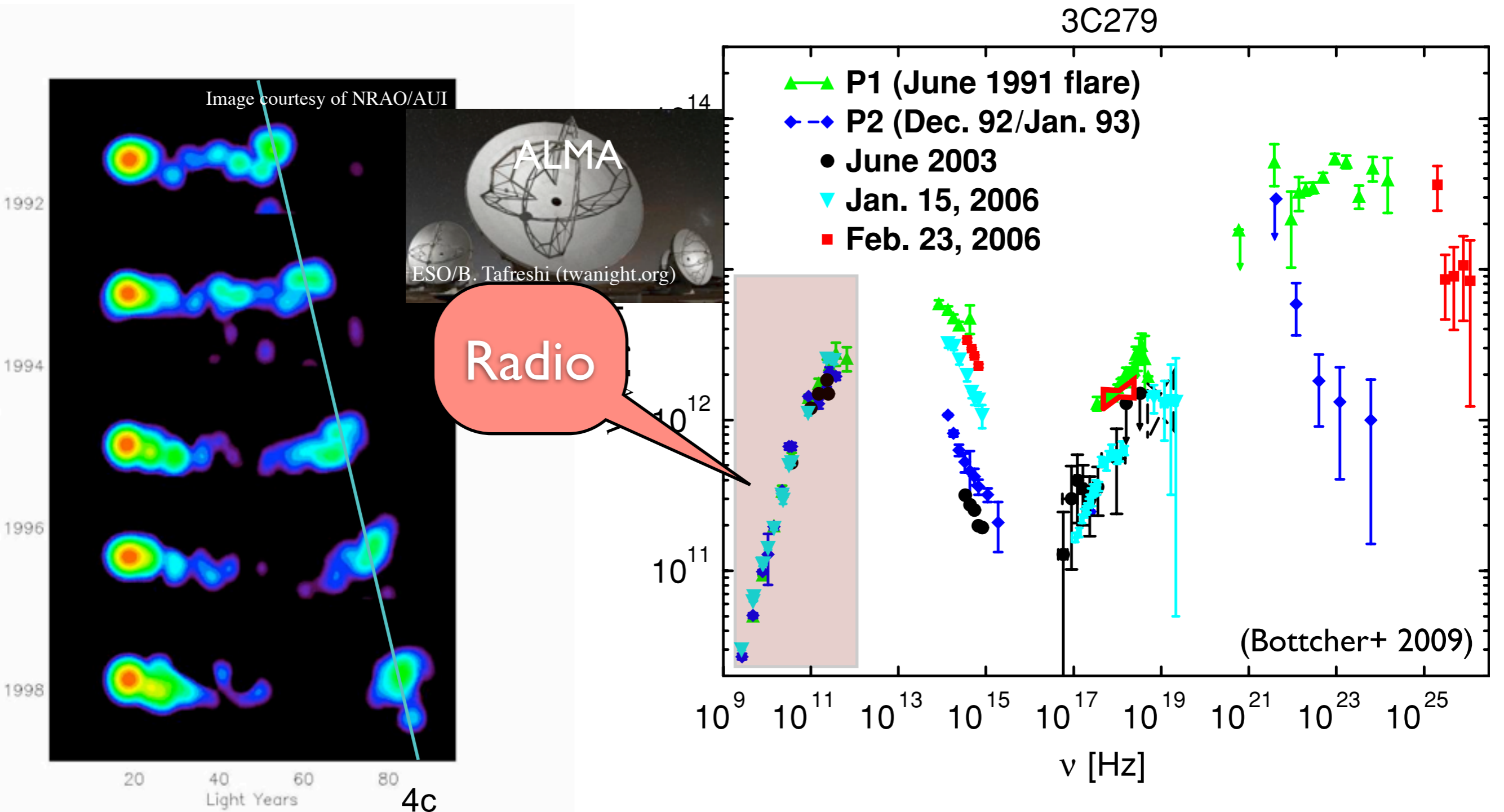
Jet Beaming Allows Observing Dim Sources

Active Galaxy with Jet Pointing at us: 3C279



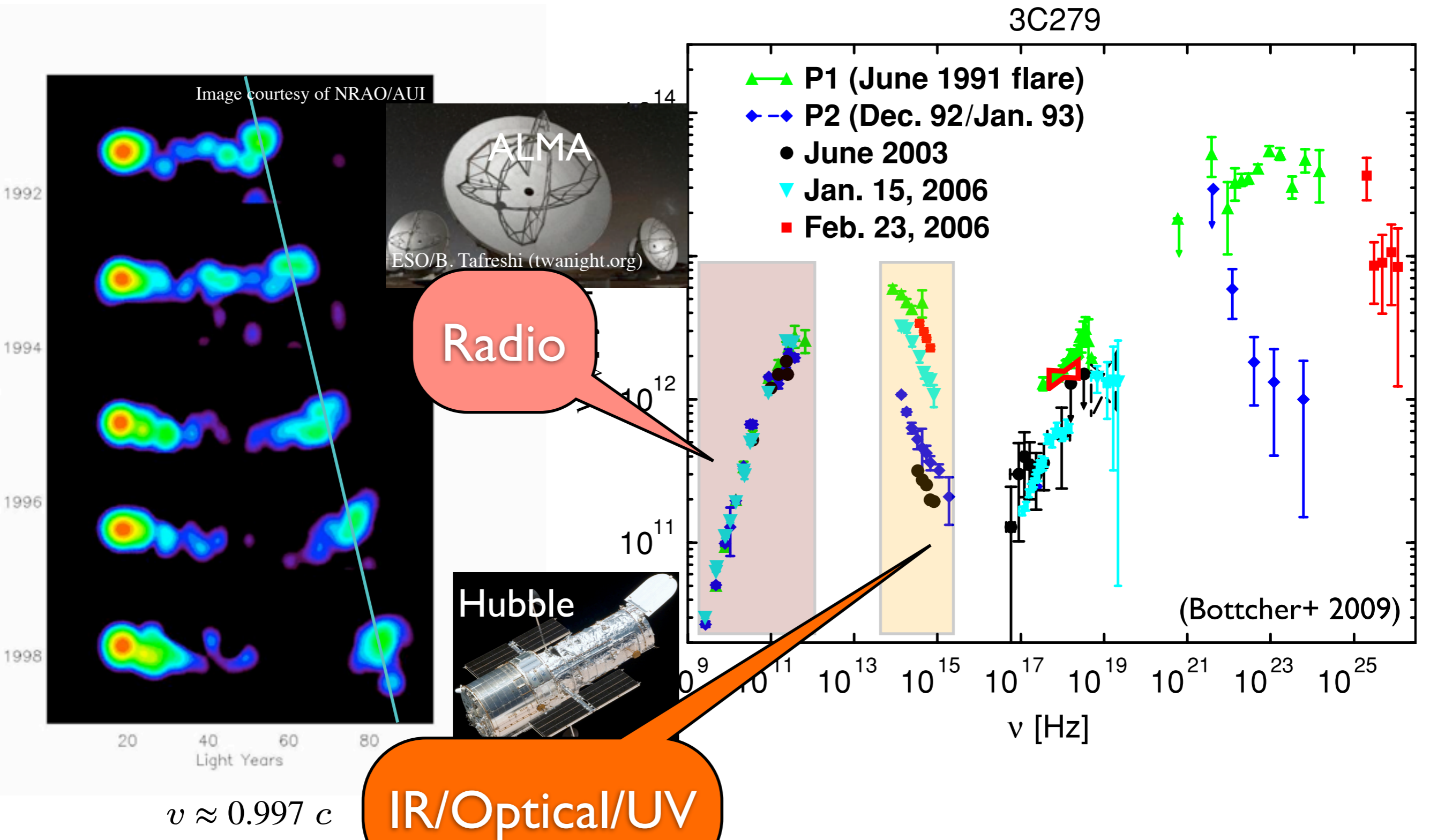
Jet Beaming Allows Observing Dim Sources

Active Galaxy with Jet Pointing at us: 3C279



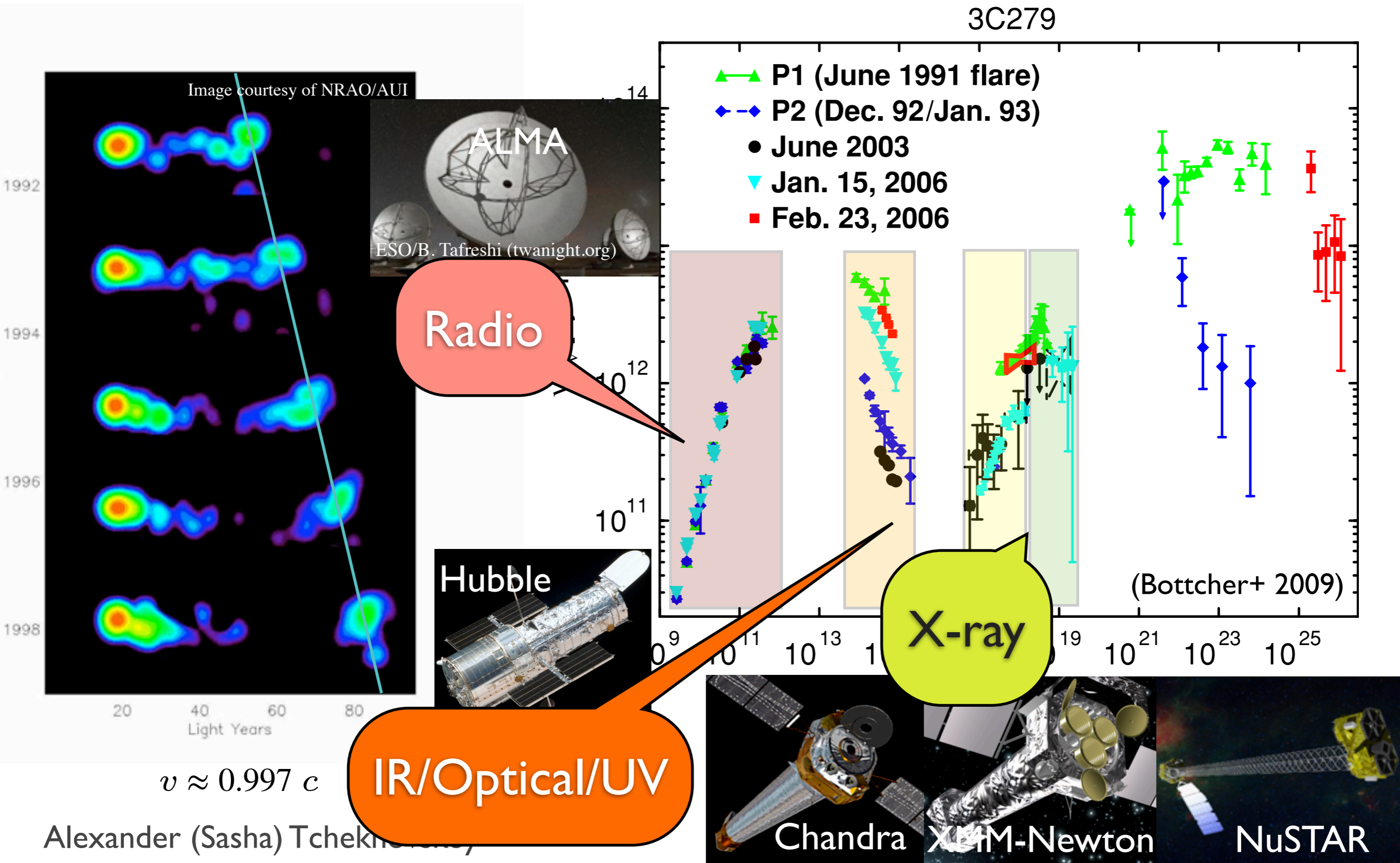
Jet Beaming Allows Observing Dim Sources

Active Galaxy with Jet Pointing at us: 3C279



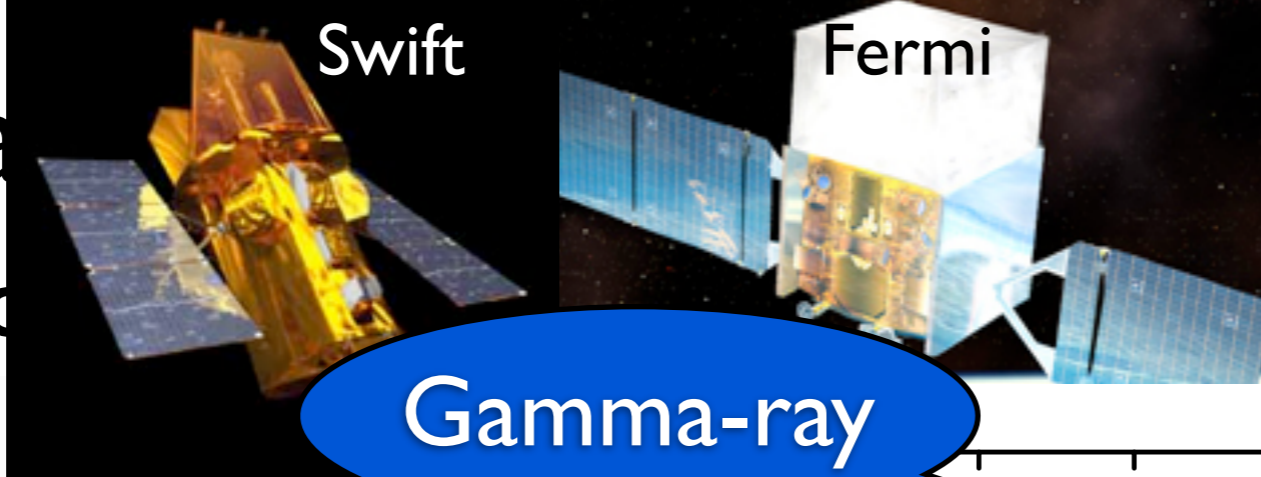
Jet Beaming Allows Observing Dim Sources

Active Galaxy with Jet Pointing at us: 3C279



Jet Beam

Ad

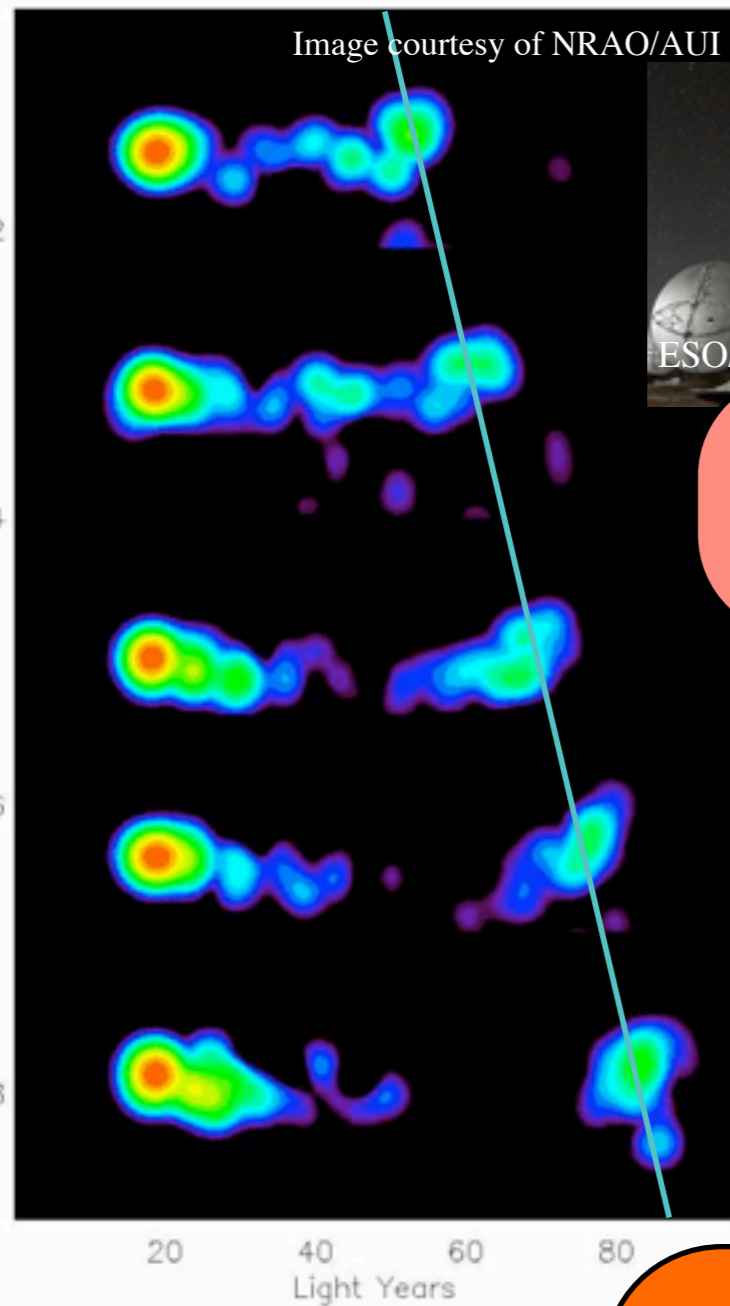


ing Dim Sources

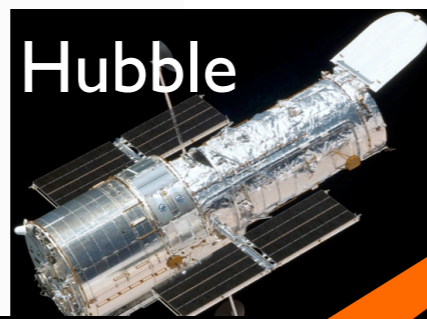
at us: 3C279

3C279

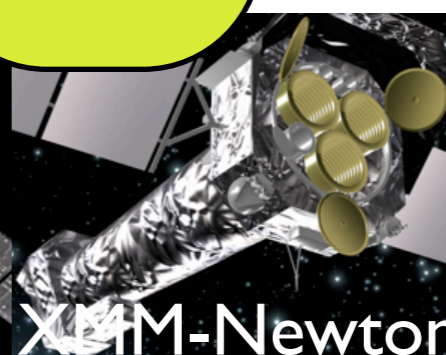
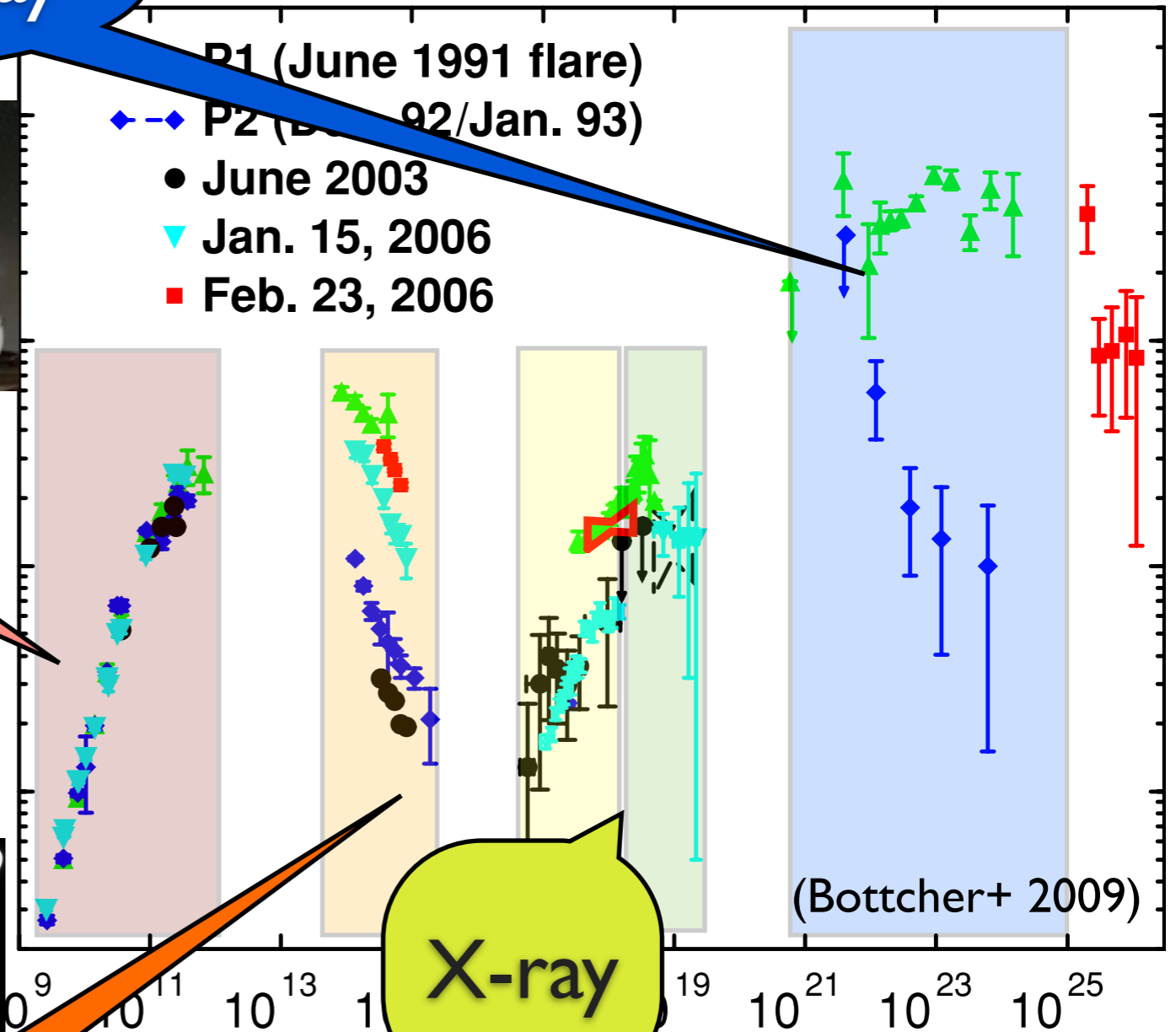
Gamma-ray



Radio



IR/Optical/UV

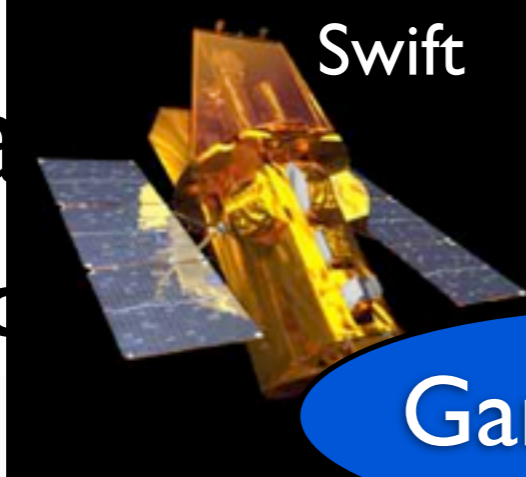


$v \approx 0.997 c$

Alexander (Sasha) Tchekhovskoy

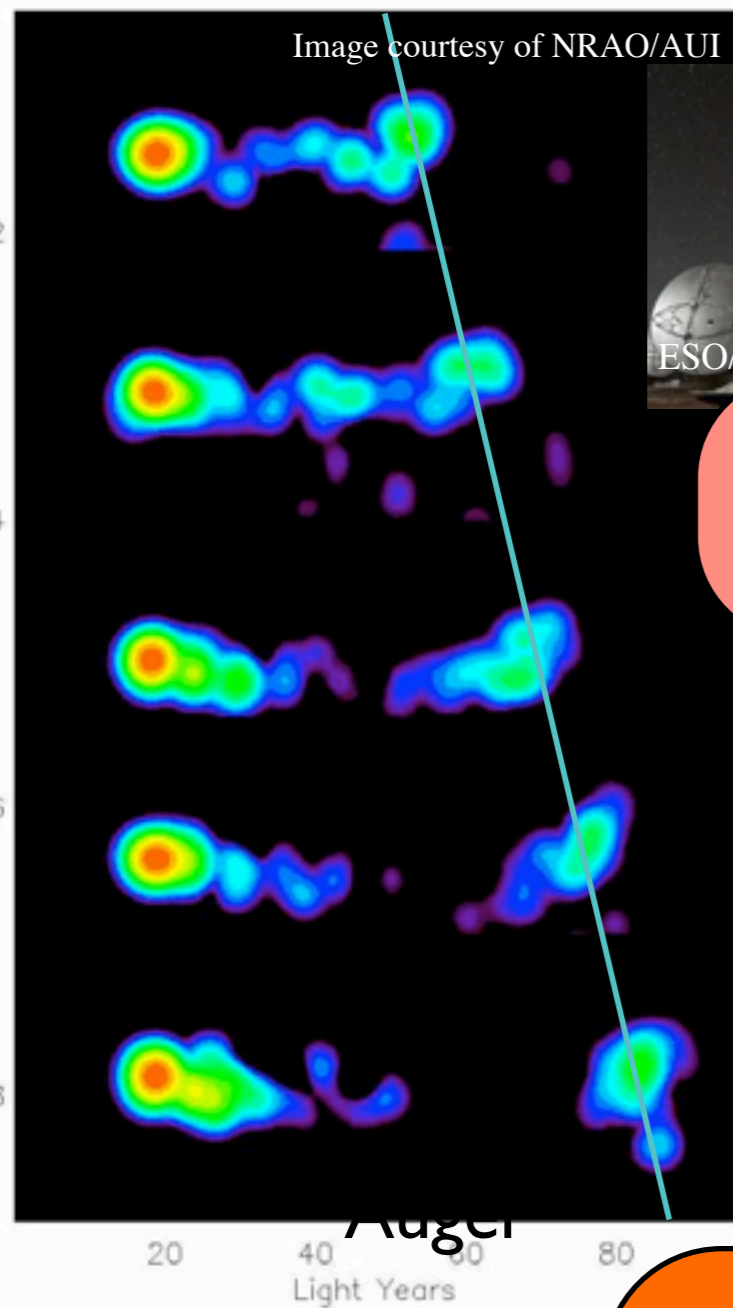
Jet Beams

Active

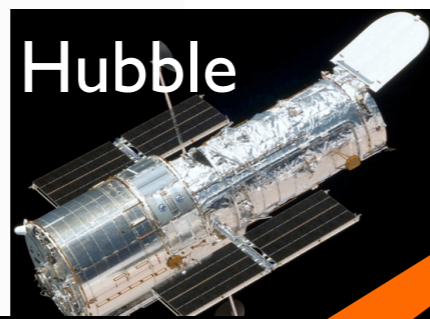


Gamma-ray

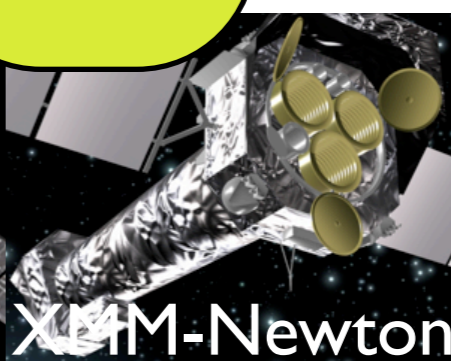
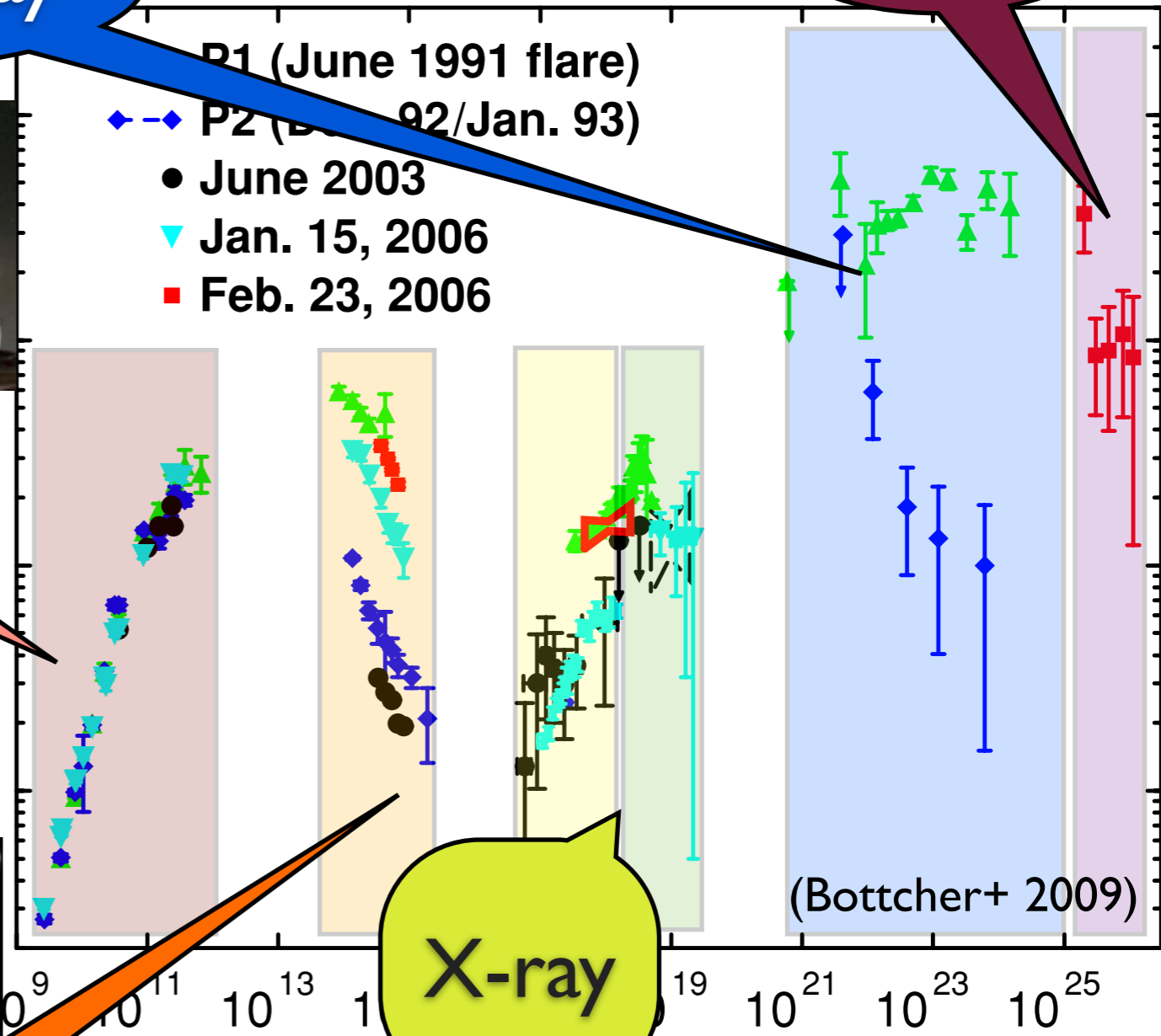
TeV



Radio



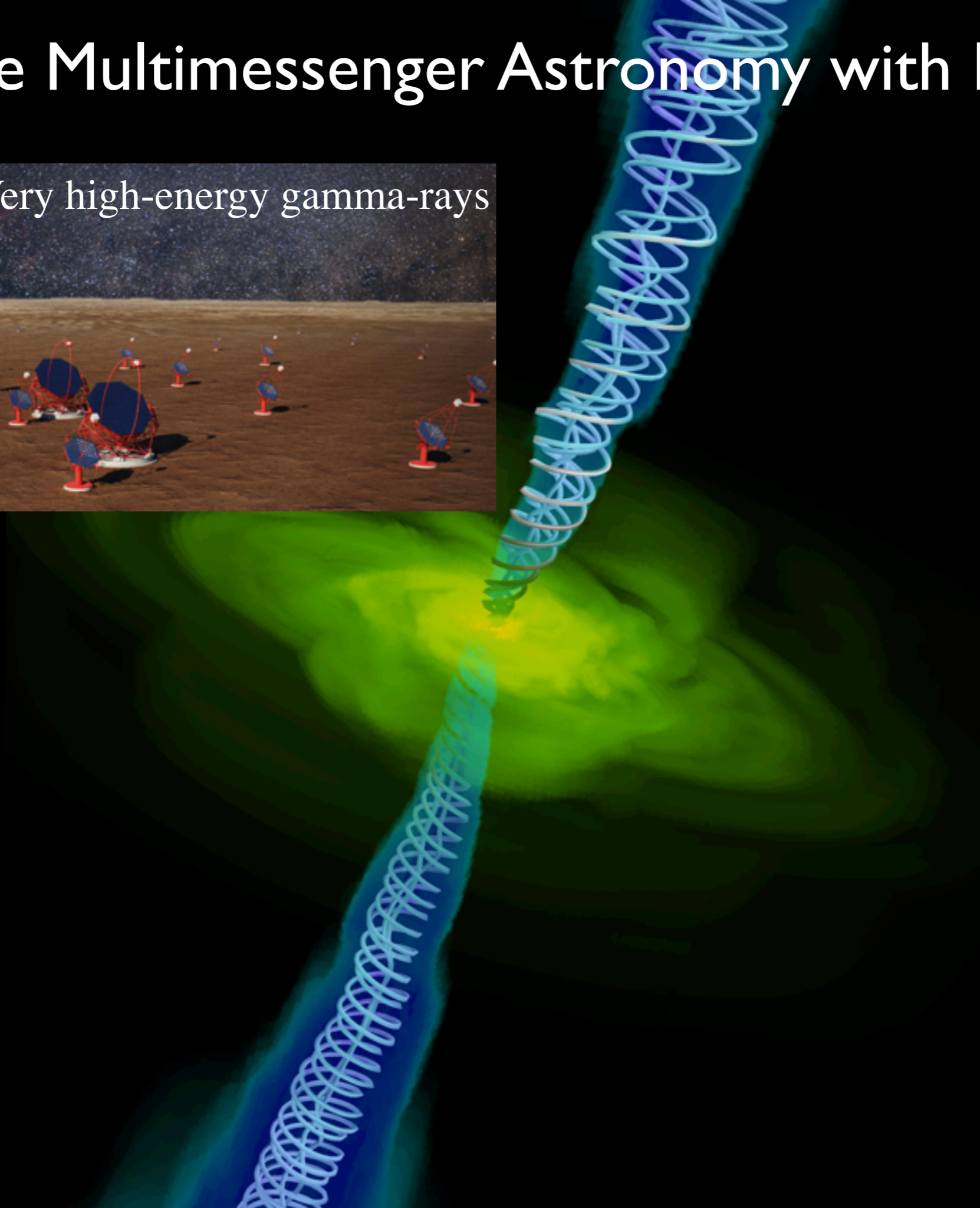
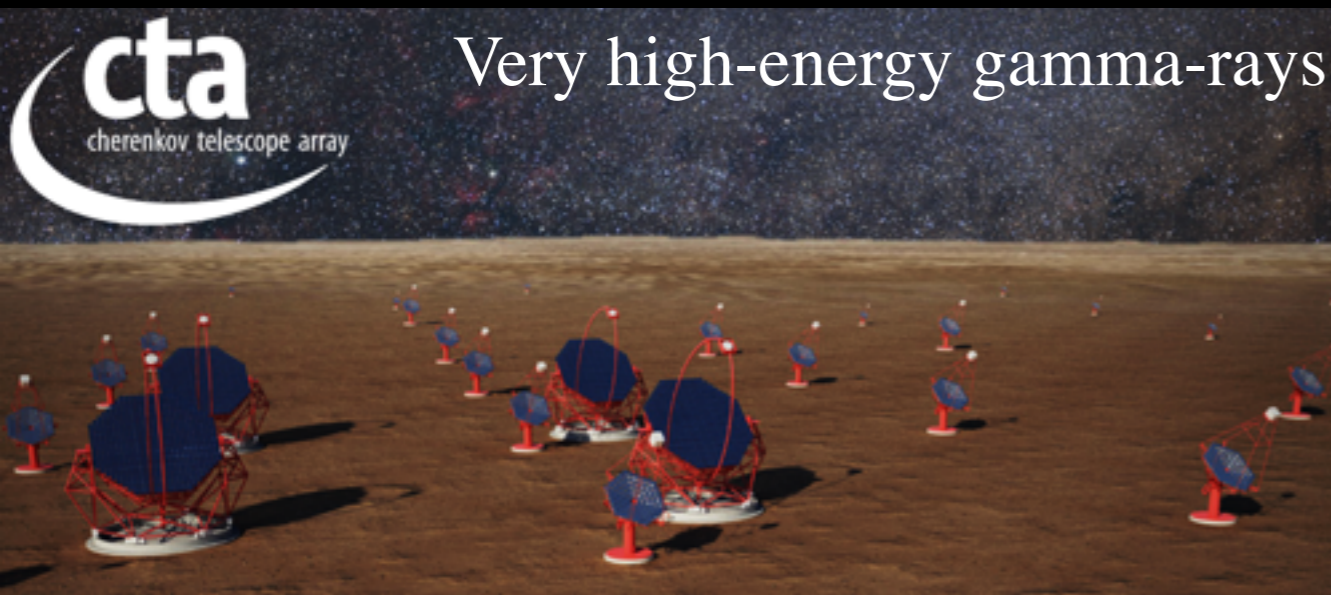
IR/Optical/UV



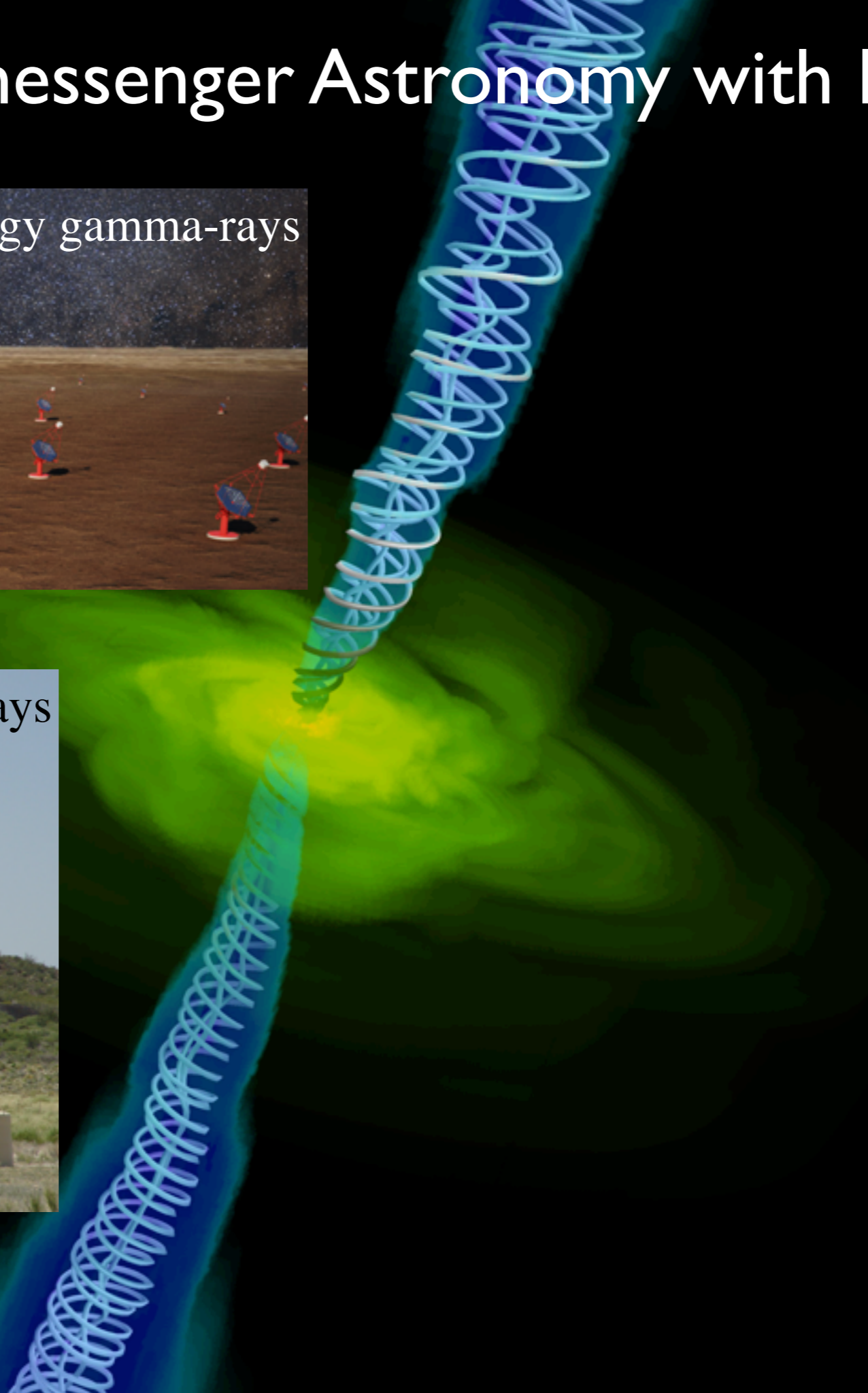
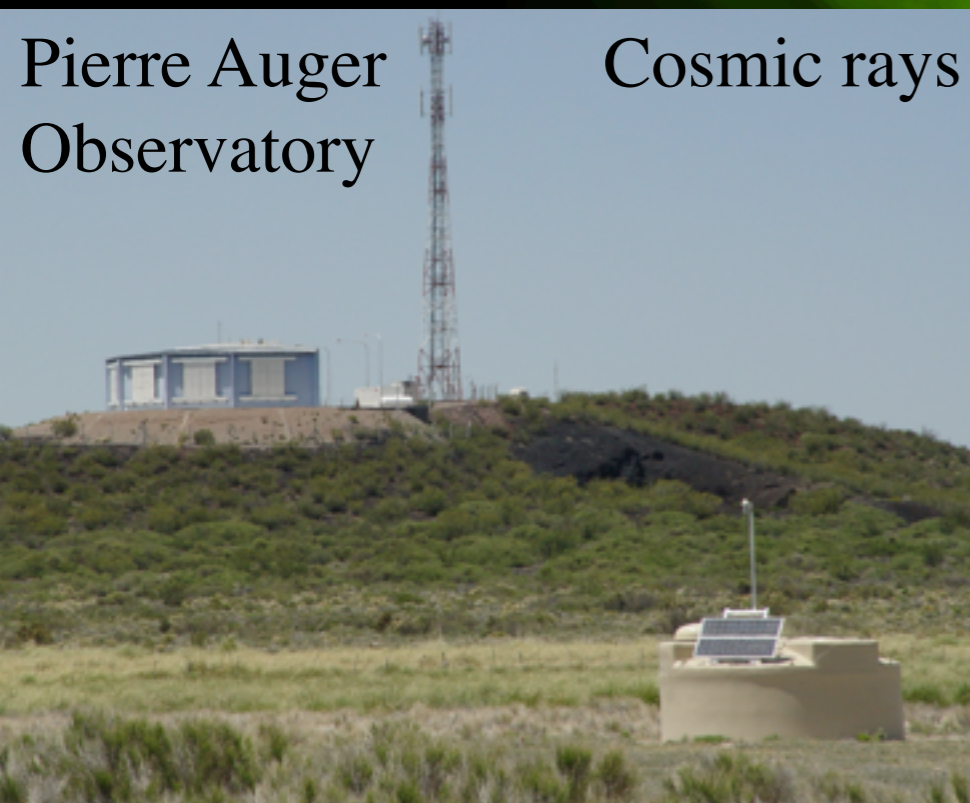
$$v \approx 0.997 c$$

Alexander (Sasha) Tchekhovskoy

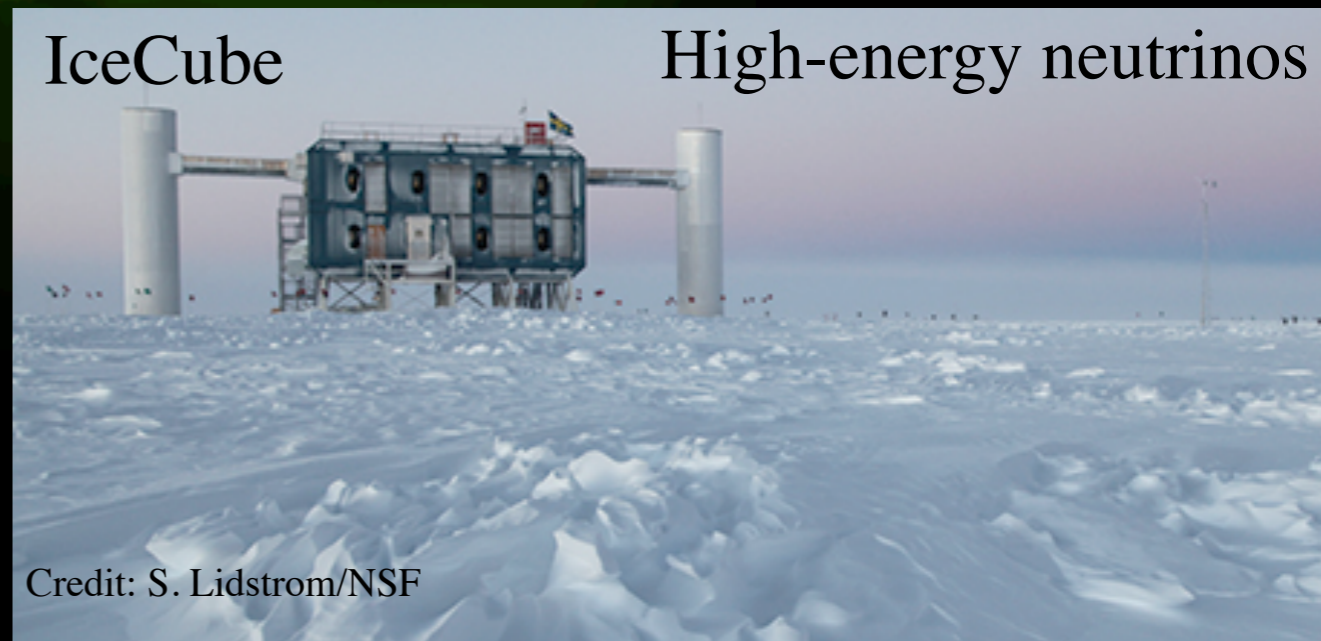
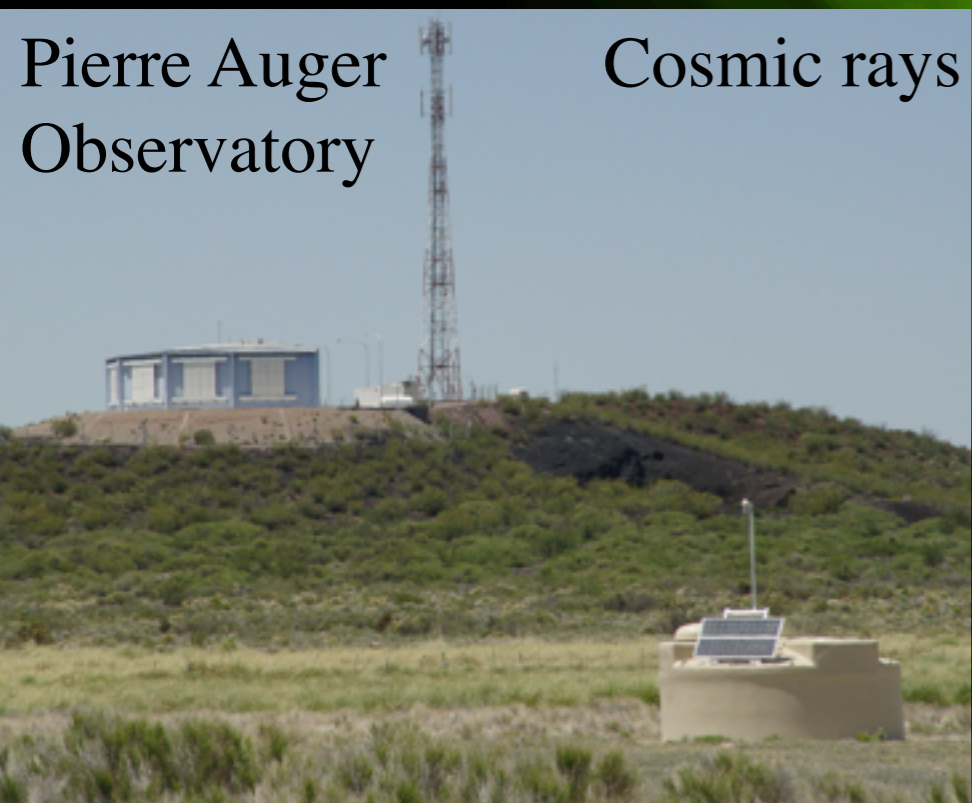
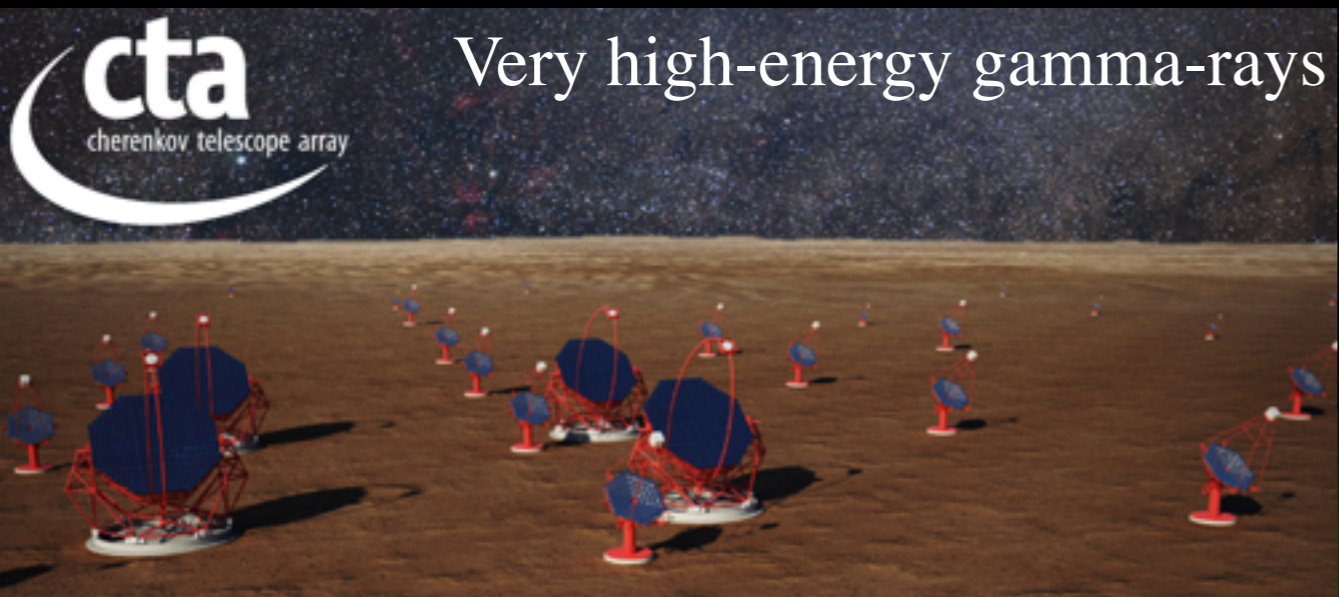
Jets Enable Multimessenger Astronomy with Black Holes



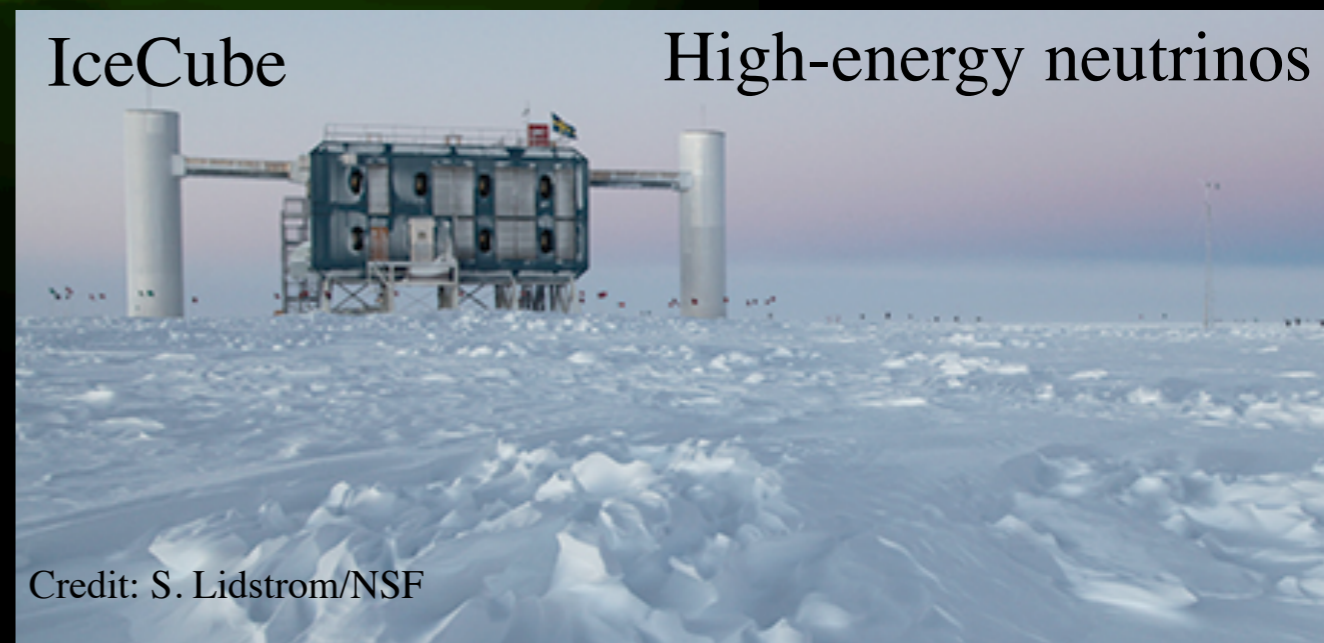
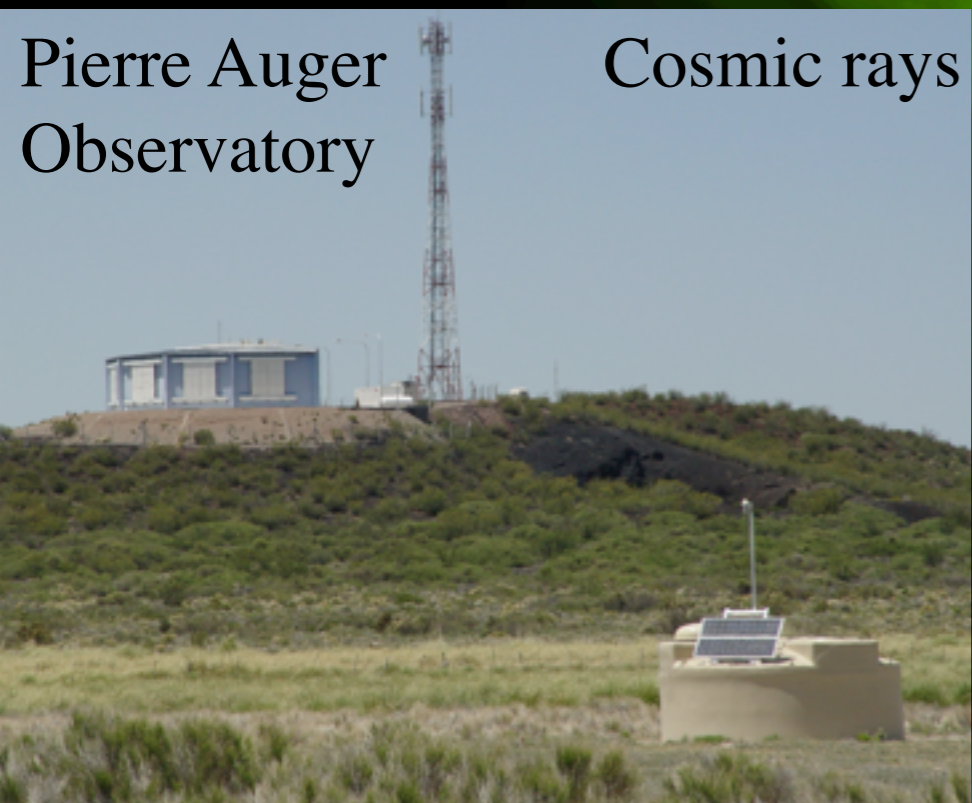
Jets Enable Multimessenger Astronomy with Black Holes



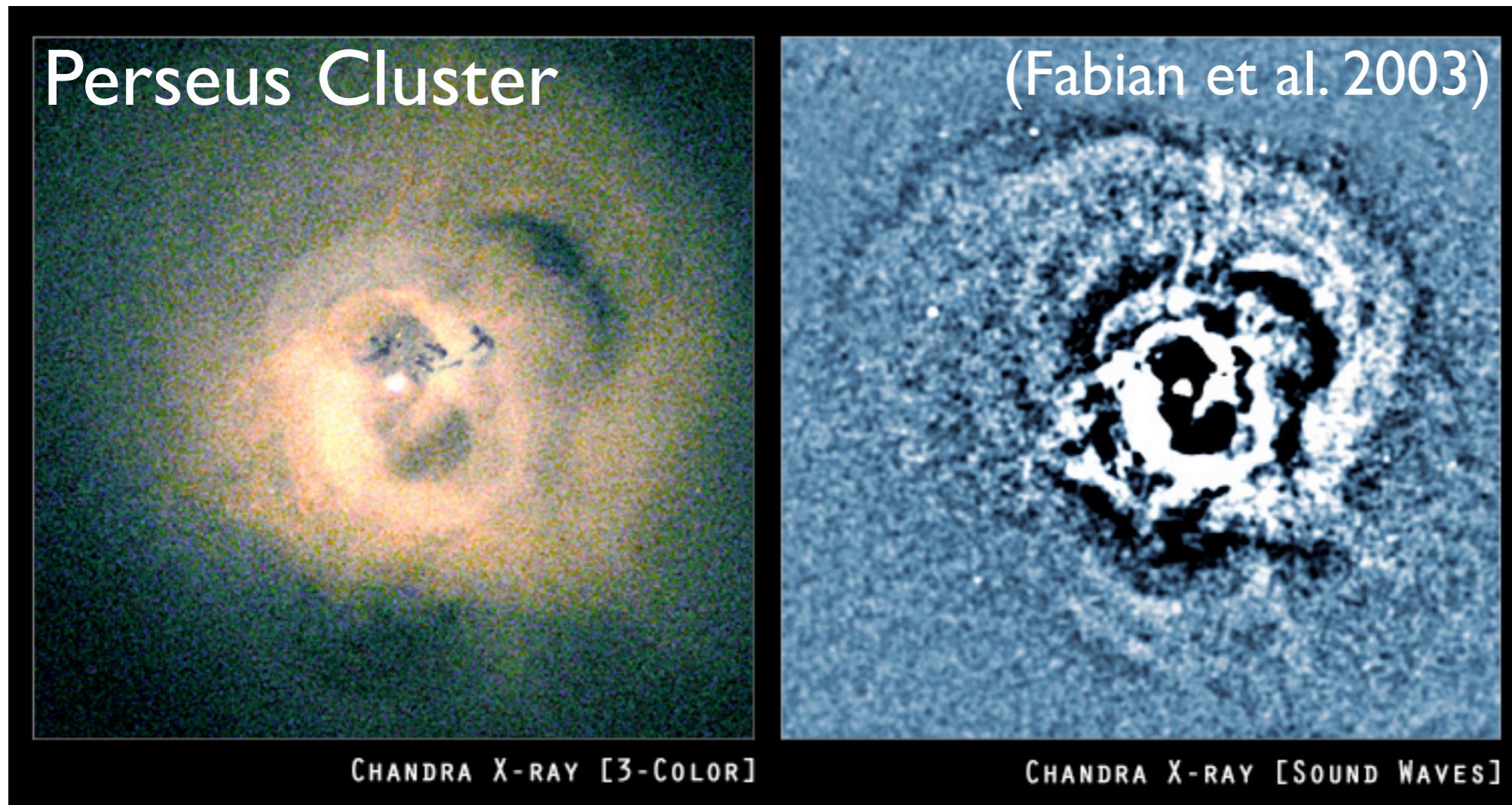
Jets Enable Multimessenger Astronomy with Black Holes



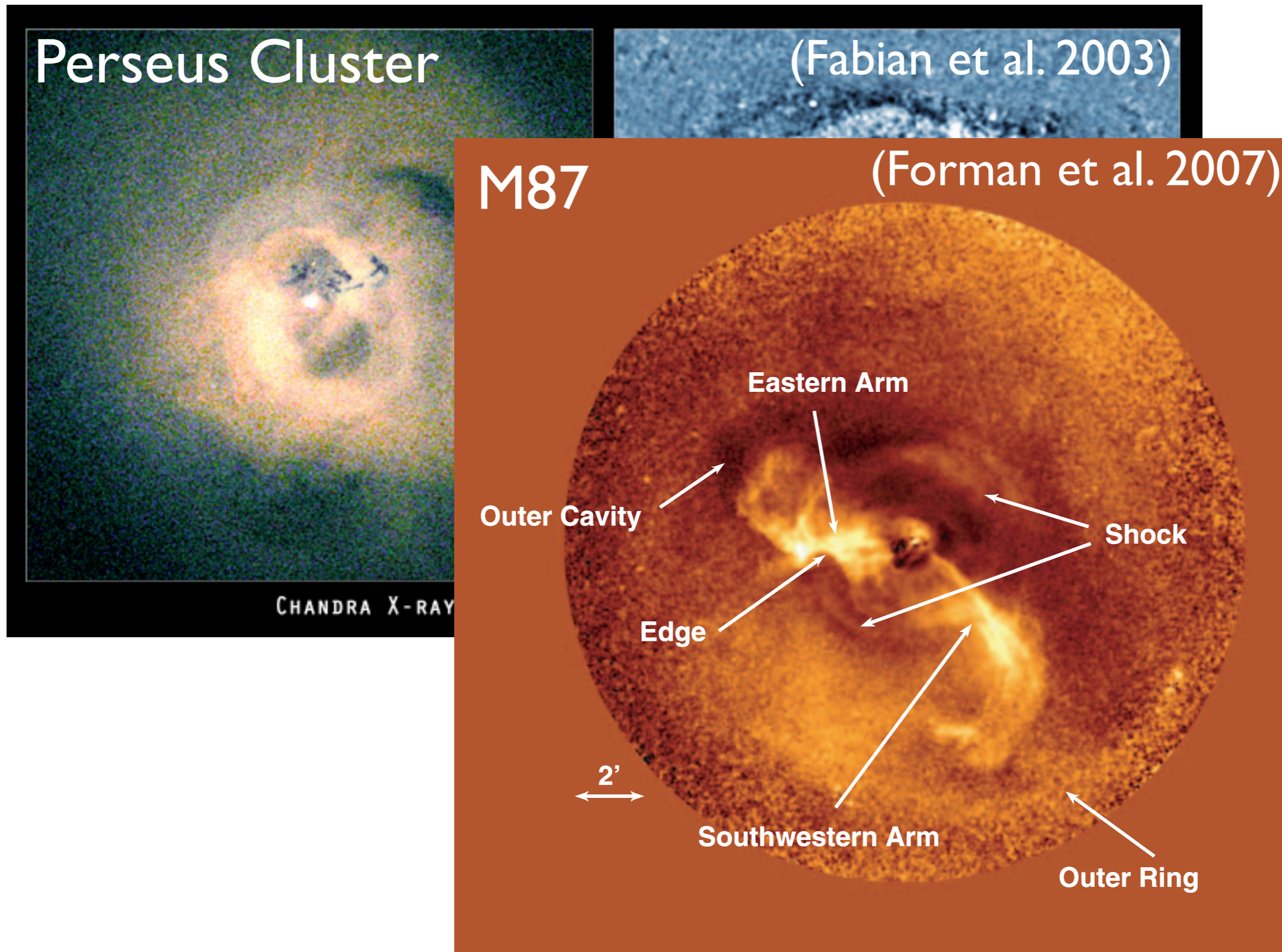
Jets Enable Multimessenger Astronomy with Black Holes



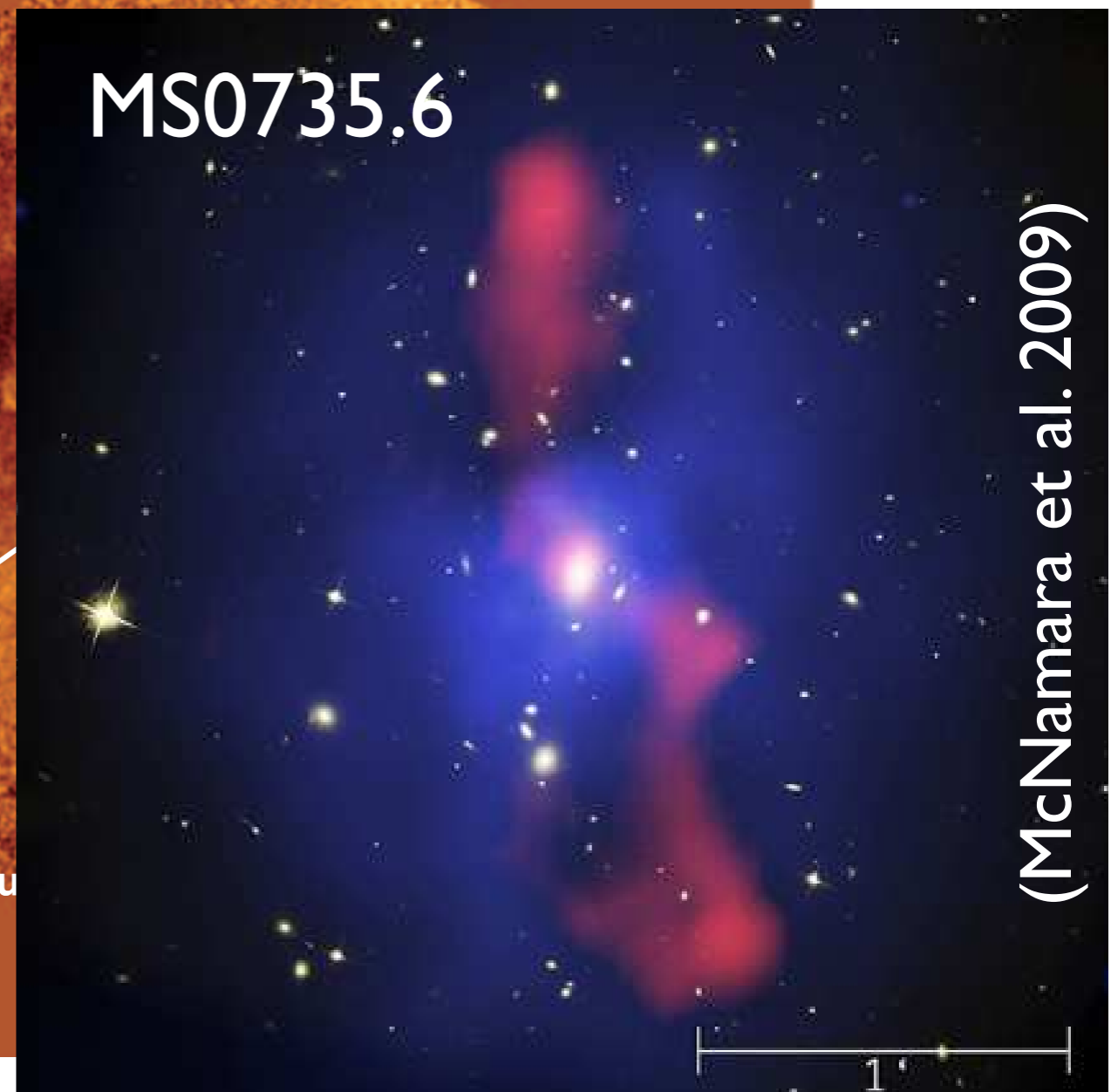
Jets Affect Galaxies/Clusters



Jets Affect Galaxies/Clusters



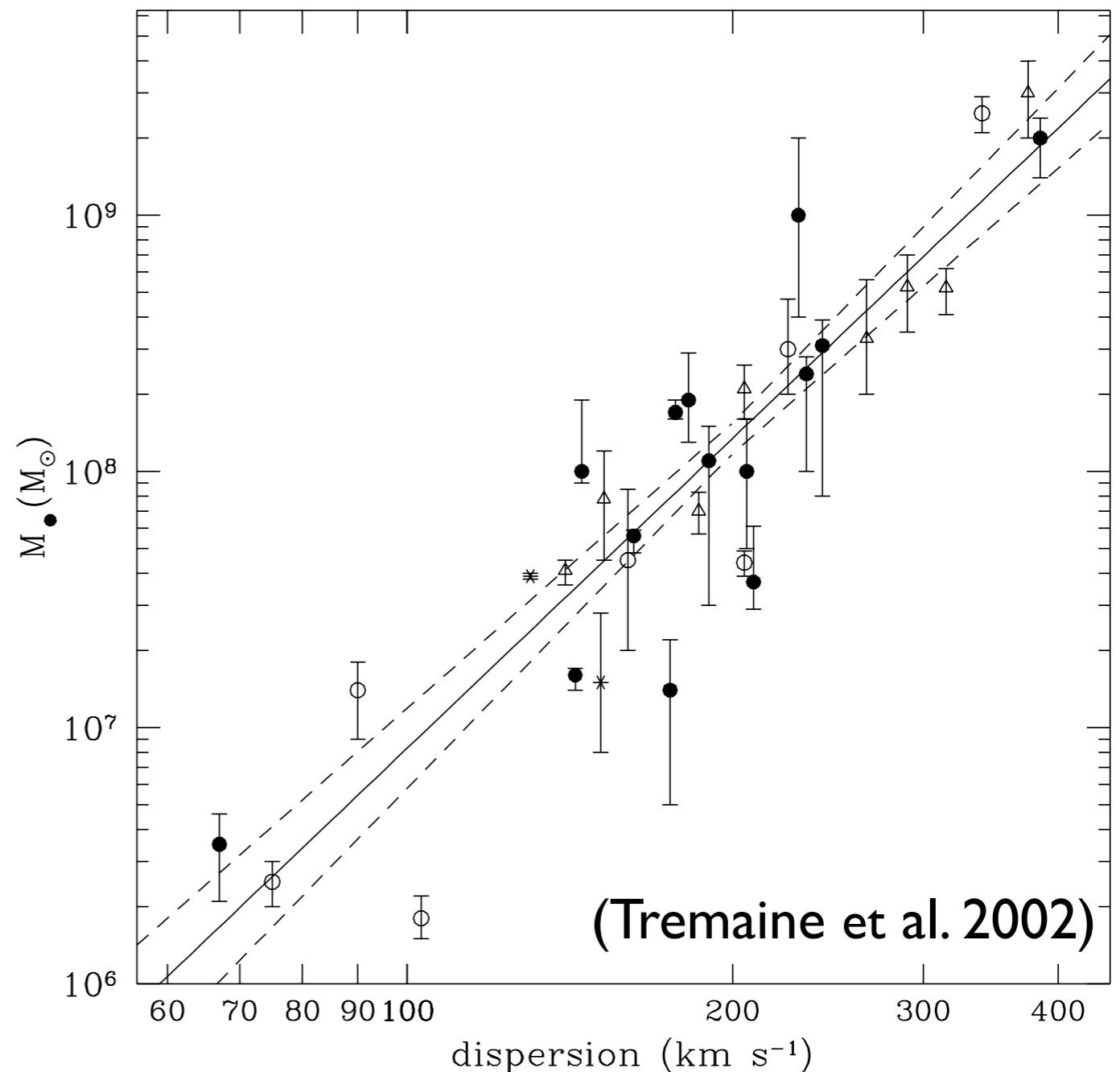
Jets Affect Galaxies/Clusters



Jets Affect Galaxies/Clusters

“M-sigma” relation: BH mass and stellar velocity dispersion are correlated

- Growth of the central BHs and their host galaxies are inter-connected
- Jet feedback?
- Radiative feedback?

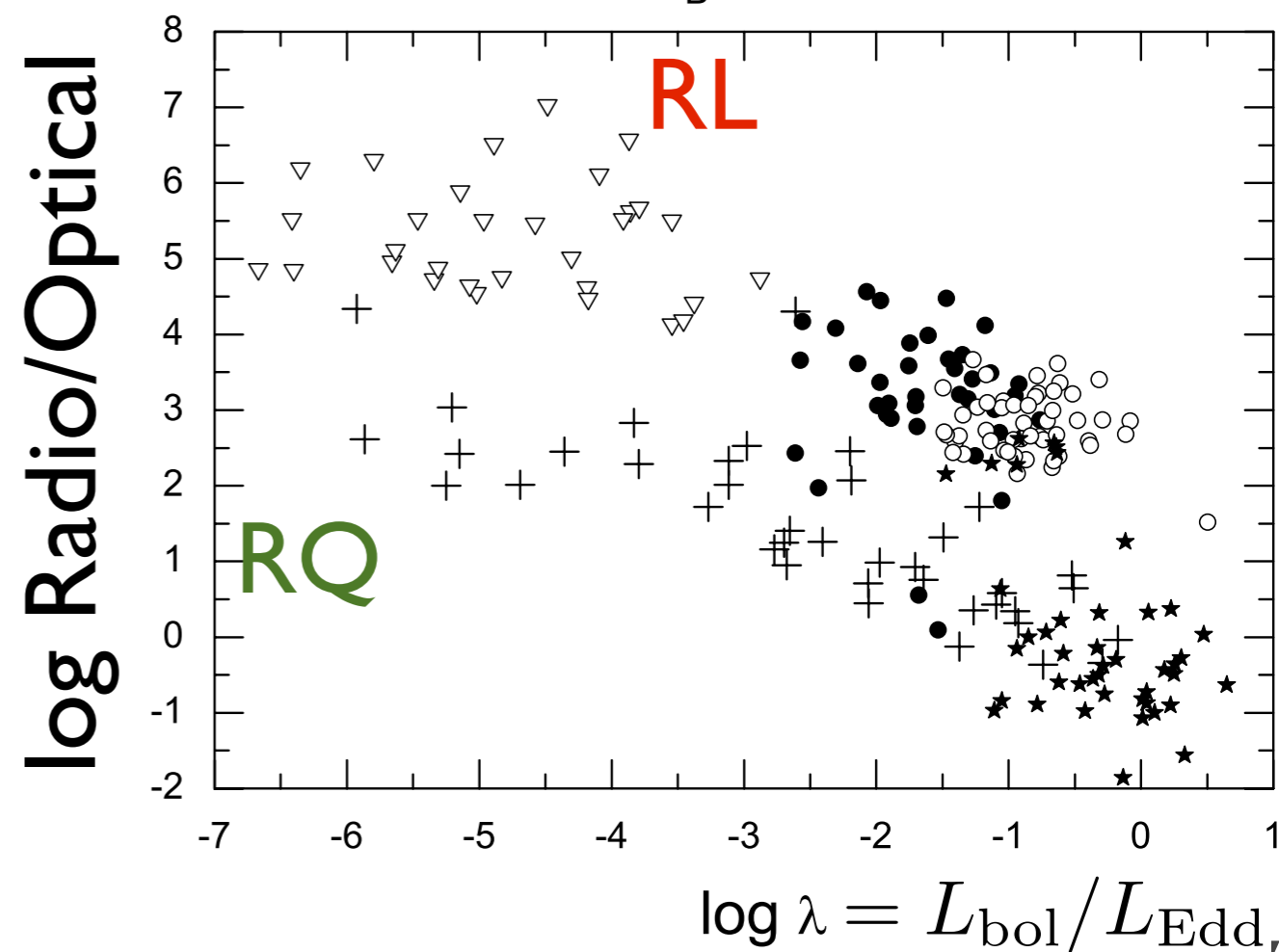
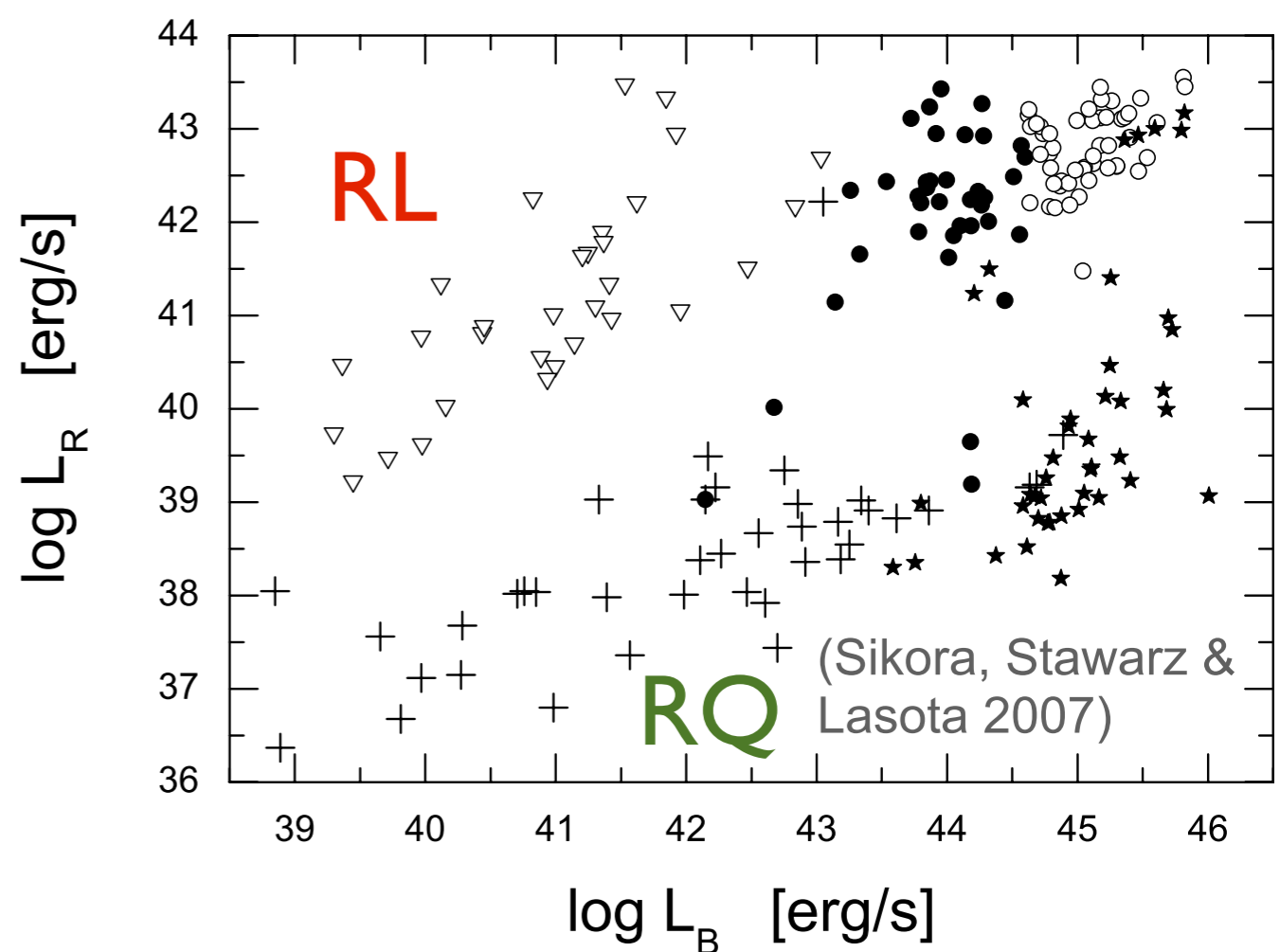


AGN Radio Loud/Quiet Dichotomy

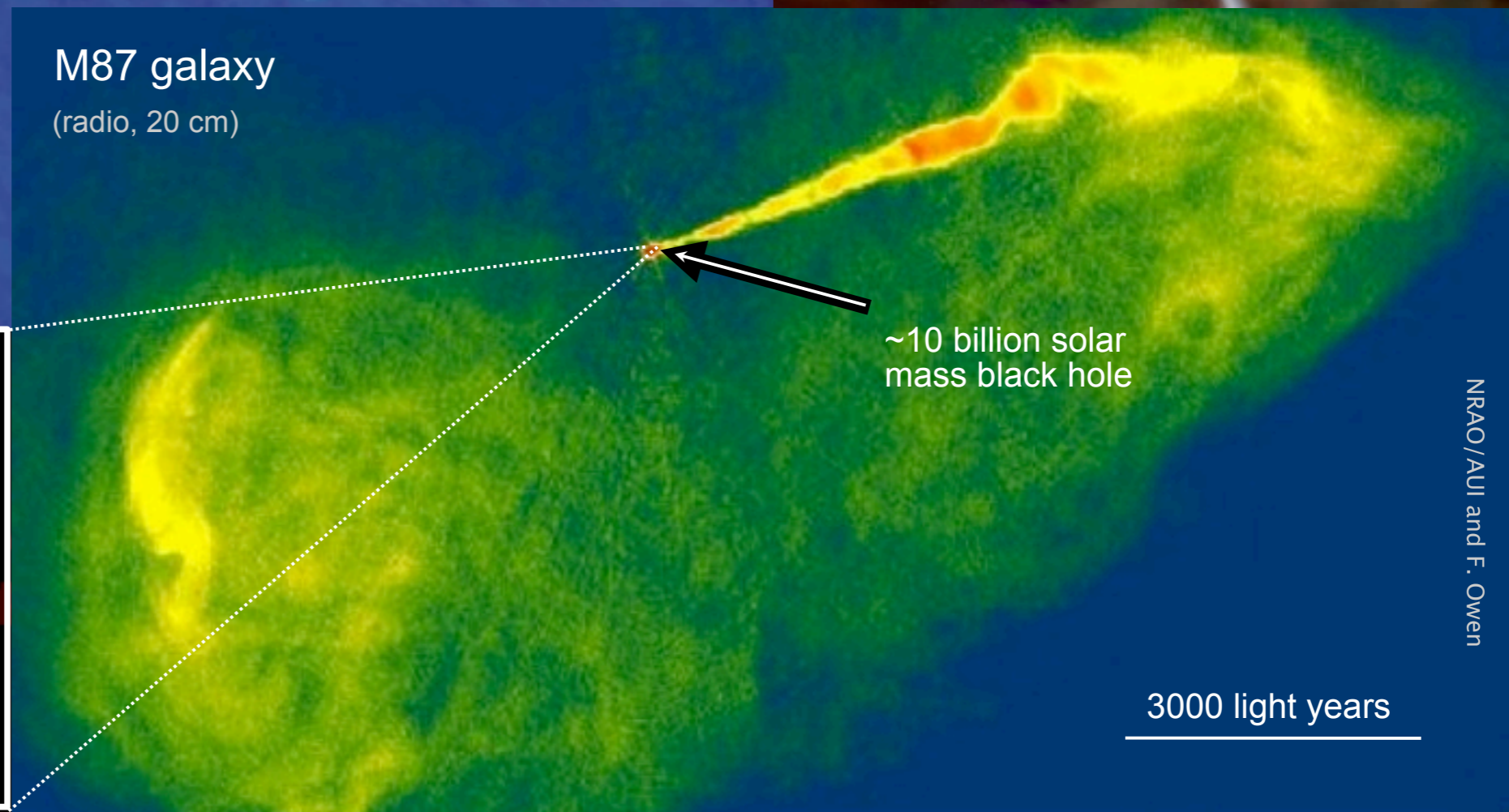
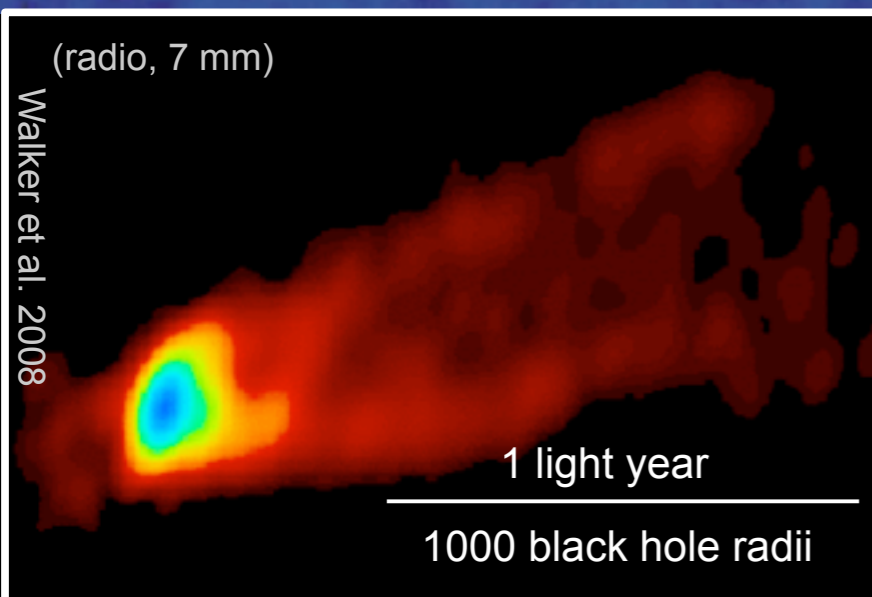
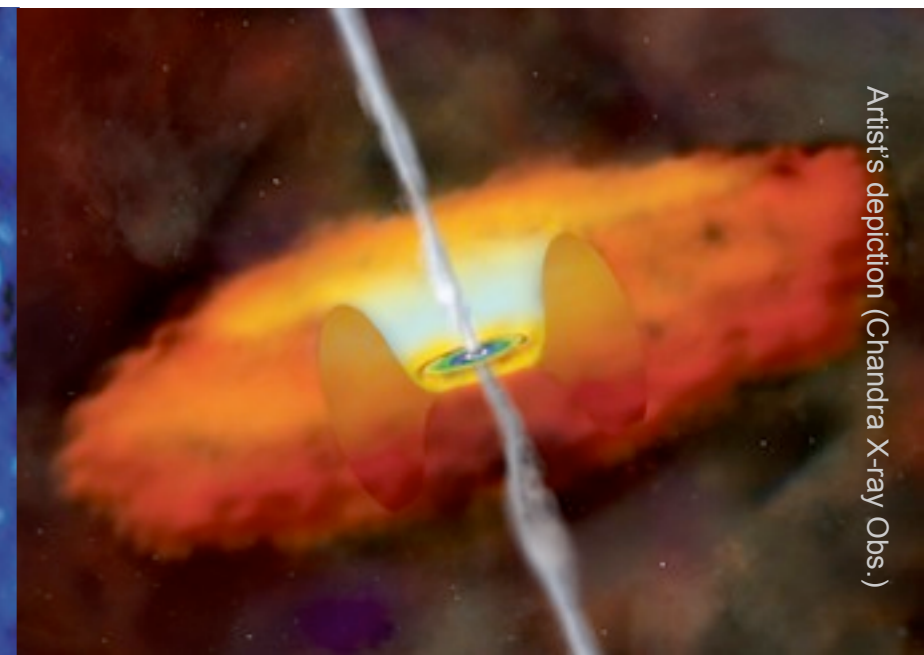
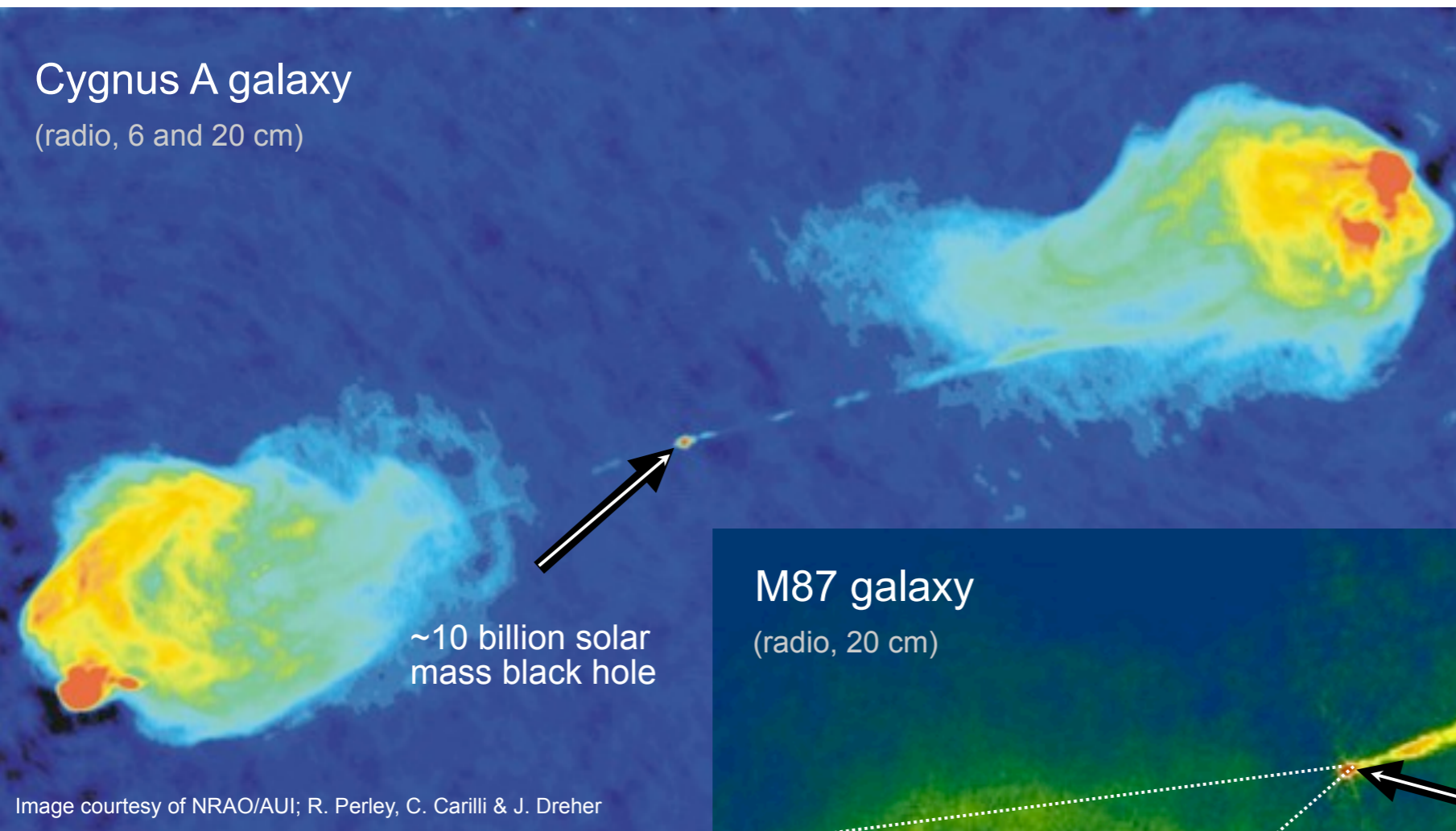
- Factor of 1000 difference in radio luminosity.
- There must be at least one other parameter in addition to M and \dot{M} :

$$P_{\text{jet}}(M, \dot{M}; ??)$$

- Magnetic flux?
- Ambient medium? (Broderick & Fender 2012)
- BH spin? (Blandford 1990, Tchekhovskoy et al. 2010)

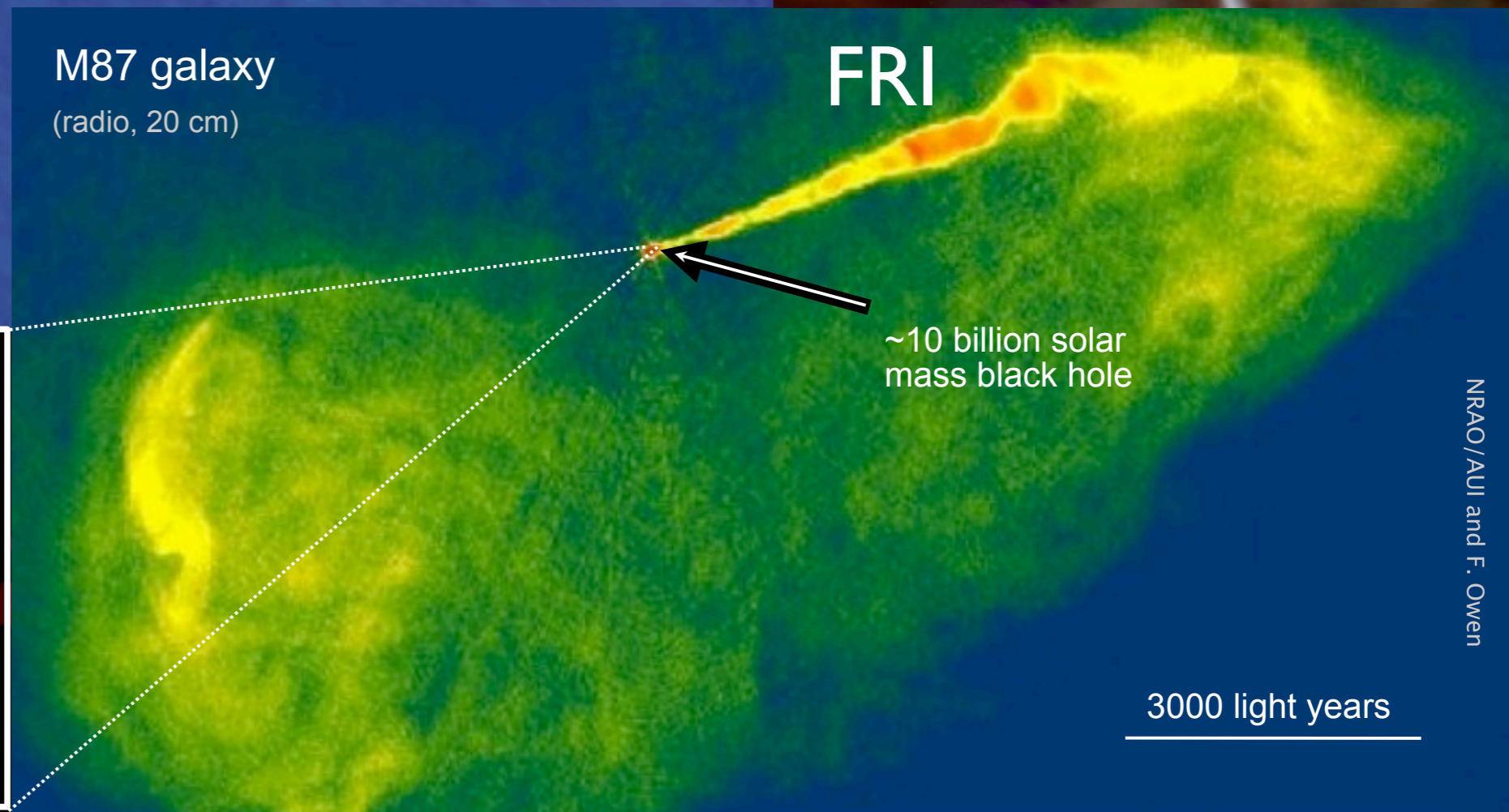
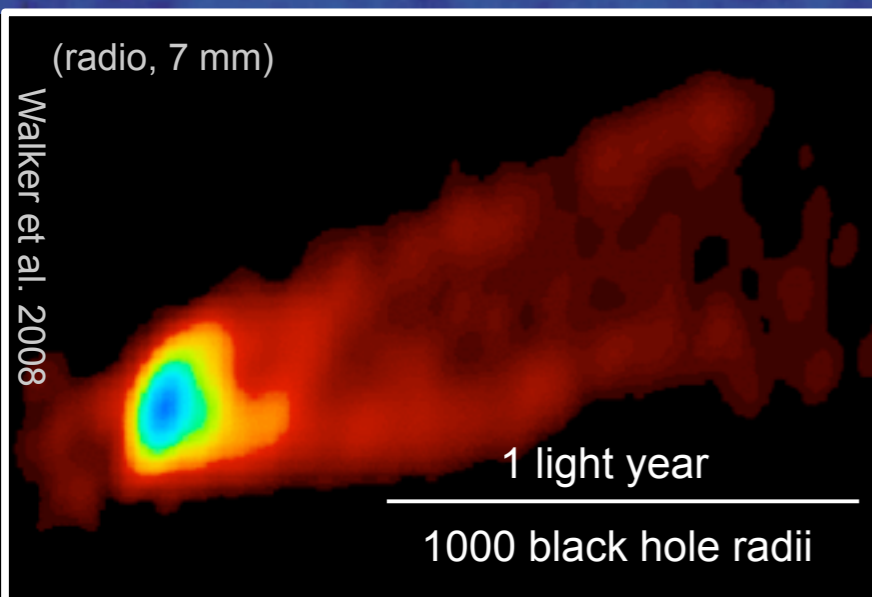
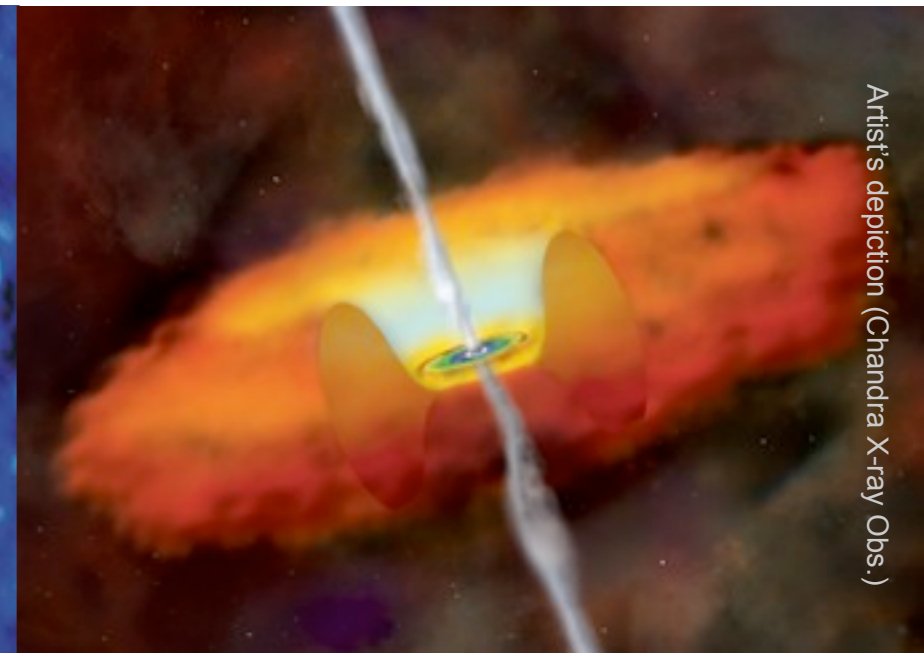
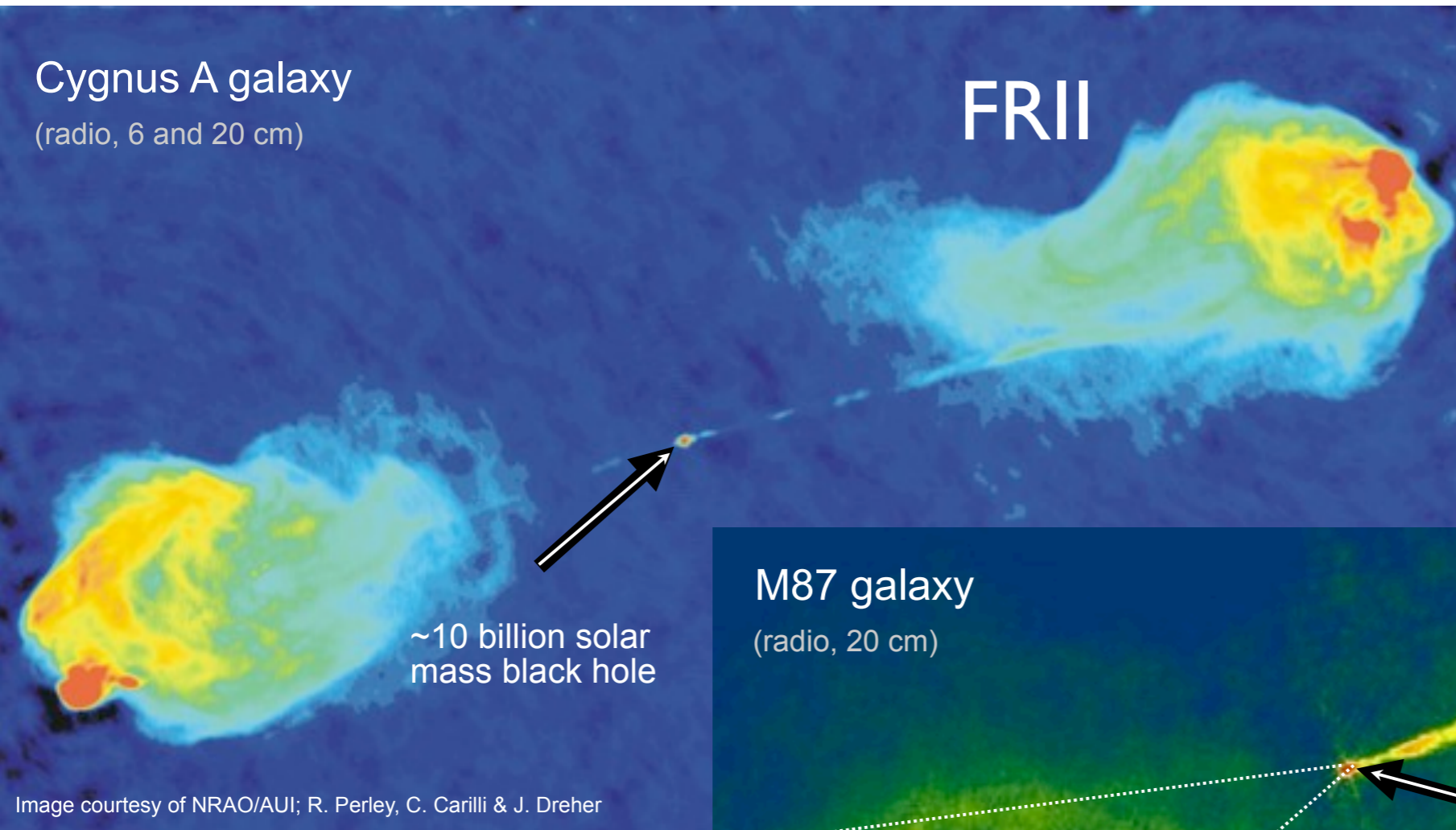


Jets: Beautiful and Challenging



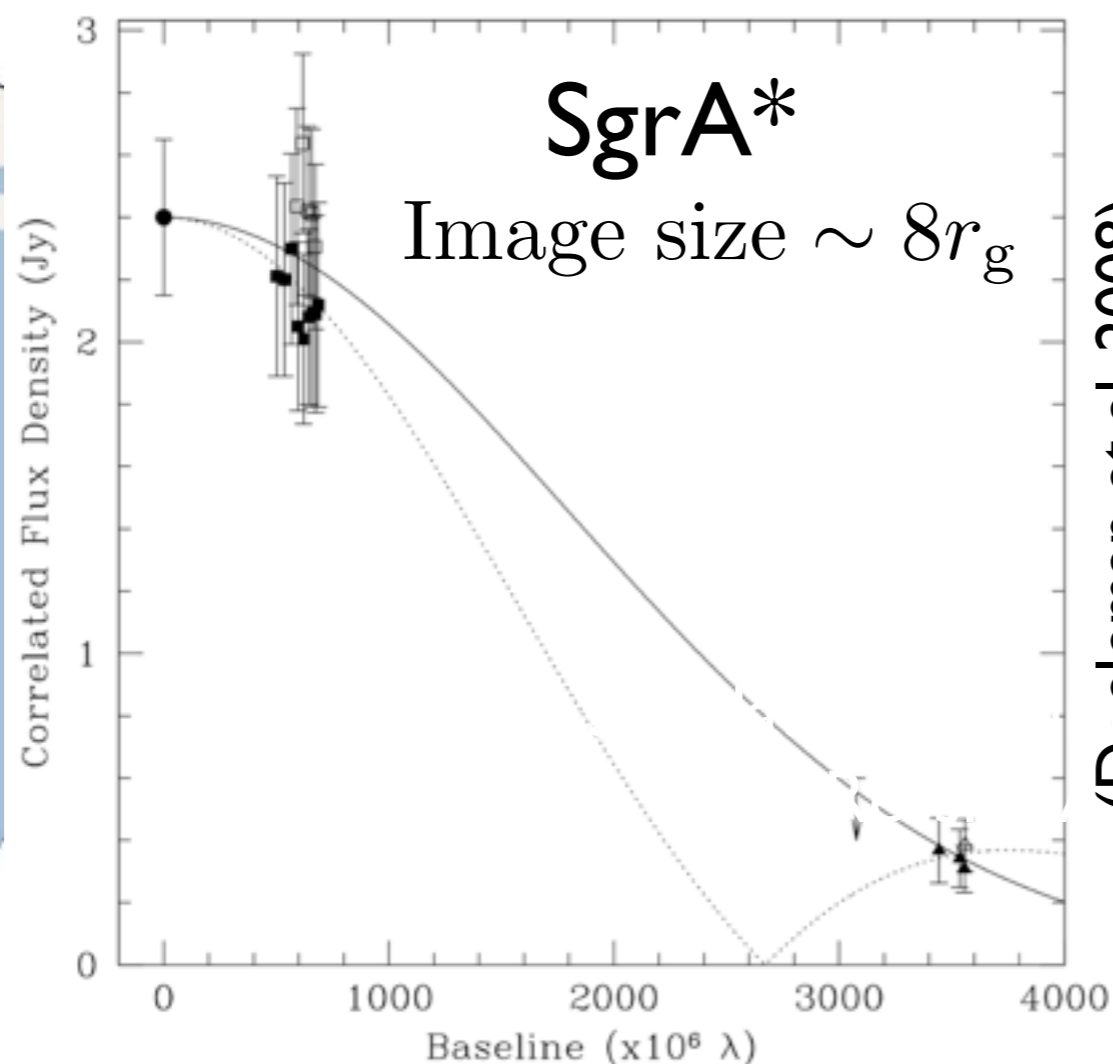
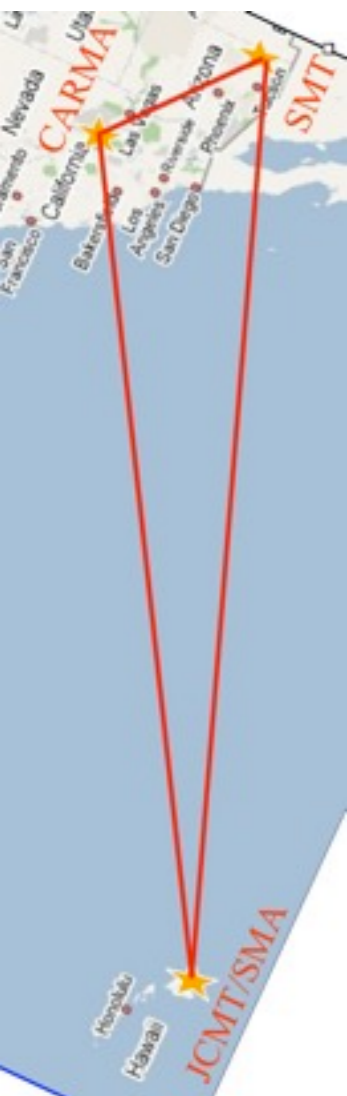
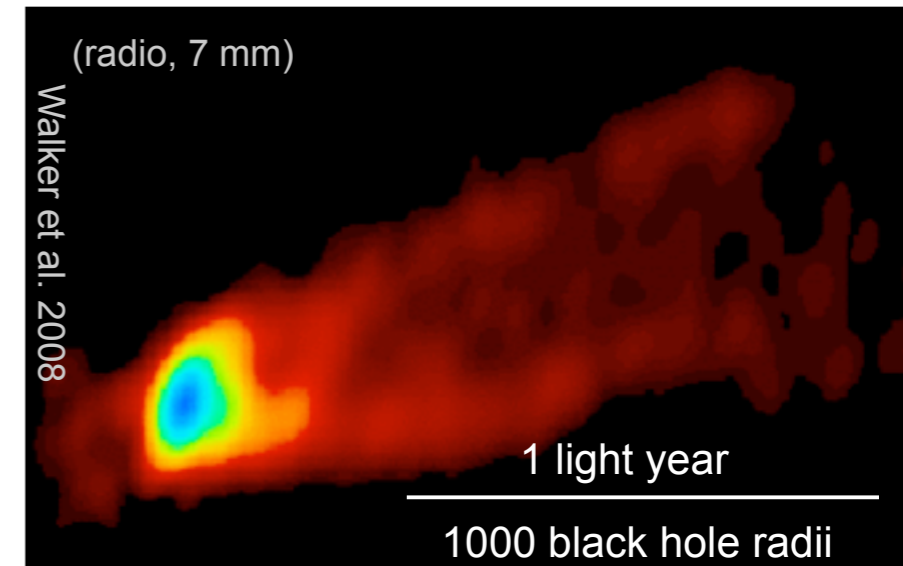
Jets: Beautiful and Challenging

FRI/FR II dichotomy (Fanaroff & Riley, 1974)

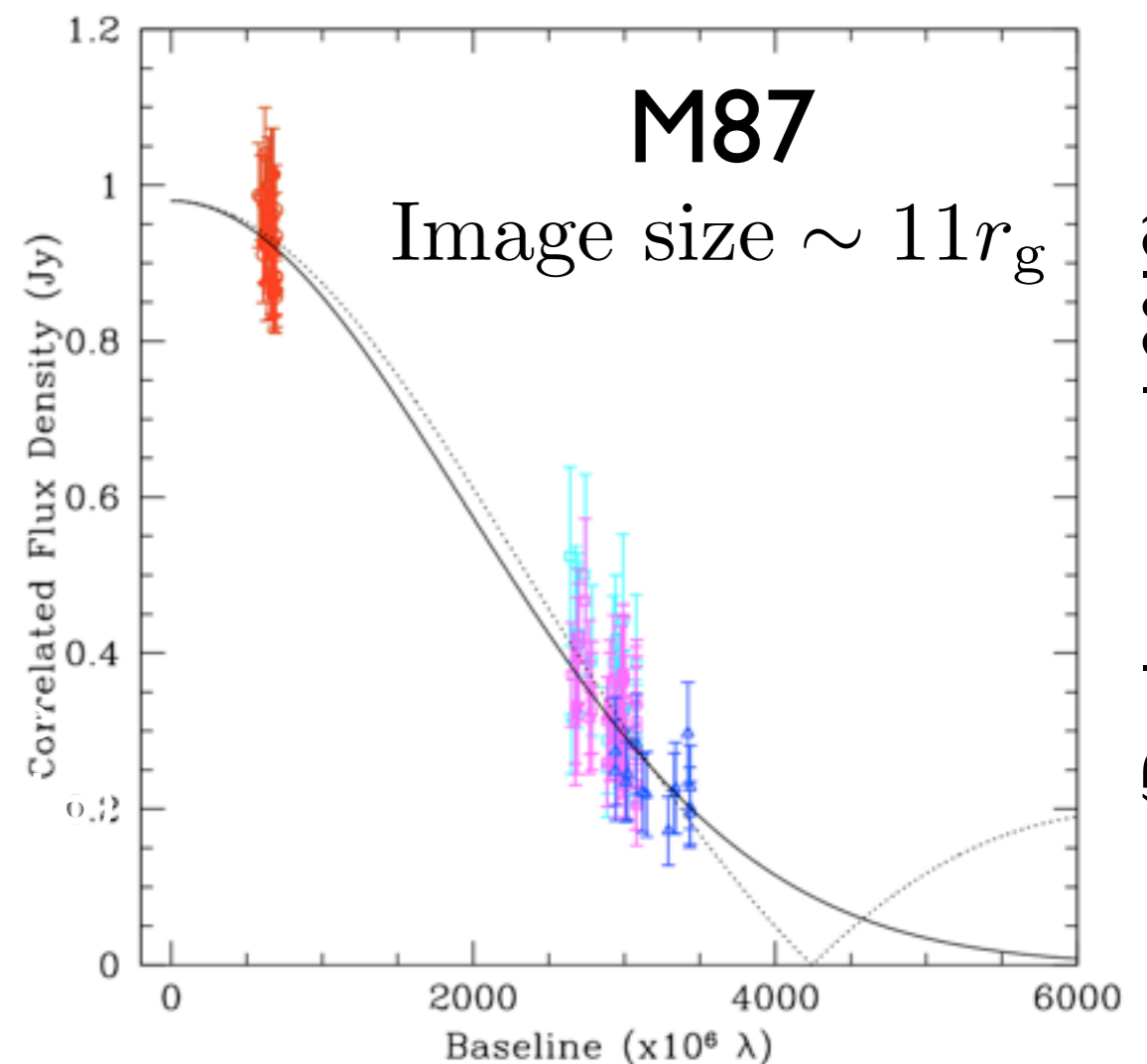


Event Horizon Telescope (EHT): VLBI images of Black Holes

- Two largest black holes on the sky
- Data is interpretation limited!

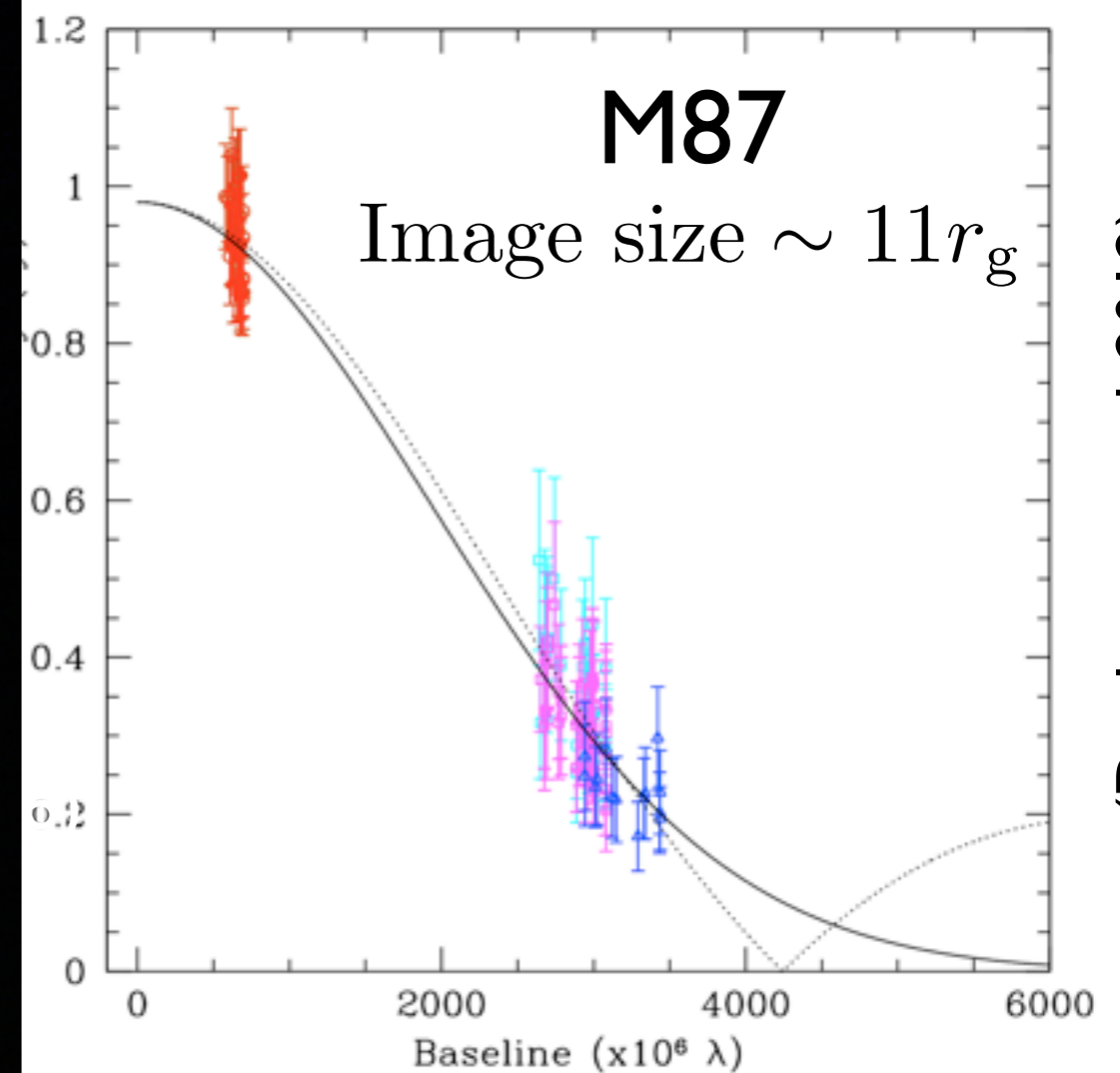
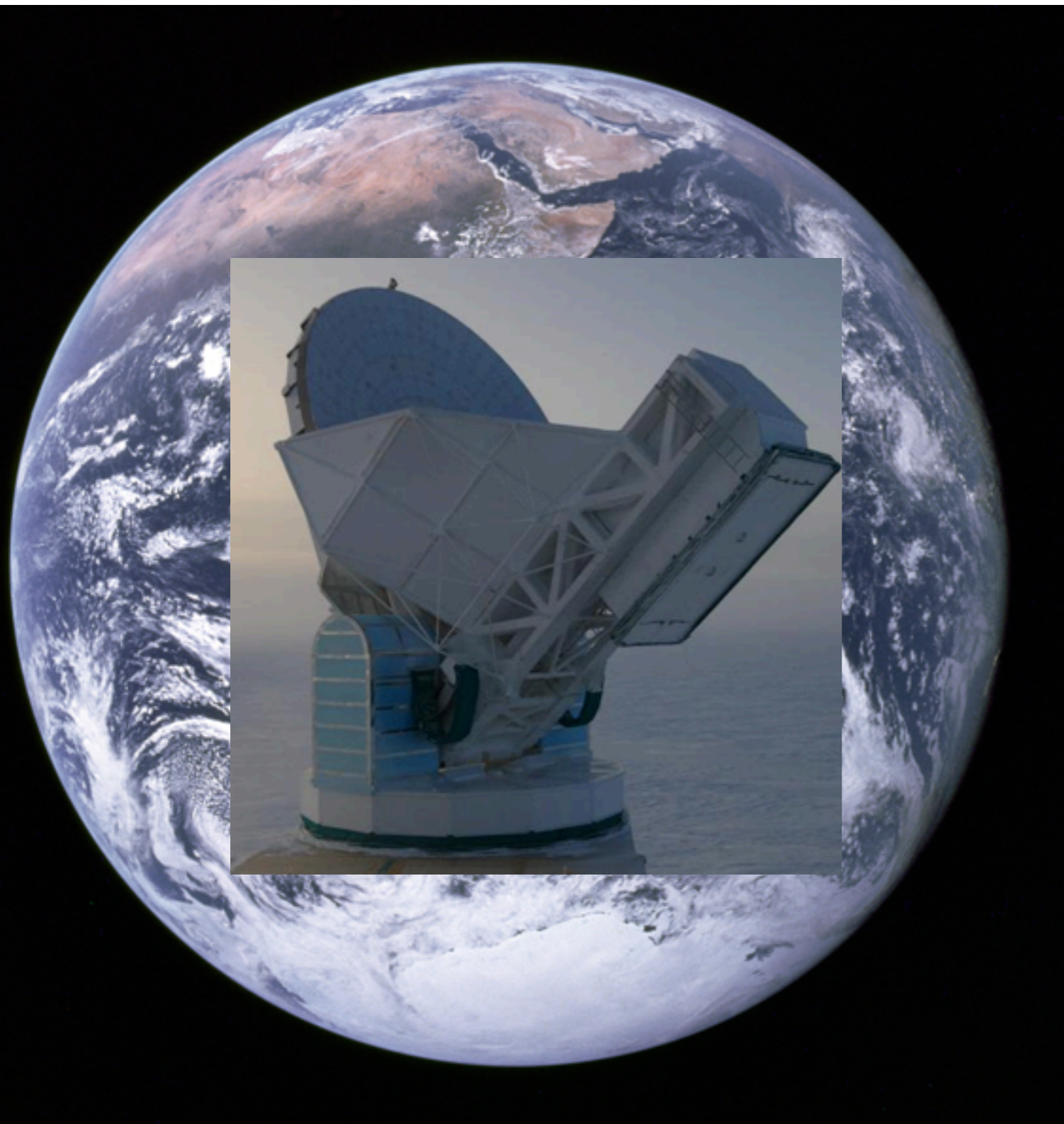
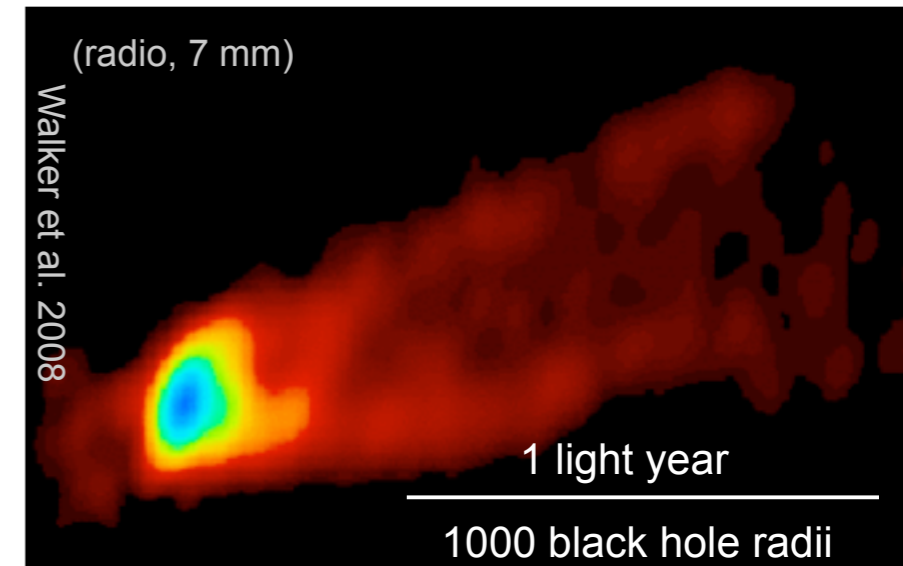


(Doeleman et al. 2008)



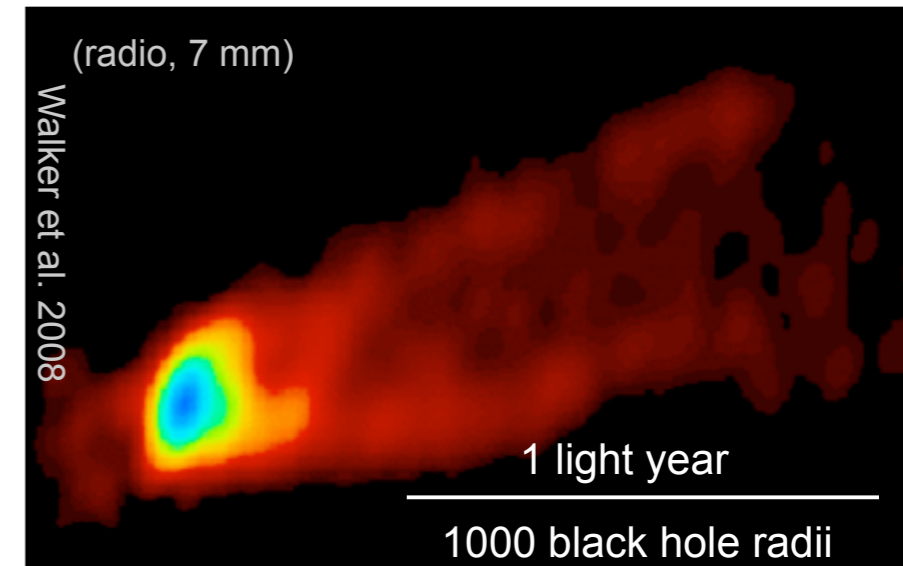
(Doeleman et al. 2012)

Event Horizon Telescope (EHT): VLBI images of Black Holes

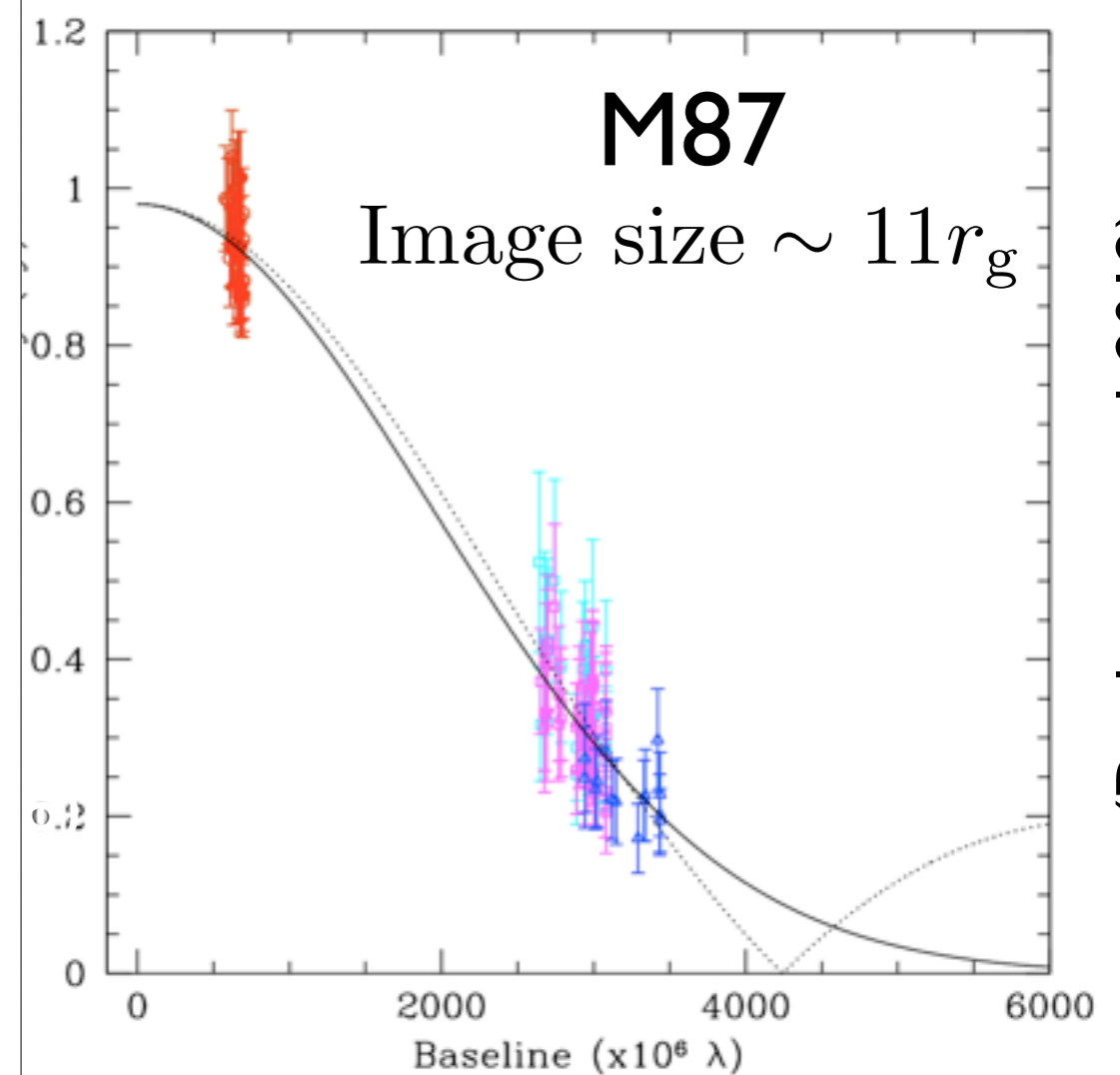


(Doeleman et al. 2012)

Event Horizon Telescope (EHT): VLBI images of Black Holes

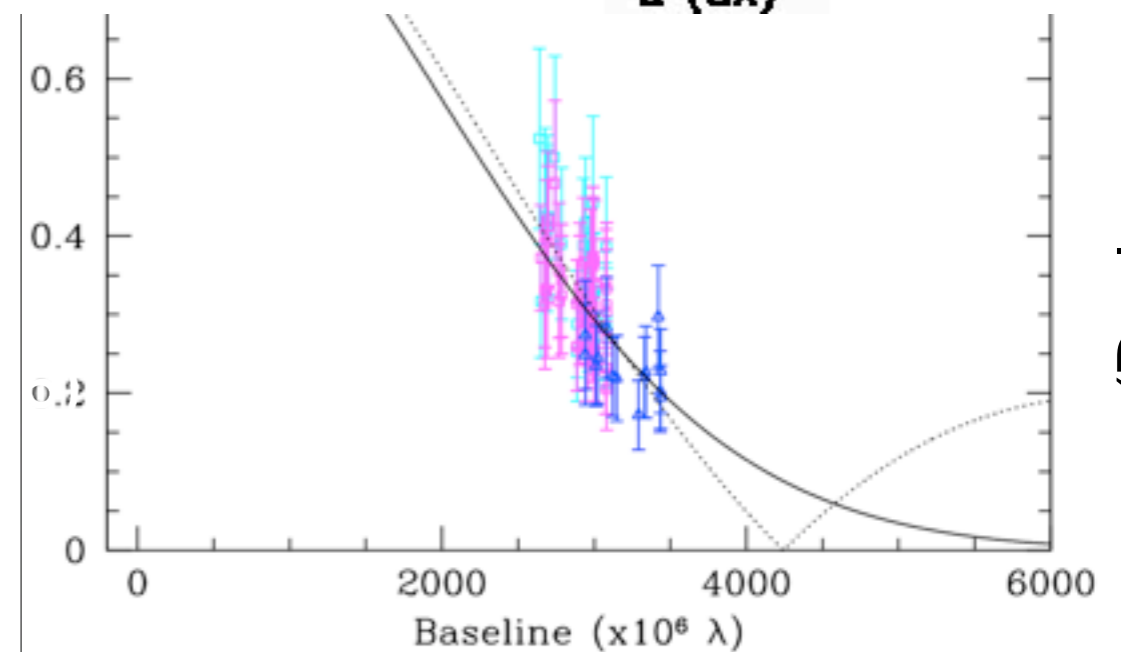
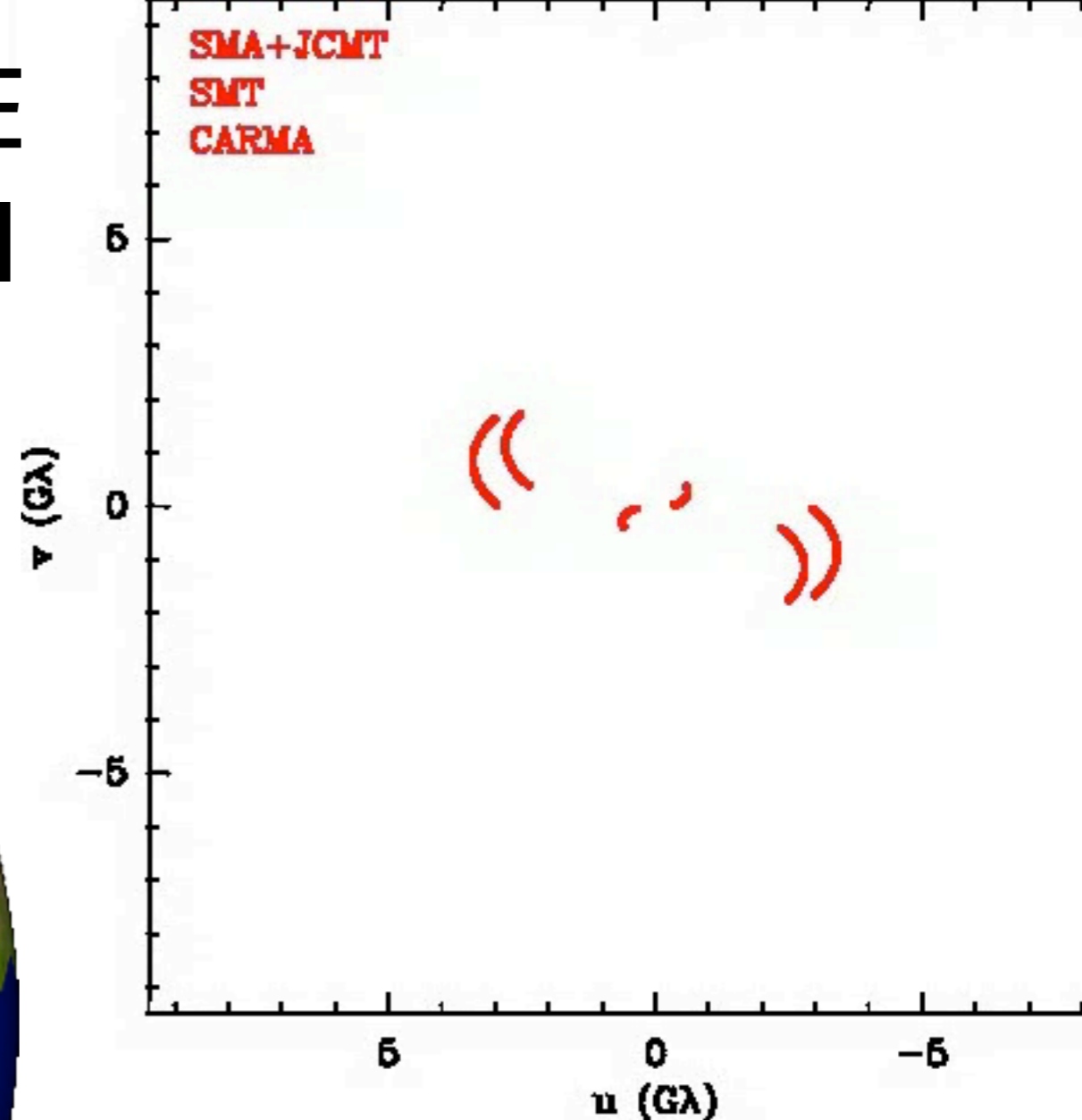


ky

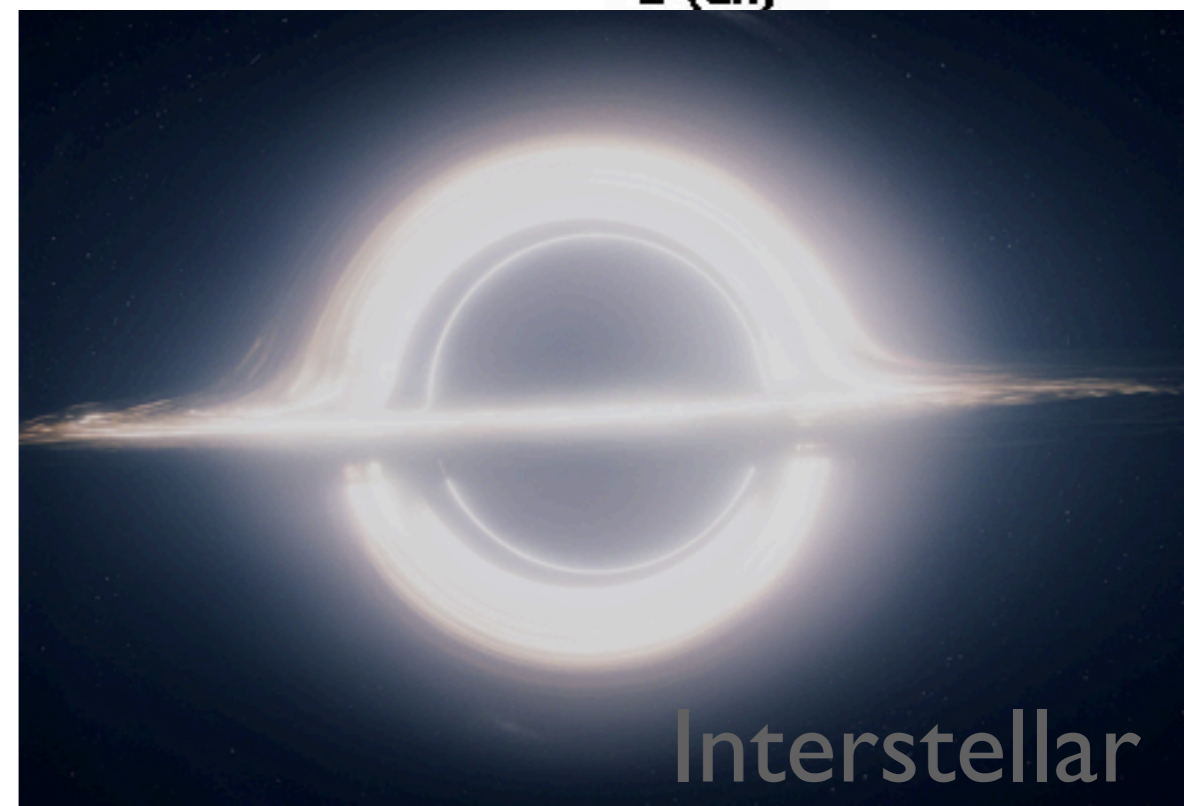
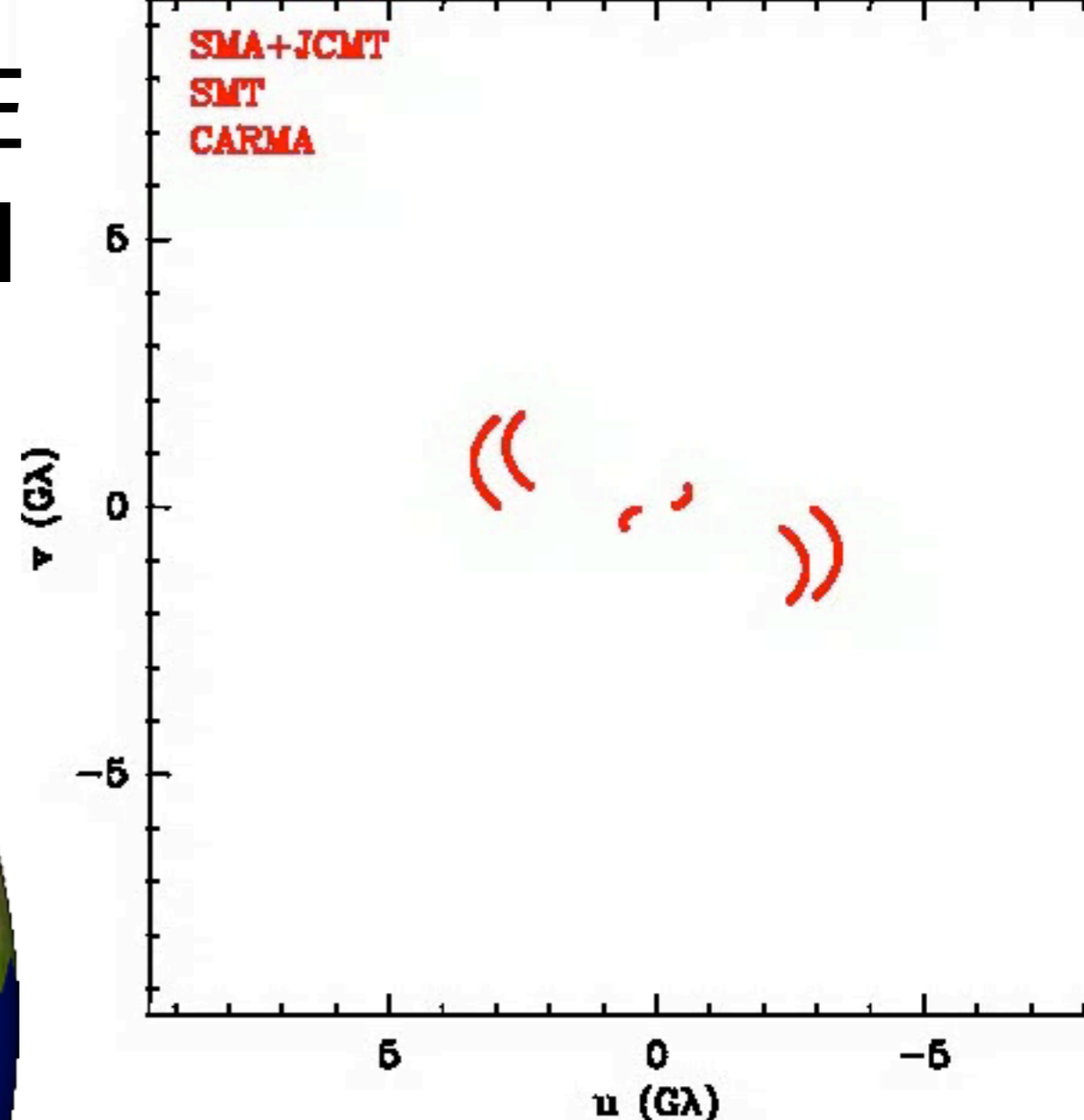


(Doeleman et al. 2012)

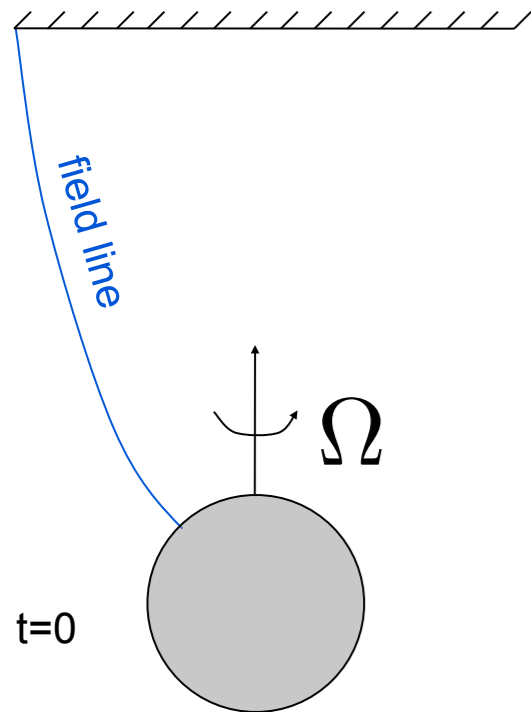
Event Horizon Telescope (EHT) VLBI images of Black Holes



Event Horizon Telescope (EHT) VLBI images of Black Holes

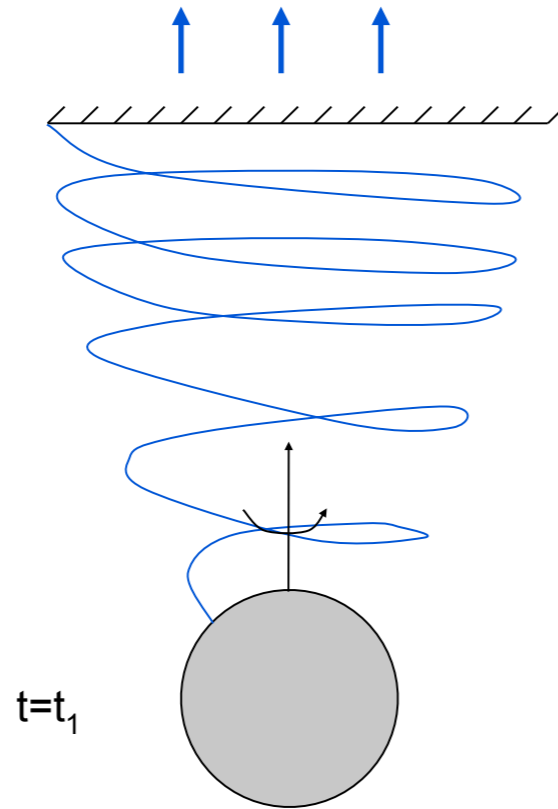
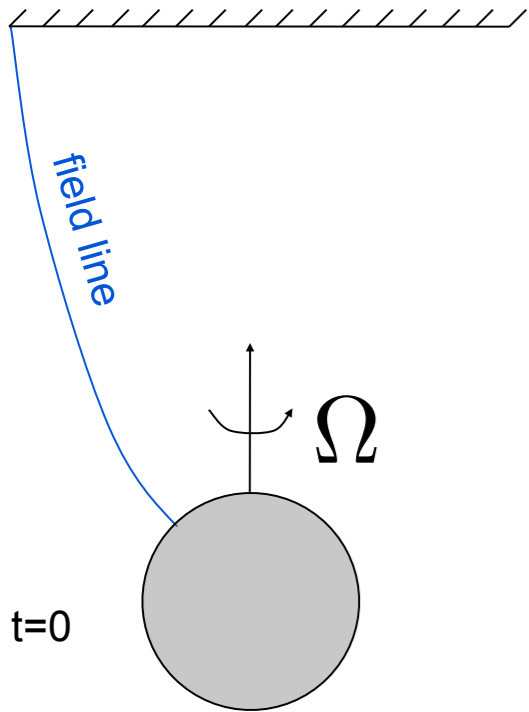


Jets 101



Jets 101

$$P = \frac{B_{\varphi}^2}{8\pi}$$

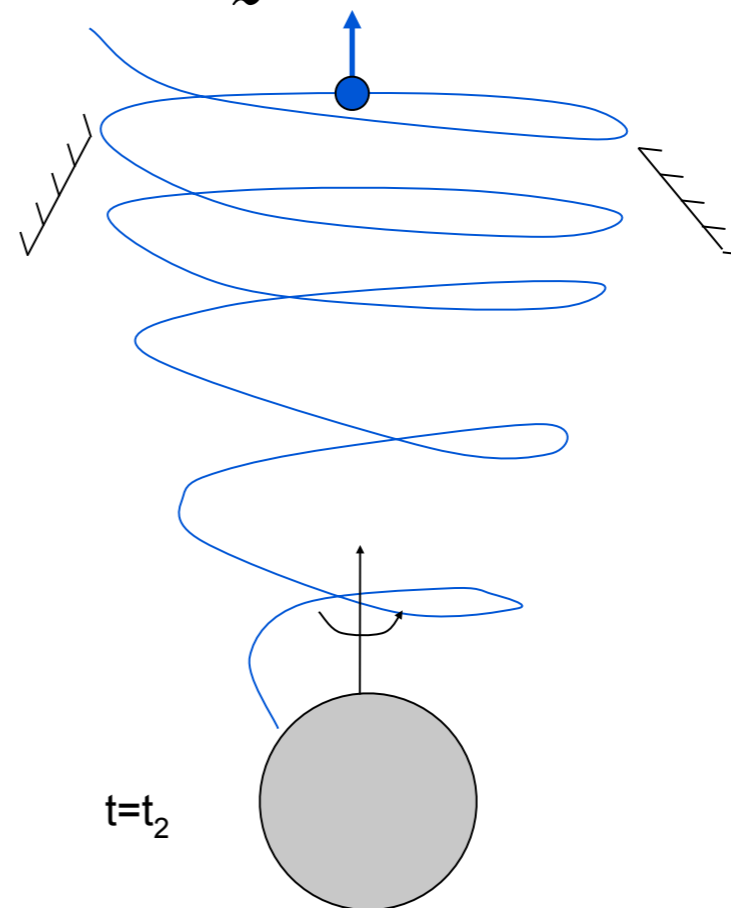
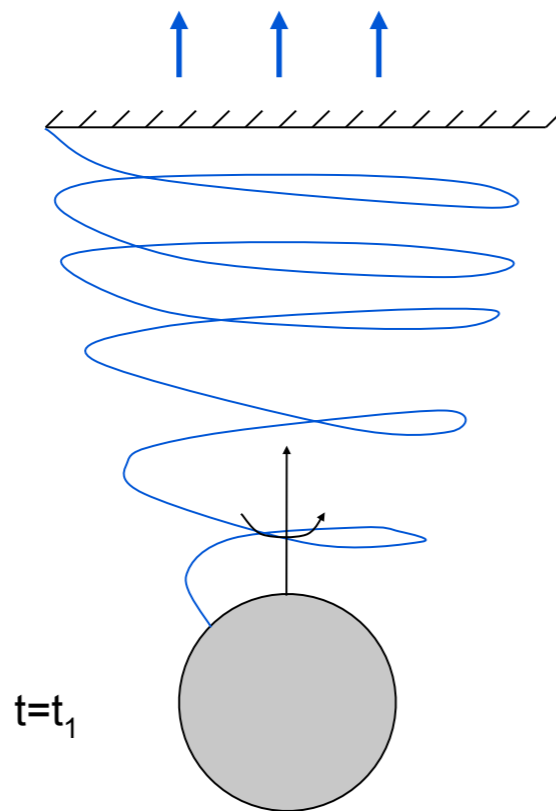
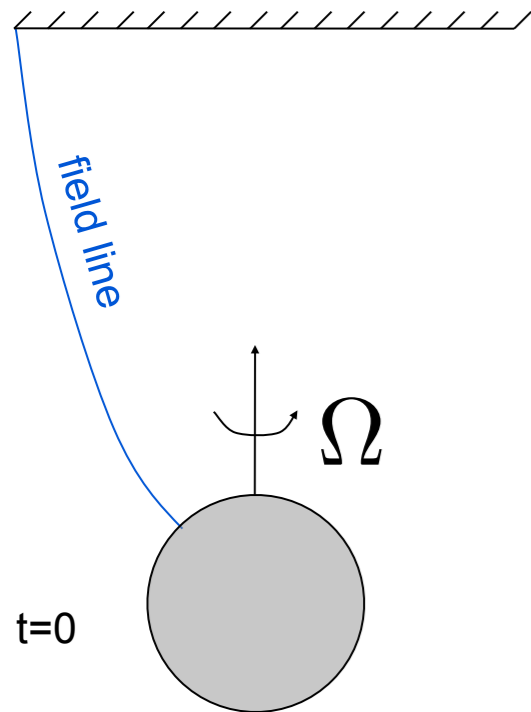


Jets 101

$$P = \frac{B_{\varphi}^2}{8\pi}$$

Field toroidally-
dominated

$$B_{\varphi} \gg B_z$$



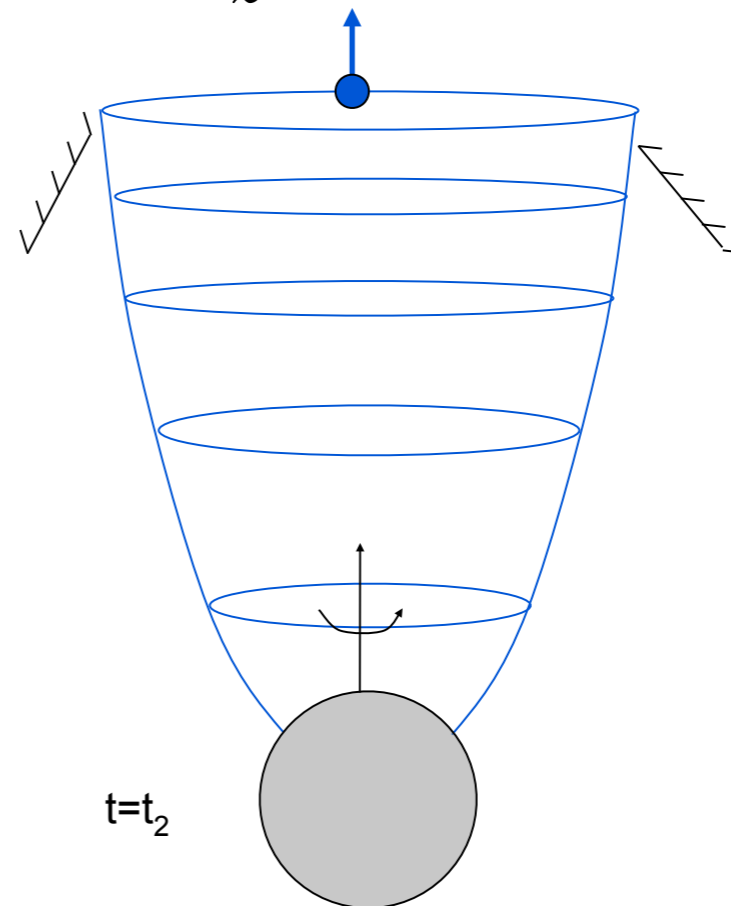
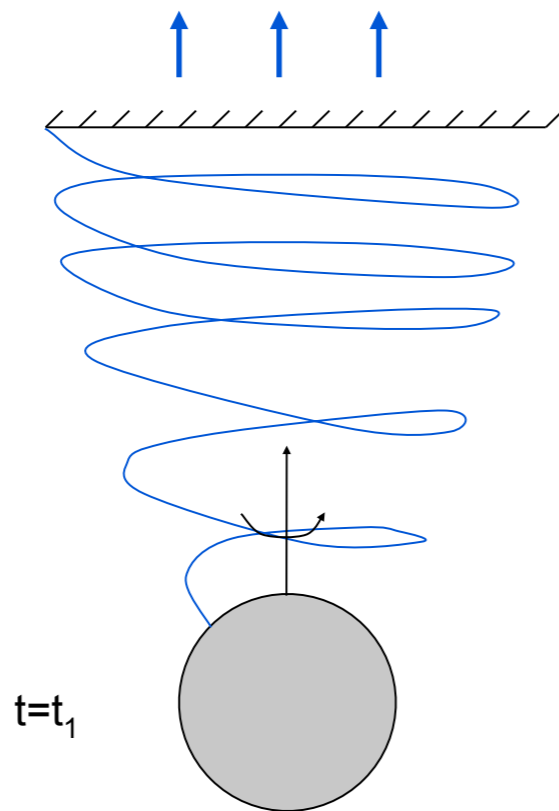
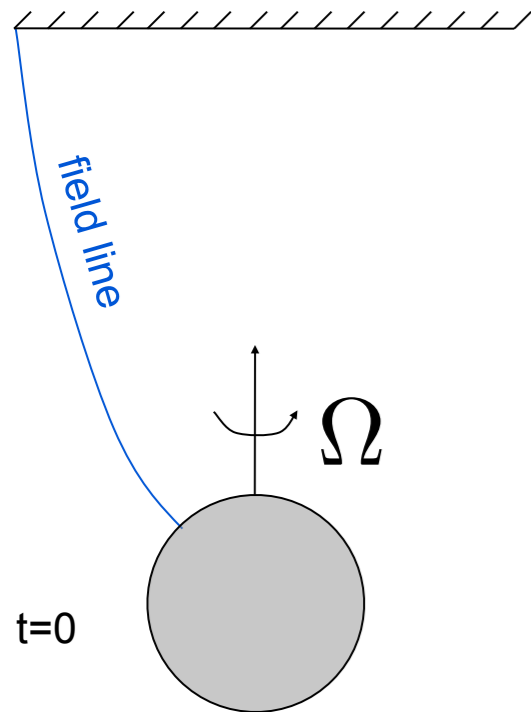
(Beskin &
Nokhrina 06,
Komissarov 07-10,
AT+08,09,10,12,
Lyubarsky 10,
AT 15)

Jets 101

$$P = \frac{B_{\varphi}^2}{8\pi}$$

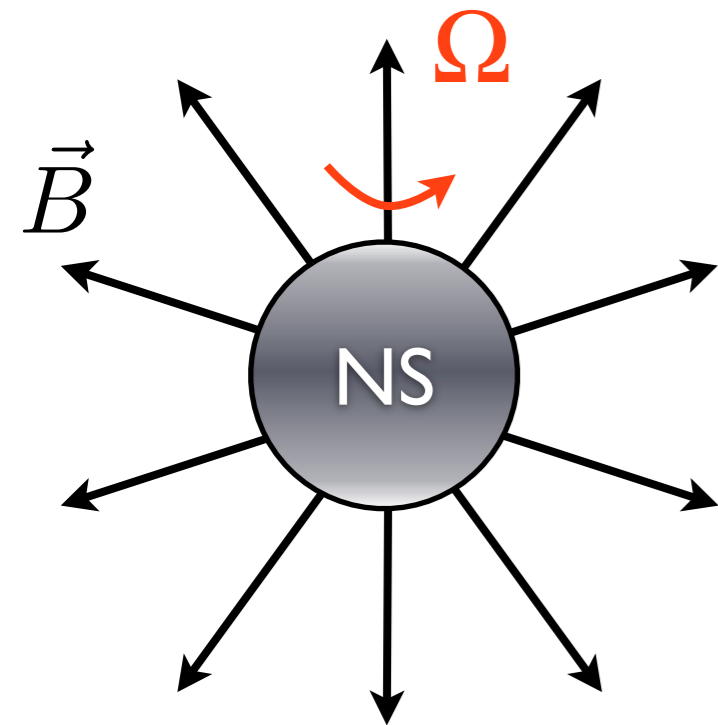
Field toroidally-
dominated

$$B_{\varphi} \gg B_z$$

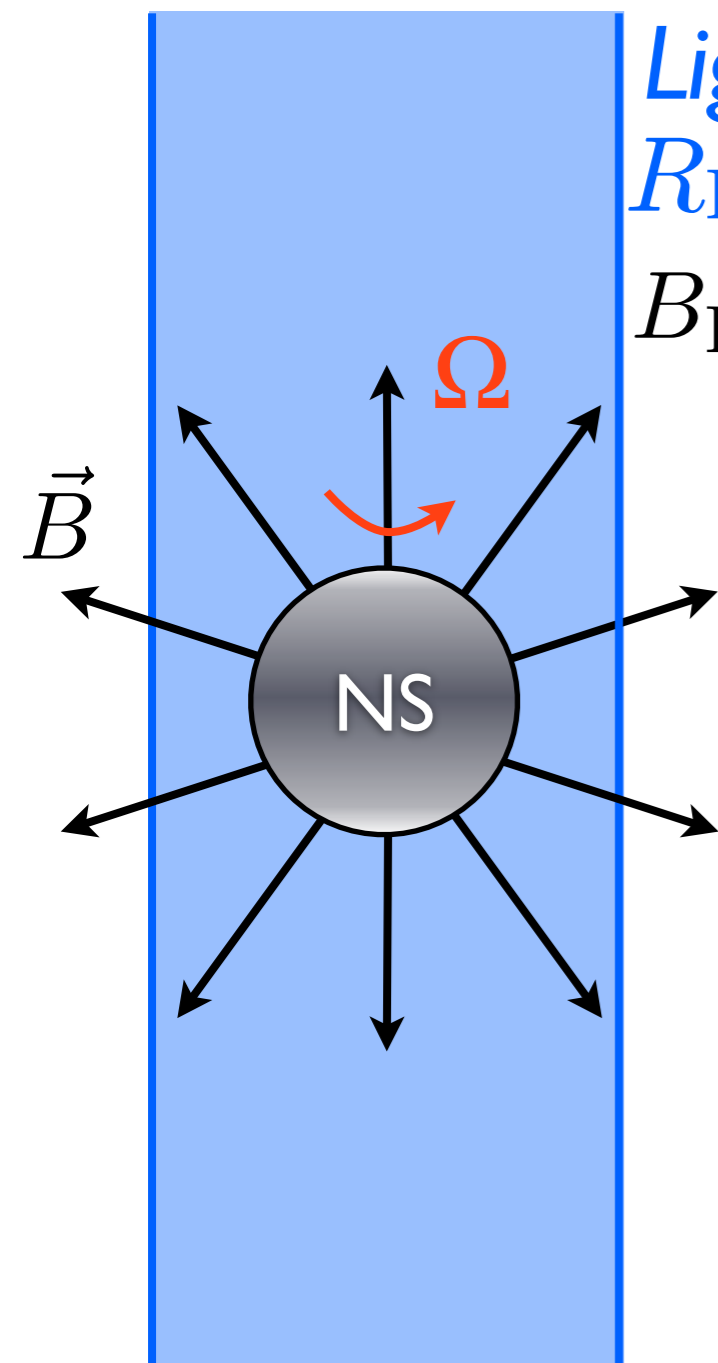


(Beskin &
Nokhrina 06,
Komissarov 07-10,
AT+08,09,10,12,
Lyubarsky 10,
AT 15)

What Powers Outflow?



What Powers Outflow?



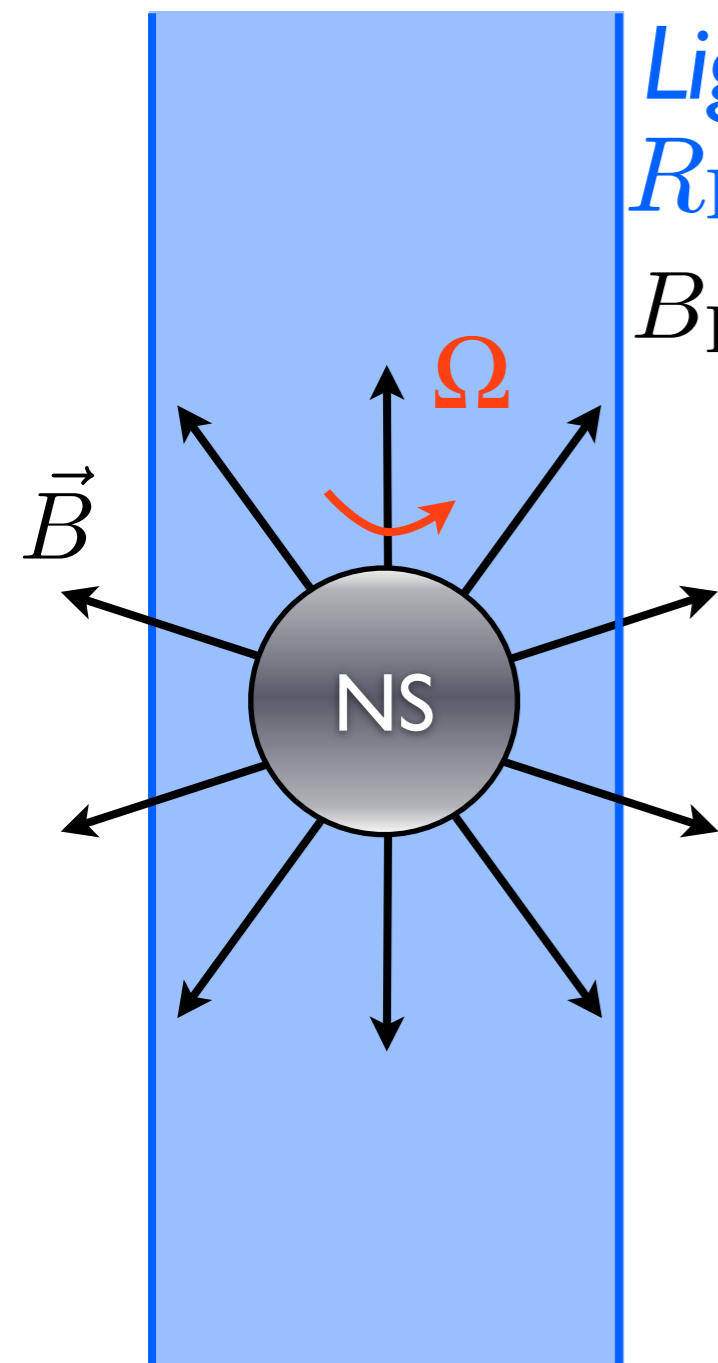
Light cylinder (LC):

$$R_L = c/\Omega$$

$$B_L = \Phi/2\pi R_L$$

- Flow separates from NS at LC

What Powers Outflow?



Light cylinder (LC):

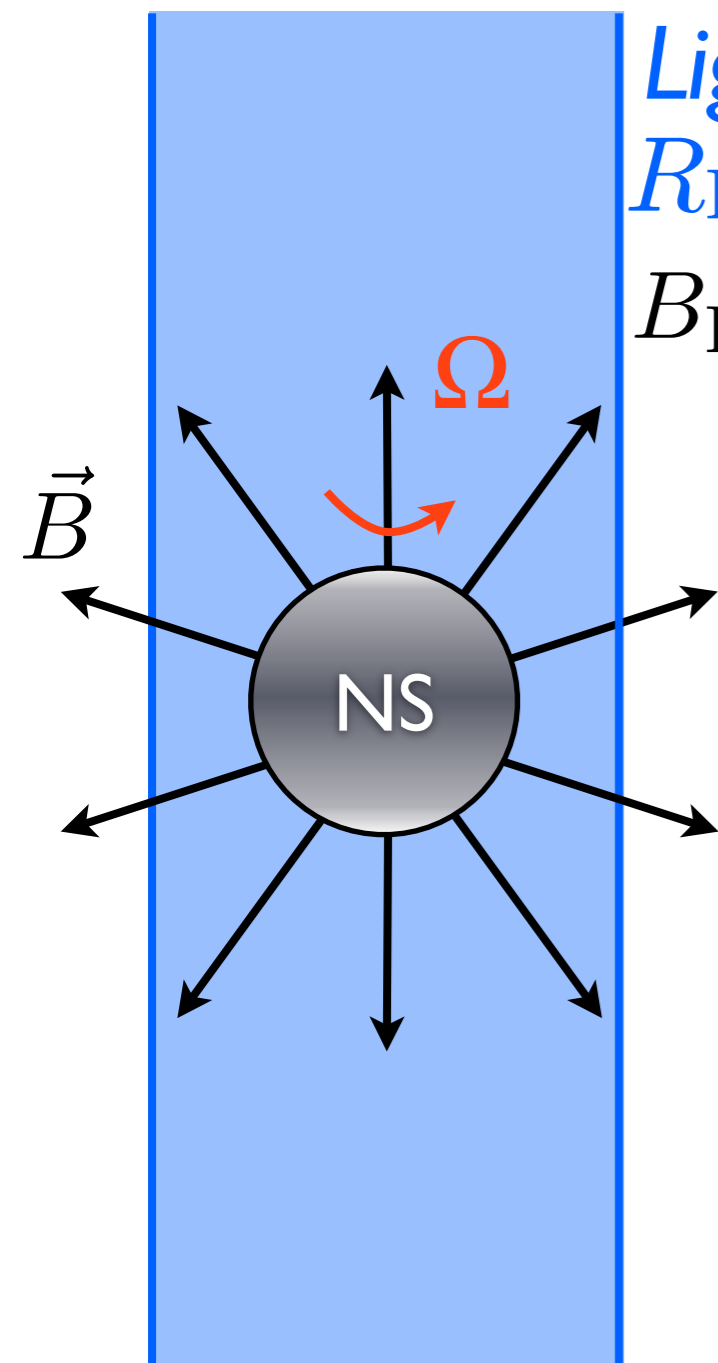
$$R_L = c/\Omega$$

$$B_L = \Phi/2\pi R_L$$

- Flow separates from NS at LC
- Spindown power

$$P \sim \frac{c}{4\pi} (\vec{E} \times \vec{B}) \times 4\pi R_L^2 = c B_L^2 R_L^2$$

What Powers Outflow?



Light cylinder (LC):

$$R_L = c/\Omega$$

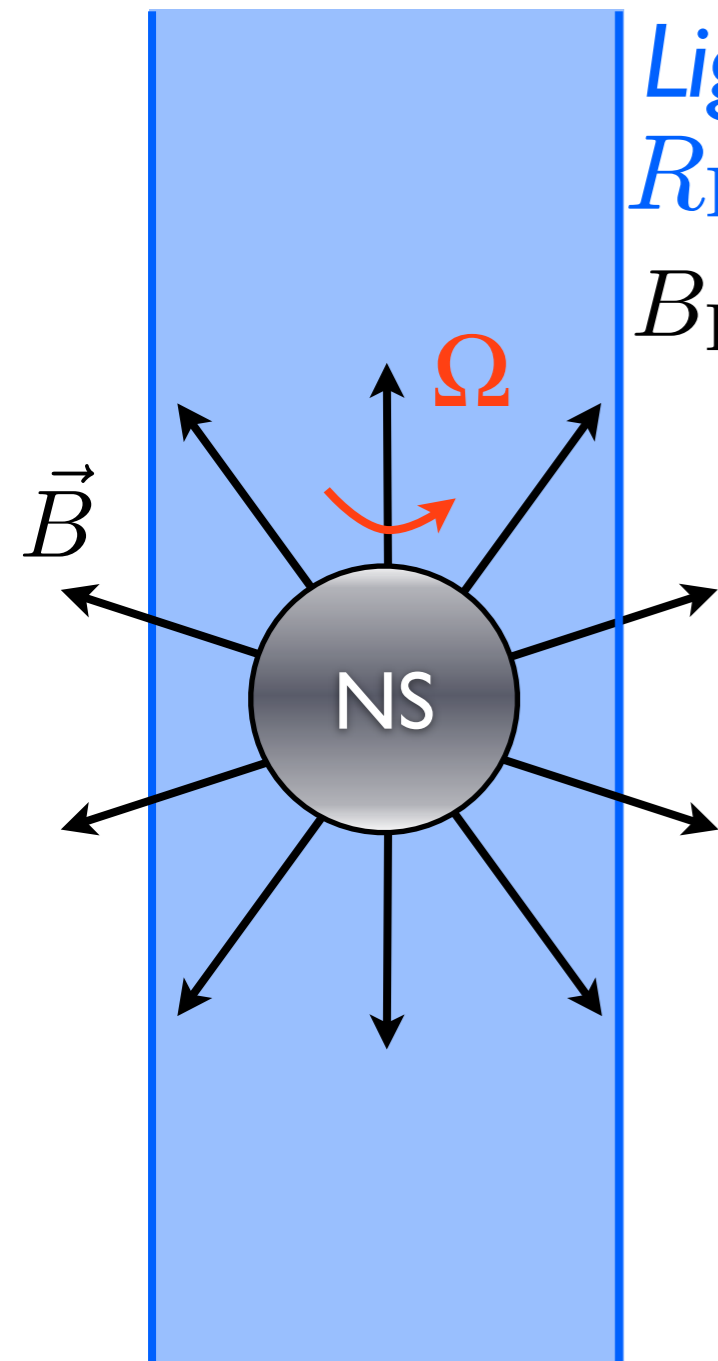
$$B_L = \Phi/2\pi R_L$$

- Flow separates from NS at LC
- Spindown power

$$P \sim \frac{c}{4\pi} (\vec{E} \times \vec{B}) \times 4\pi R_L^2 = c B_L^2 R_L^2$$

$$P \sim \frac{1}{4\pi^2 c} \Phi^2 \Omega^2$$

What Powers Outflow?



Light cylinder (LC):

$$R_L = c/\Omega$$

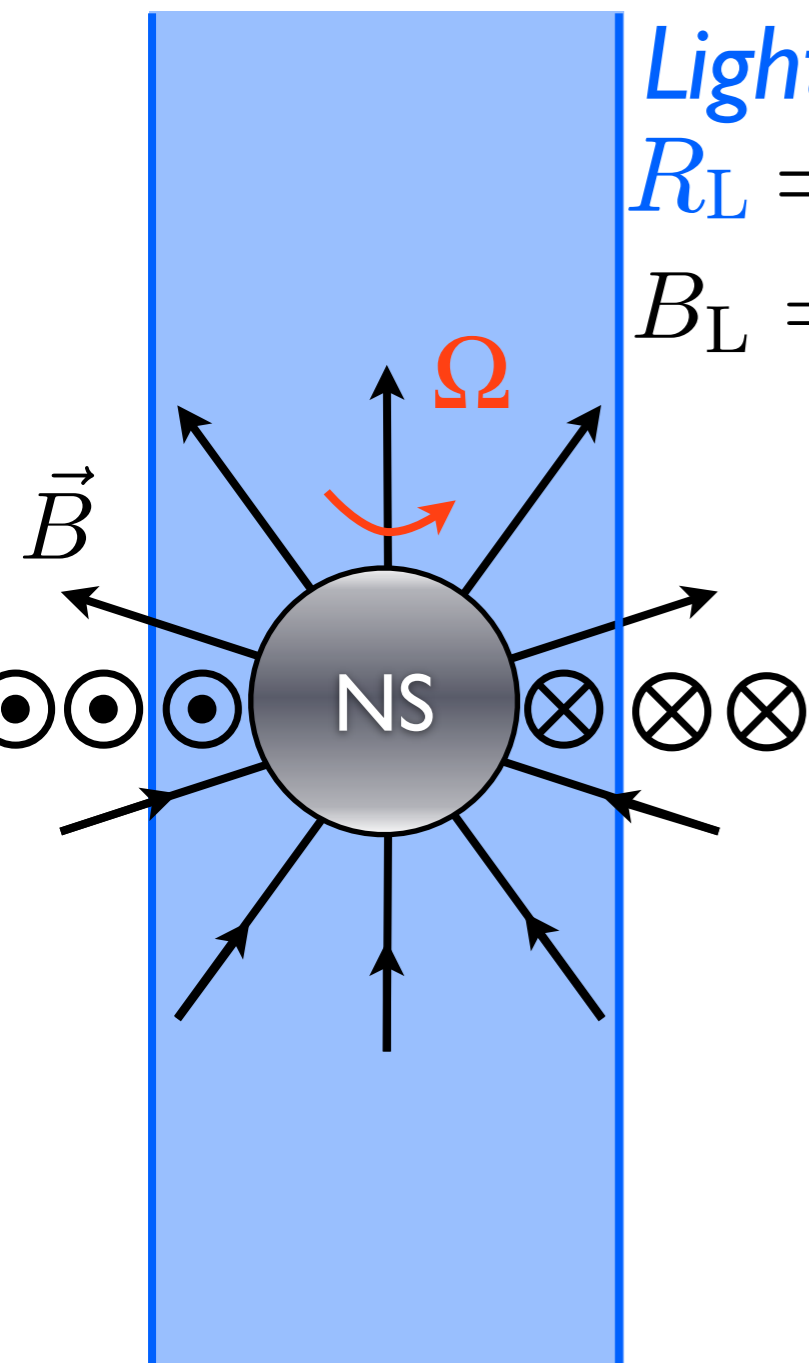
$$B_L = \Phi/2\pi R_L$$

- Flow separates from NS at LC
- Spindown power

$$P \sim \frac{c}{4\pi} (\vec{E} \times \vec{B}) \times 4\pi R_L^2 = c B_L^2 R_L^2$$

$$P \sim \frac{1}{64\pi^2 c} \Phi^2 \Omega^2$$

What Powers Outflow?



Light cylinder (LC):

$$R_L = c/\Omega$$

$$B_L = \Phi/2\pi R_L$$

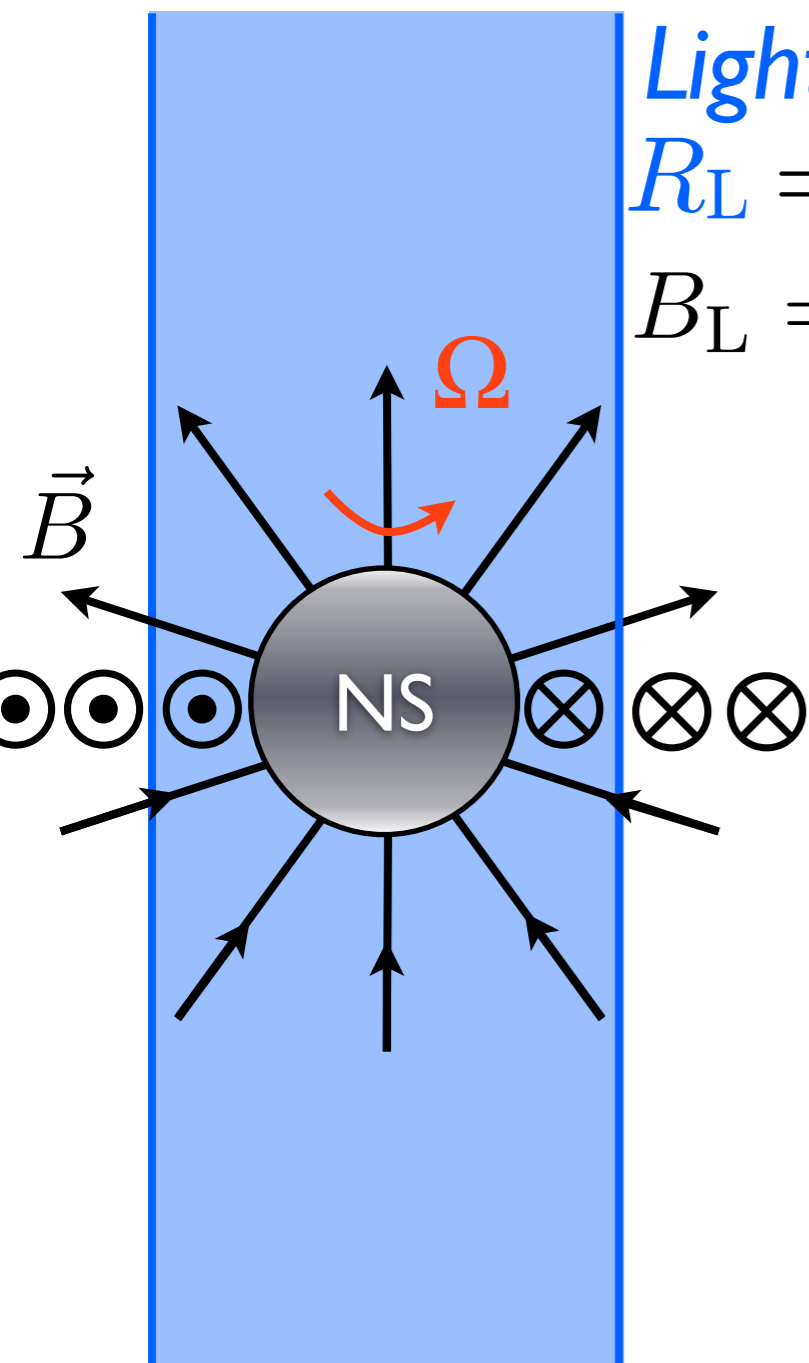
- Flow separates from NS at LC
- Spindown power

$$P \sim \frac{c}{4\pi} (\vec{E} \times \vec{B}) \times 4\pi R_L^2 = c B_L^2 R_L^2$$

$$P \sim \frac{1}{64\pi^2 c} \Phi^2 \Omega^2$$

- Split-monopole

What Powers Outflow?



Light cylinder (LC):

$$R_L = c/\Omega$$

$$B_L = \Phi/2\pi R_L$$

- Flow separates from NS at LC
- Spindown power

$$P \sim \frac{c}{4\pi} (\vec{E} \times \vec{B}) \times 4\pi R_L^2 = c B_L^2 R_L^2$$

$$P \sim \frac{1}{64\pi^2 c} \Phi^2 \Omega^2$$

- Split-monopole
- What about black holes?

A Black Hole is VERY Simple

- Mass: **M**
- Spin: a ($J=a \mathbf{GM}^2/c$)
- Charge: Q

**A Black Hole has no Hair! (No Hair
Theorem)**

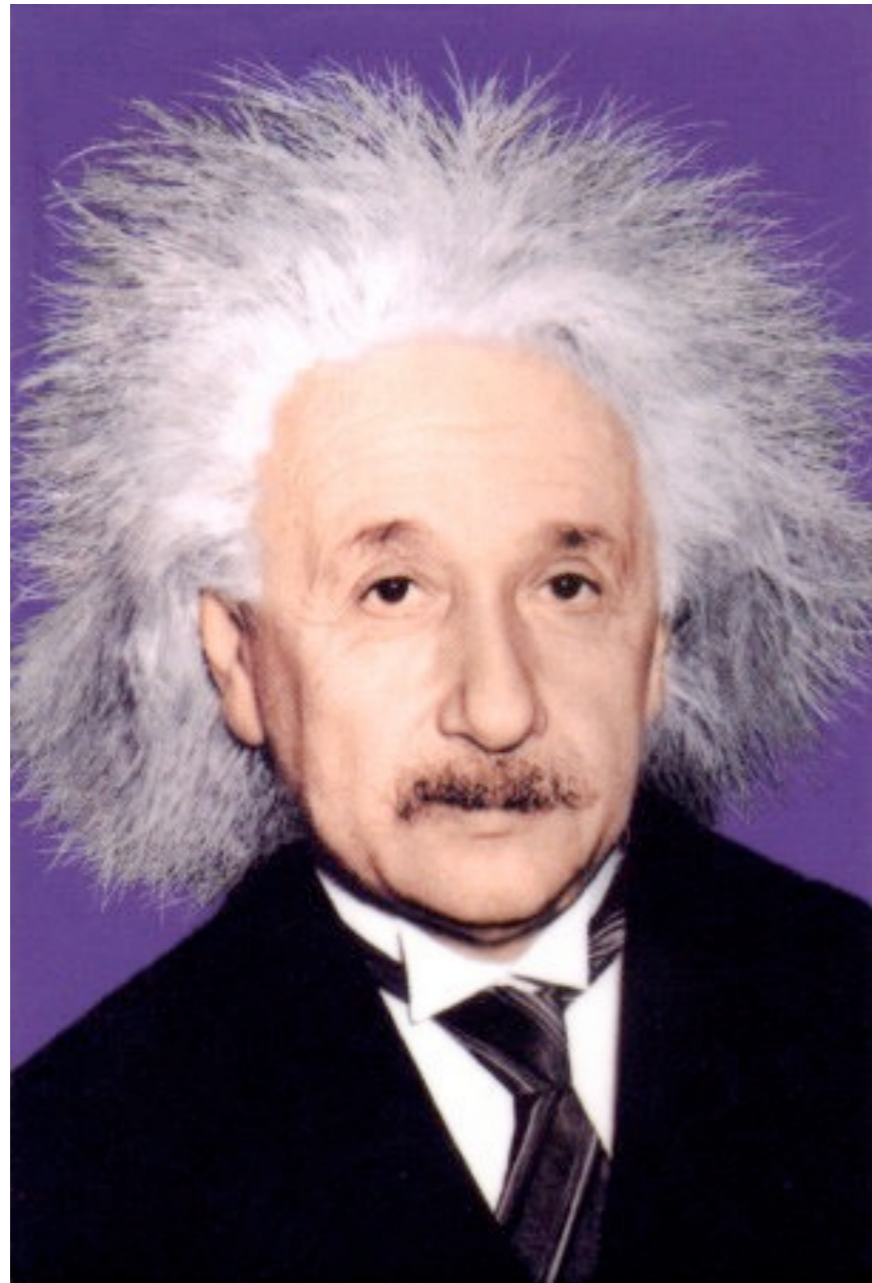
To be precise, a BH has 2 (at most 3) hairs

A Black Hole is VERY Simple

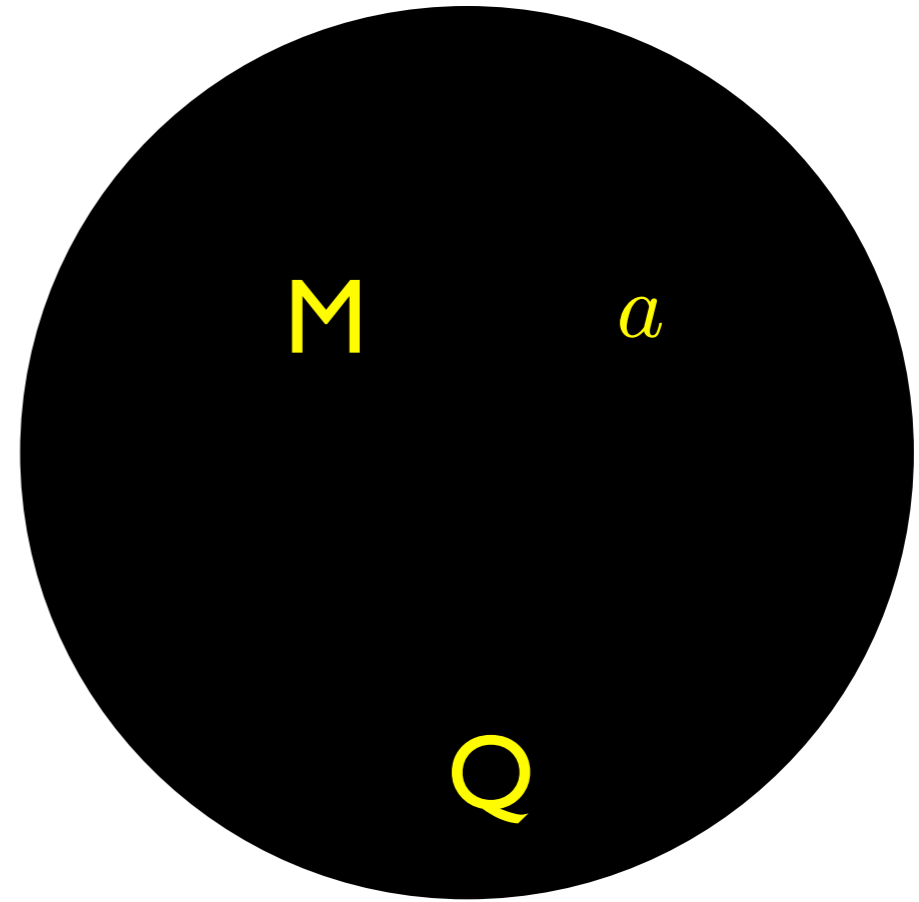
- Mass: **M**
- Spin: a ($J=a \mathbf{GM}^2/c$)
- ~~Charge: Q~~

**A Black Hole has no Hair! (No Hair
Theorem)**

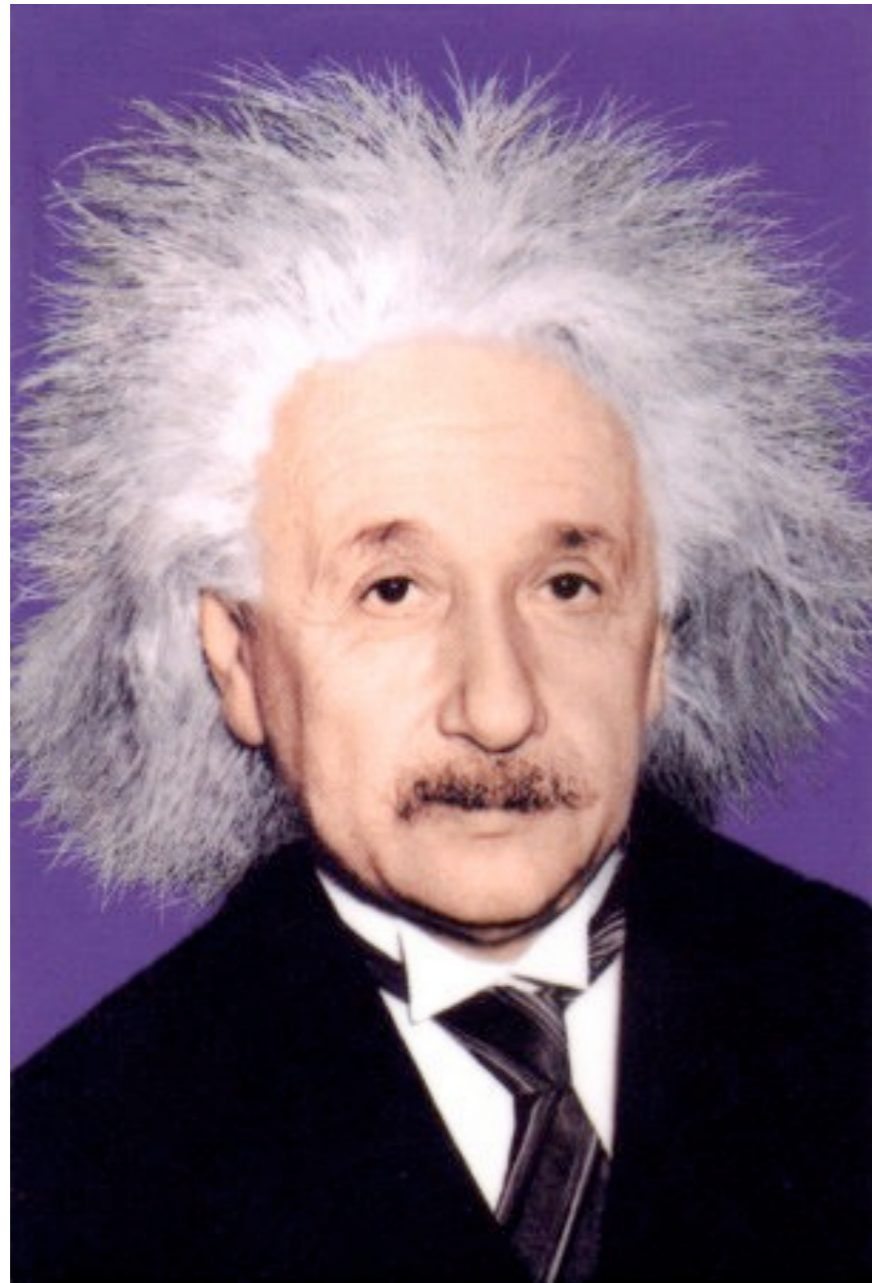
To be precise, a BH has 2 (at most 3) hairs



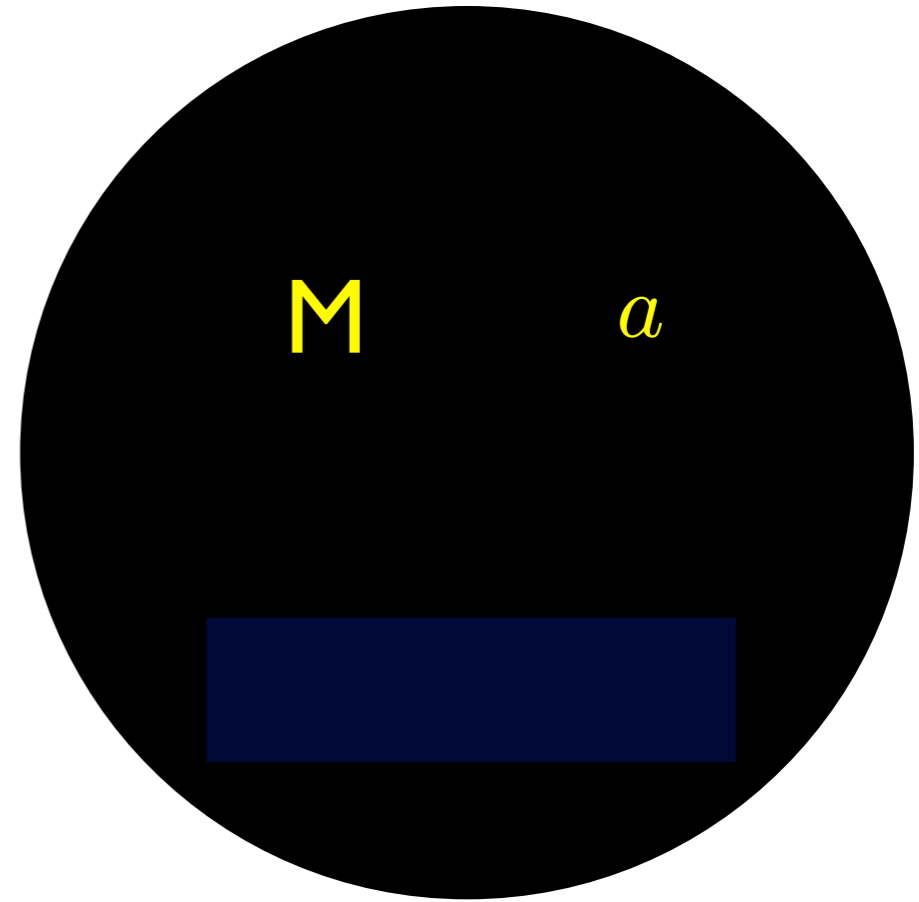
Einstein had a lot
of hair!



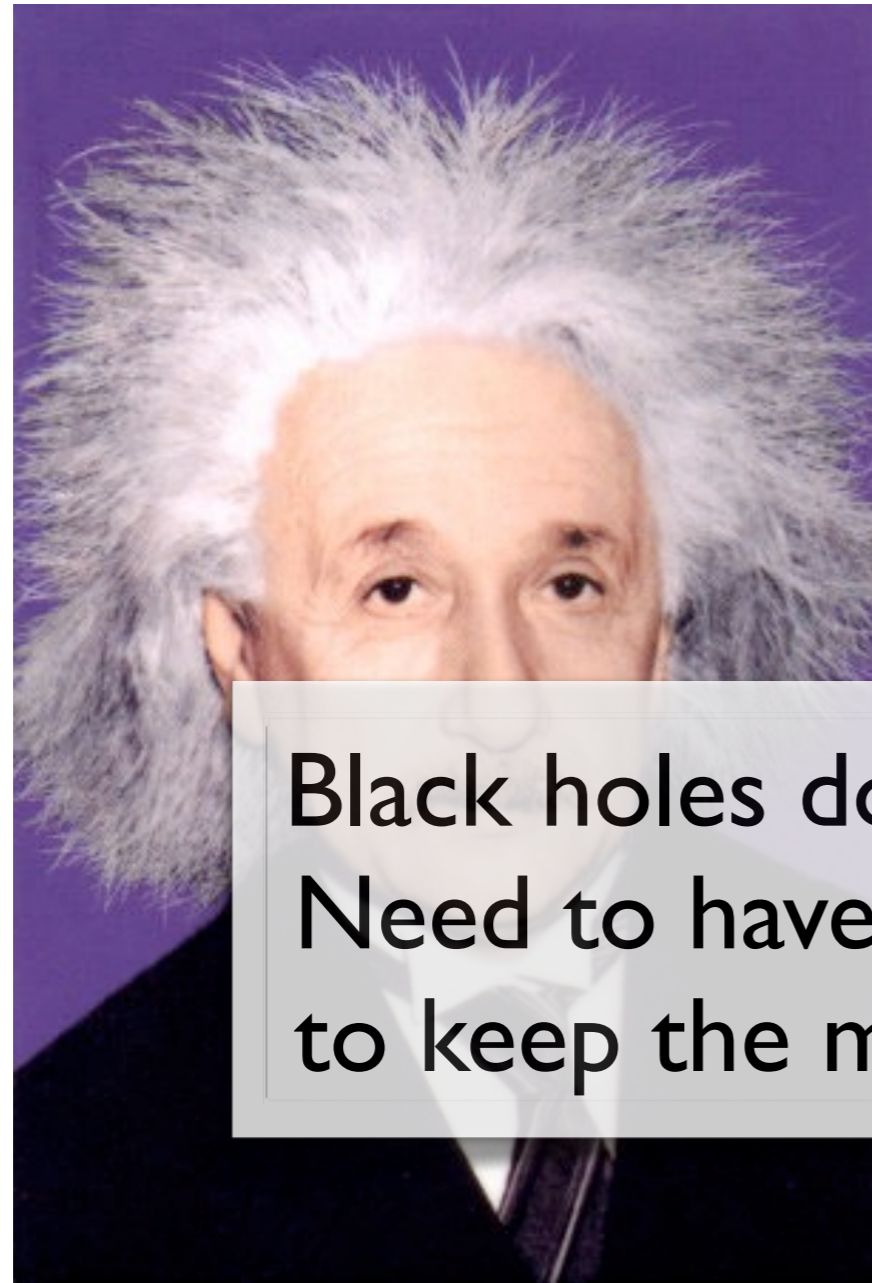
Black Hole has
3 hairs!



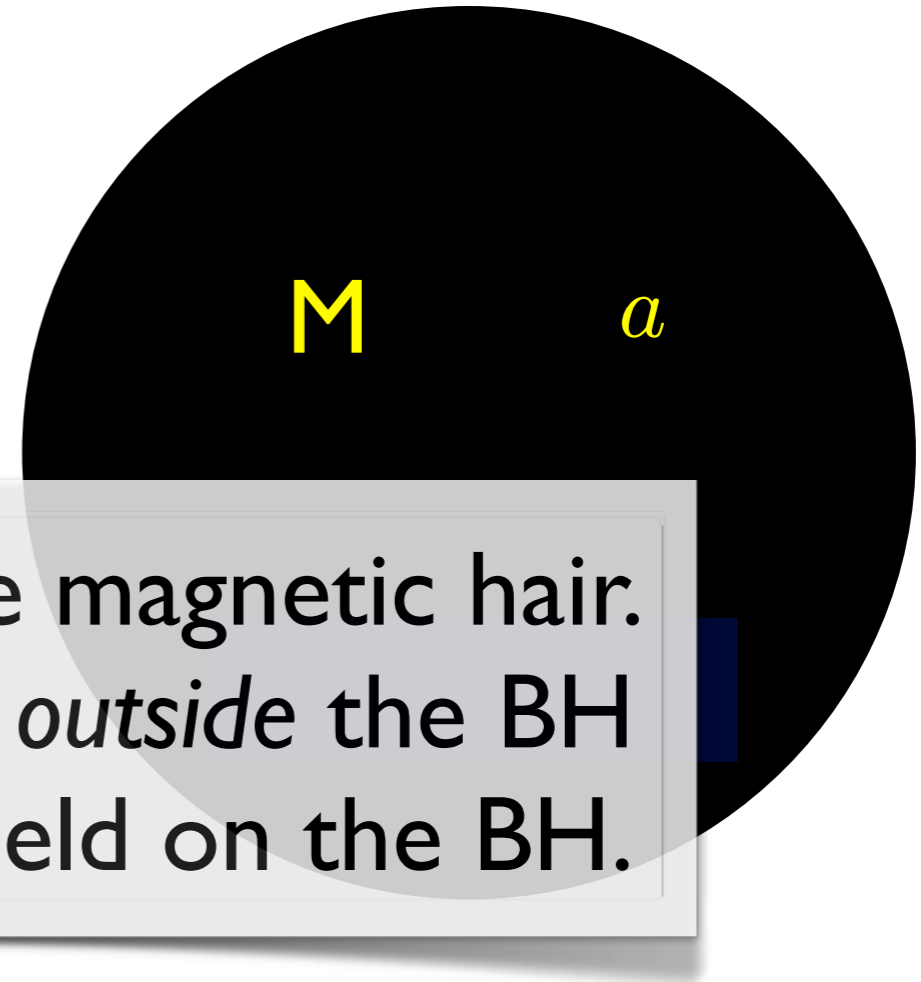
Einstein had a lot
of hair!



A Black Hole
has only
2 hairs



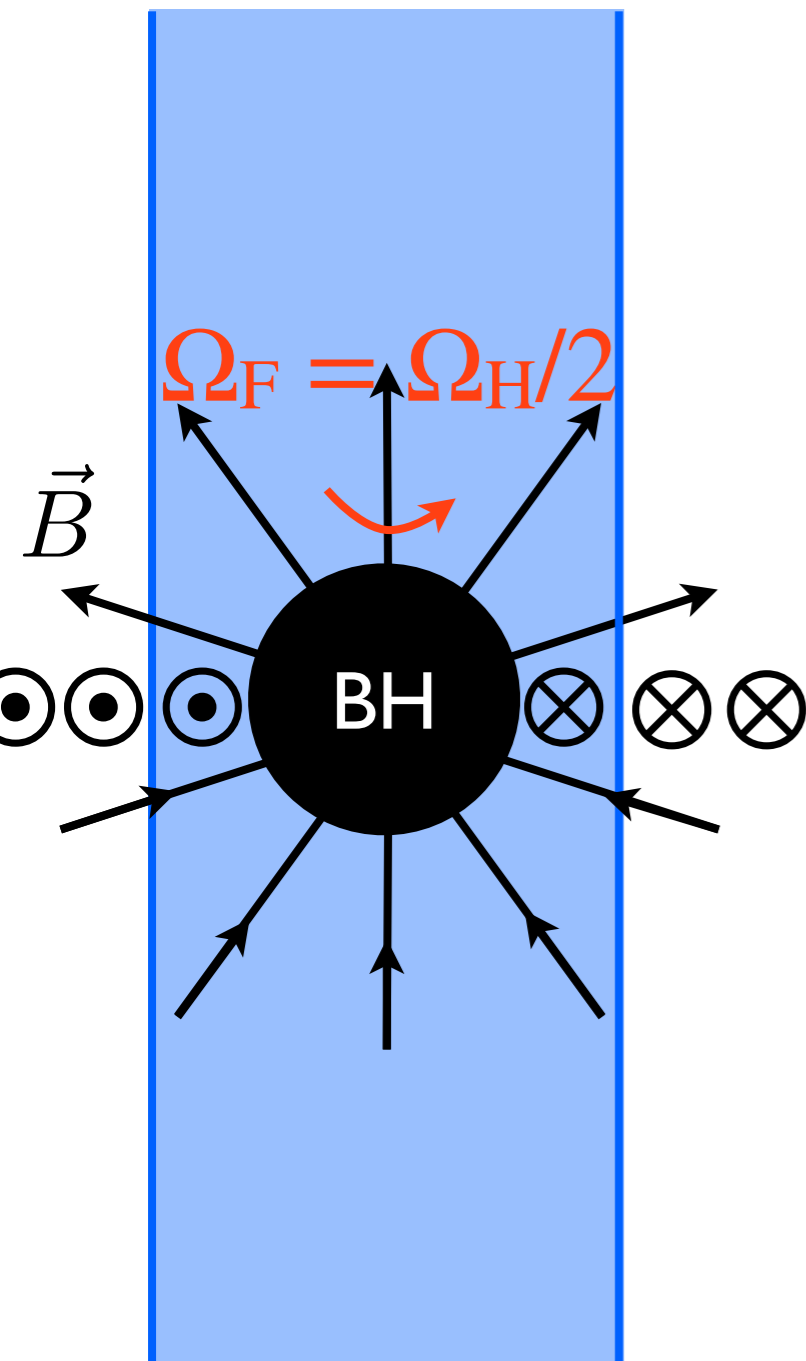
Einstein had a lot
of hair!



Black holes do not have magnetic hair.
Need to have currents *outside* the BH
to keep the magnetic field on the BH.

A Black Hole
has only
2 hairs

What about Black Holes?



- Black hole drags space-time at

$$\omega \simeq \Omega_H (r/r_H)^{-3}, \quad \Omega_H = ac/2r_H$$

- At the event horizon $\omega = \Omega_H$

- At infinity $\omega = 0$

- Field line tries to please both:

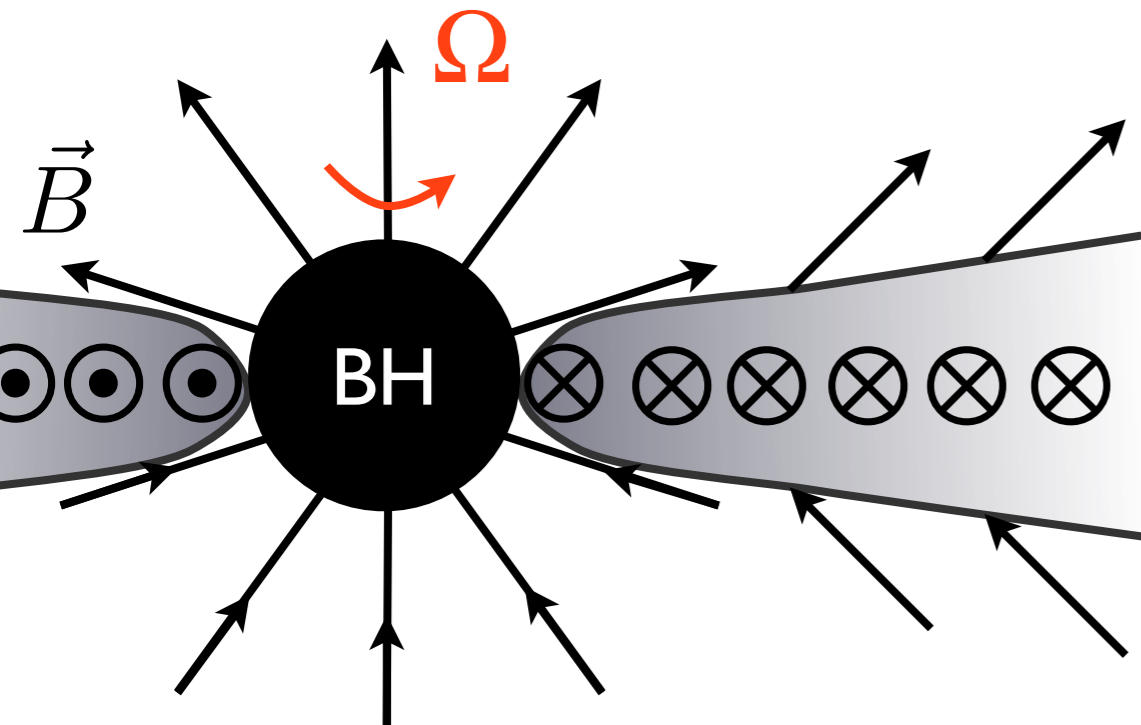
$$\Omega_F = \Omega_H/2$$

- Otherwise, behaves almost like a NS!

$$P \sim \frac{1}{64\pi^2 c} \Phi^2 \Omega_F^2 \sim \frac{1}{2416\pi^2 c} \Phi^2 \Omega_H^2$$

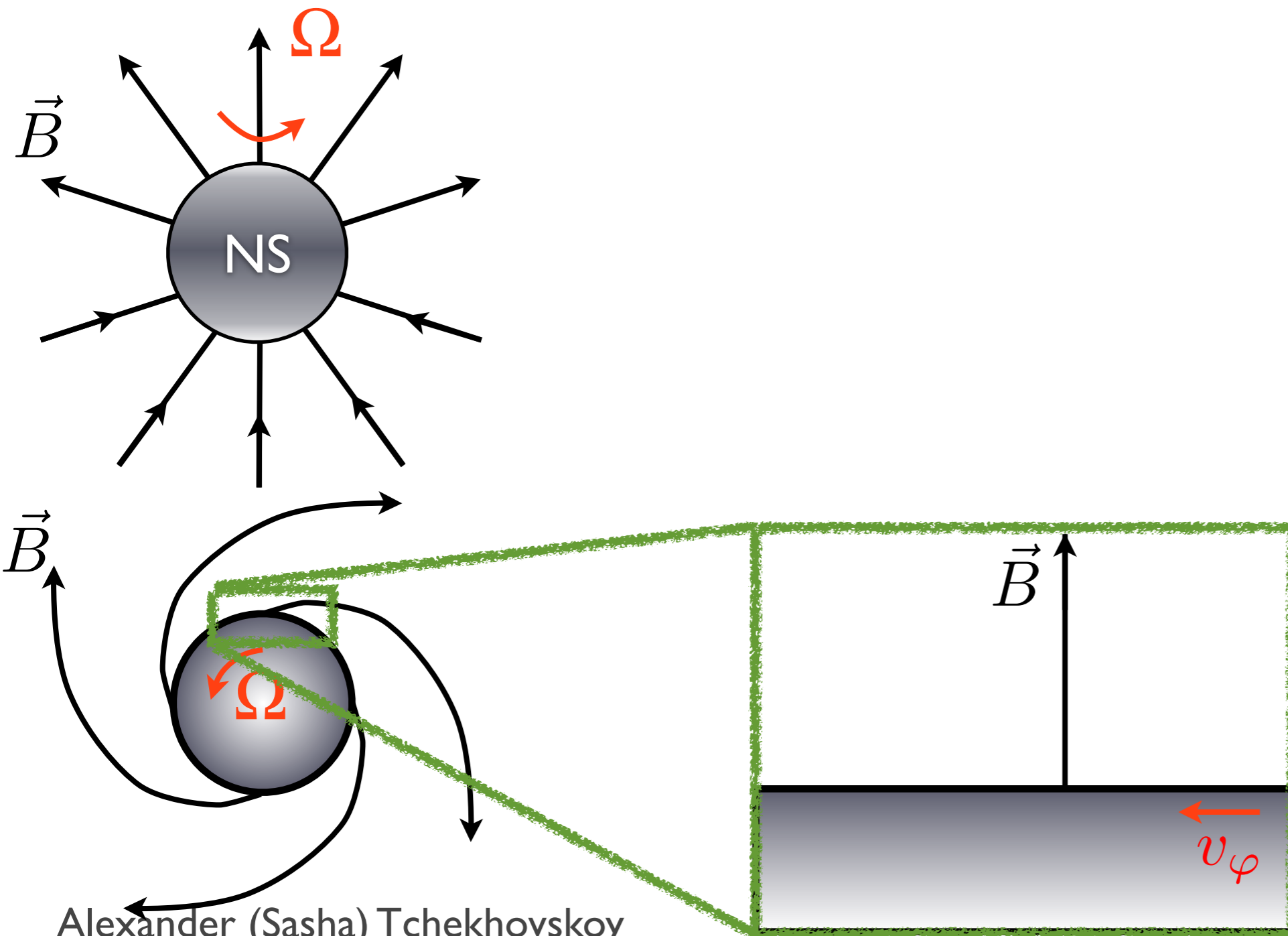
(~10% corrections for other field geometries, AT+10, AT15)

Where Does Φ Come from?

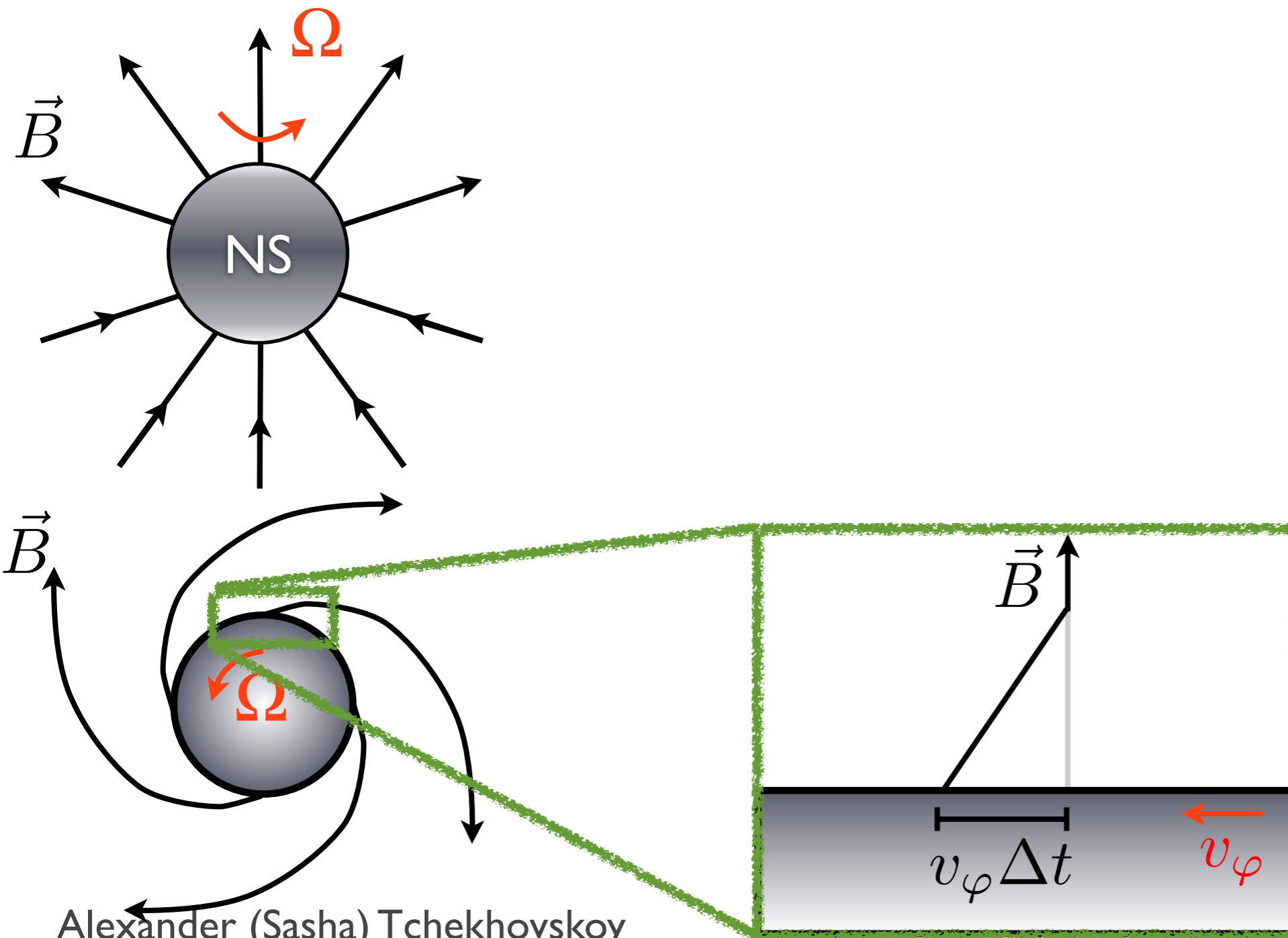


- Accretion disk:
 - either drags B from large scales
 - or generates B in situ
 - presently unsolved problem
- *Black hole must be accreting in order to form magnetosphere and produce jets*

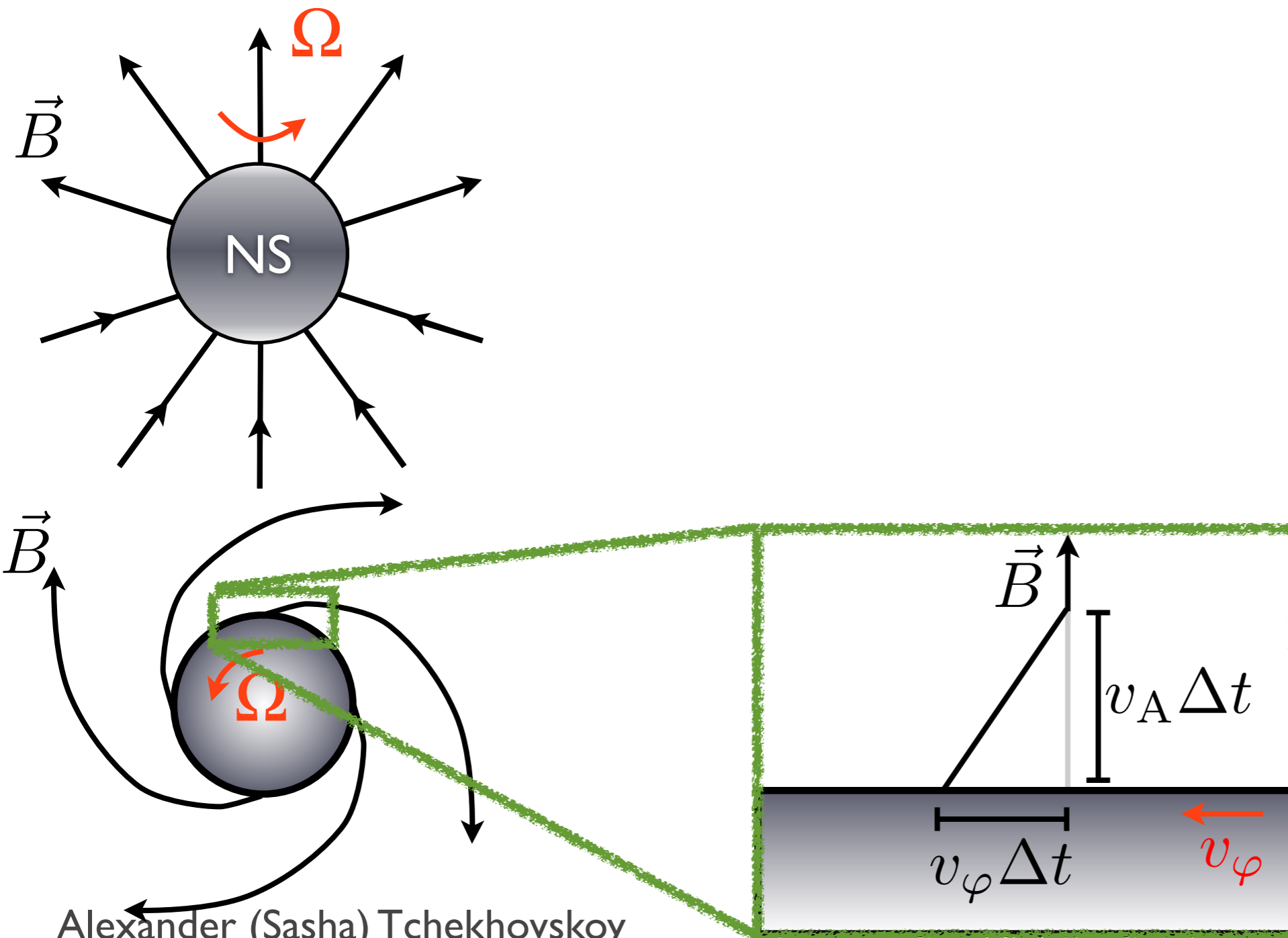
How do Jets Accelerate?



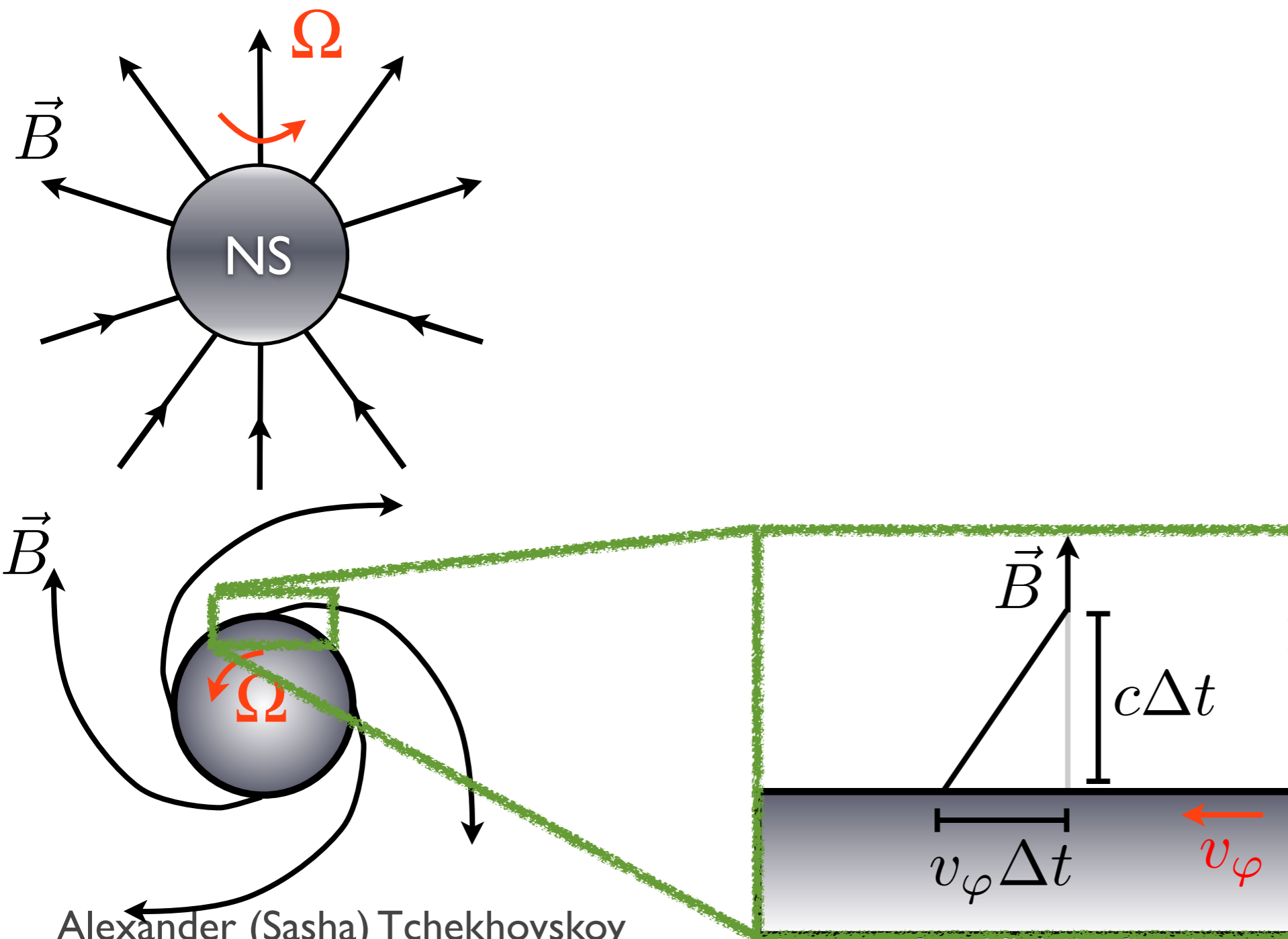
How do Jets Accelerate?



How do Jets Accelerate?

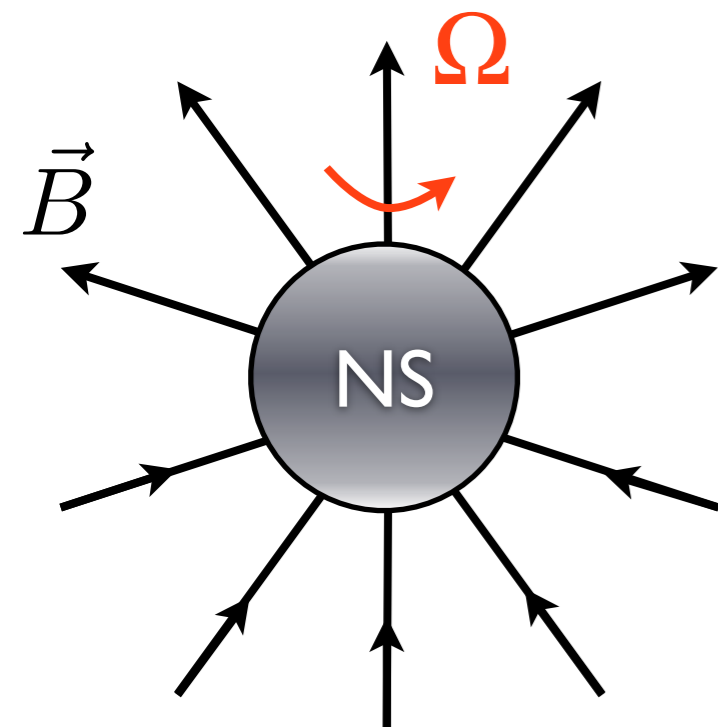


How do ^{force-free} Jets Accelerate?

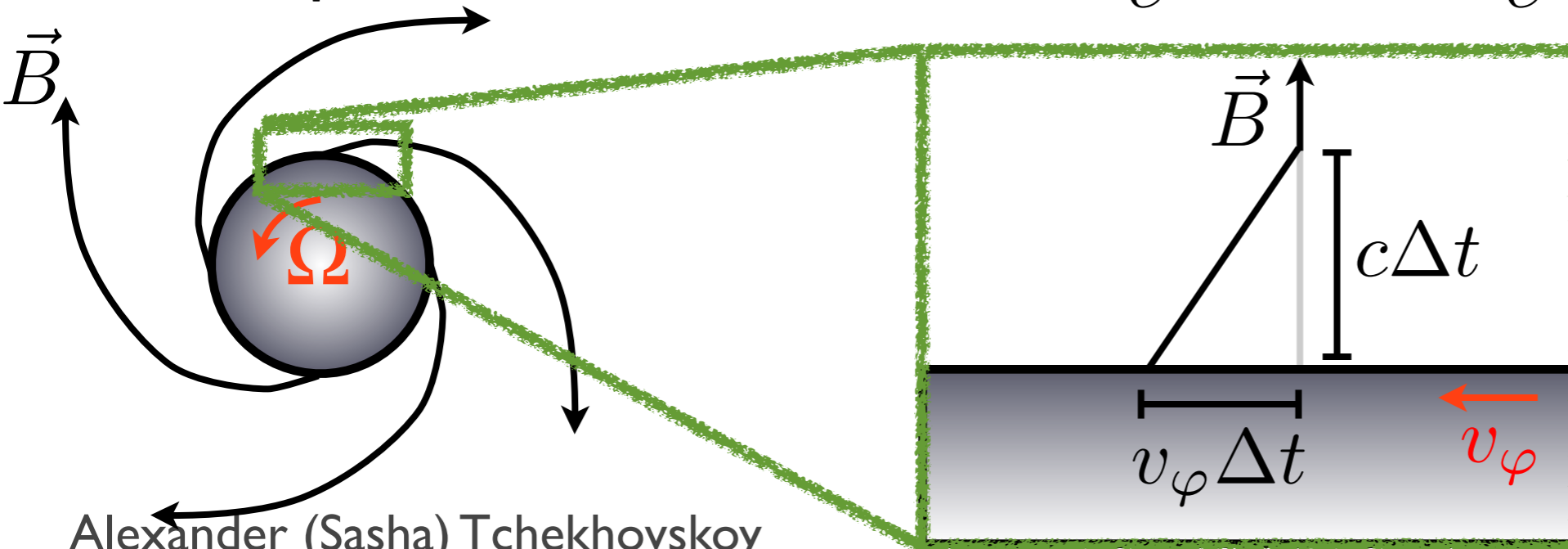


Assume the jets are **massless (force-free)** for simplicity

How do ^{force-free} Jets Accelerate?

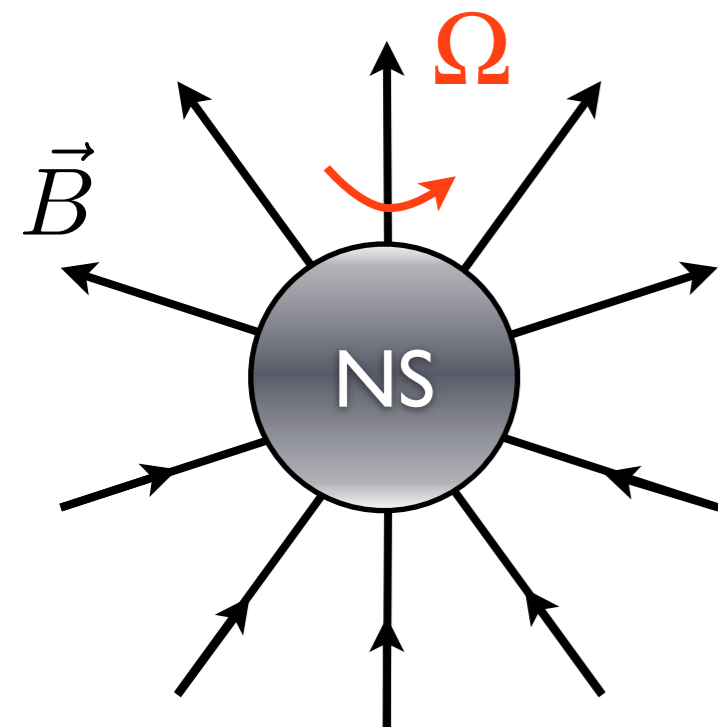


$$B_\varphi = -\frac{v_\varphi}{c} B_r = -\frac{\Omega R}{c} B_r$$



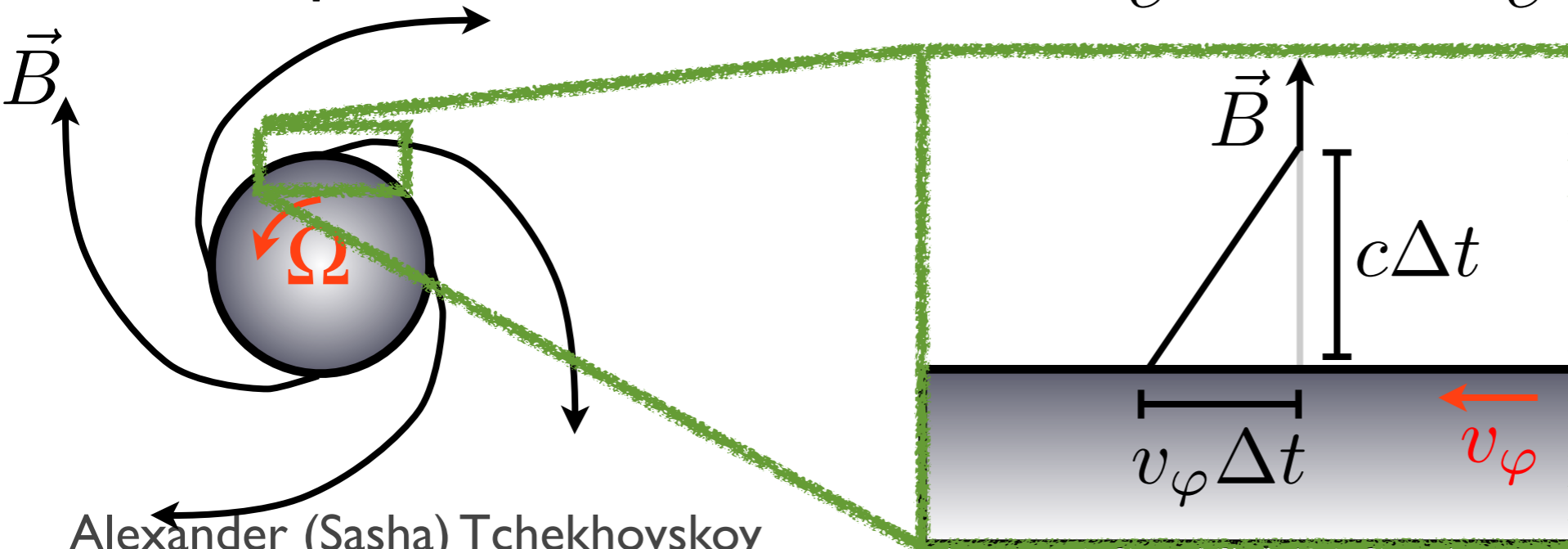
Assume the jets
are **massless**
(**force-free**)
for simplicity

How do ^{force-free} Jets Accelerate?



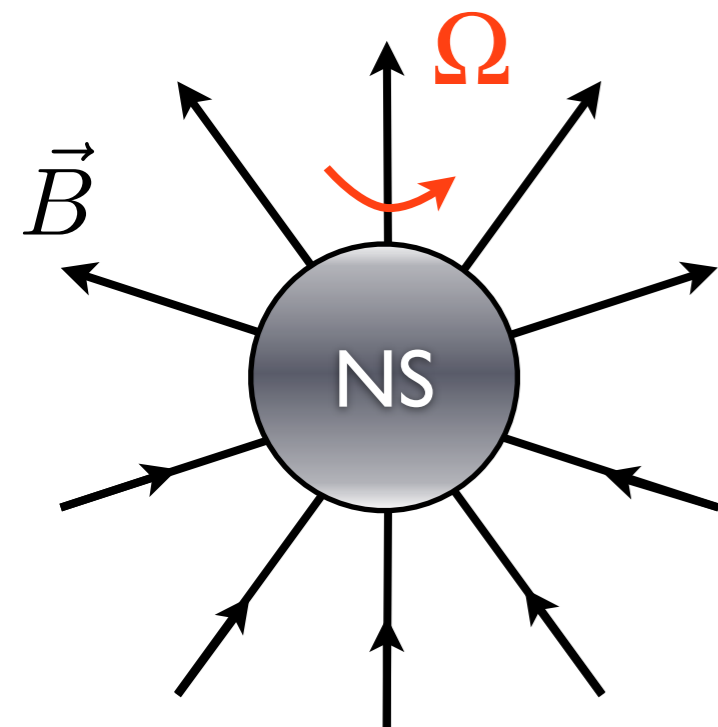
$$E = \left| -\frac{\vec{v}}{c} \times \vec{B} \right| = +\frac{\Omega R}{c} B_r$$

$$B_\varphi = -\frac{v_\varphi}{c} B_r = -\frac{\Omega R}{c} B_r$$



Assume the jets are **massless (force-free)** for simplicity

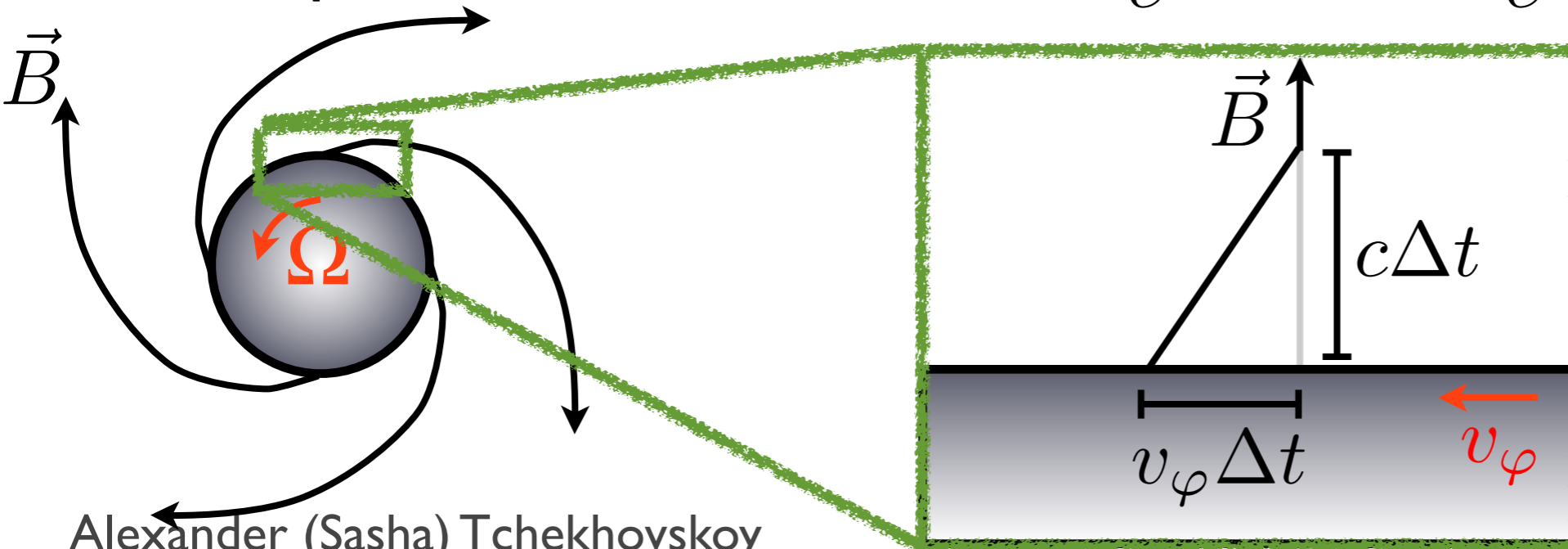
How do ^{force-free} Jets Accelerate?



$$\frac{v}{c} = \left| \frac{\vec{E} \times \vec{B}}{B^2} \right| = \frac{E}{B}$$

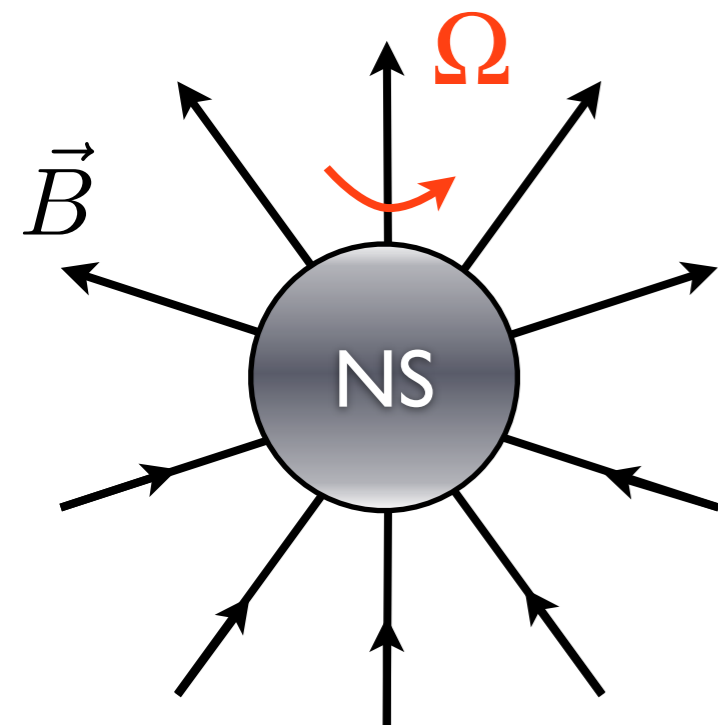
$$E = \left| -\frac{\vec{v}}{c} \times \vec{B} \right| = +\frac{\Omega R}{c} B_r$$

$$B_\varphi = -\frac{v_\varphi}{c} B_r = -\frac{\Omega R}{c} B_r$$



Assume the jets are **massless (force-free)** for simplicity

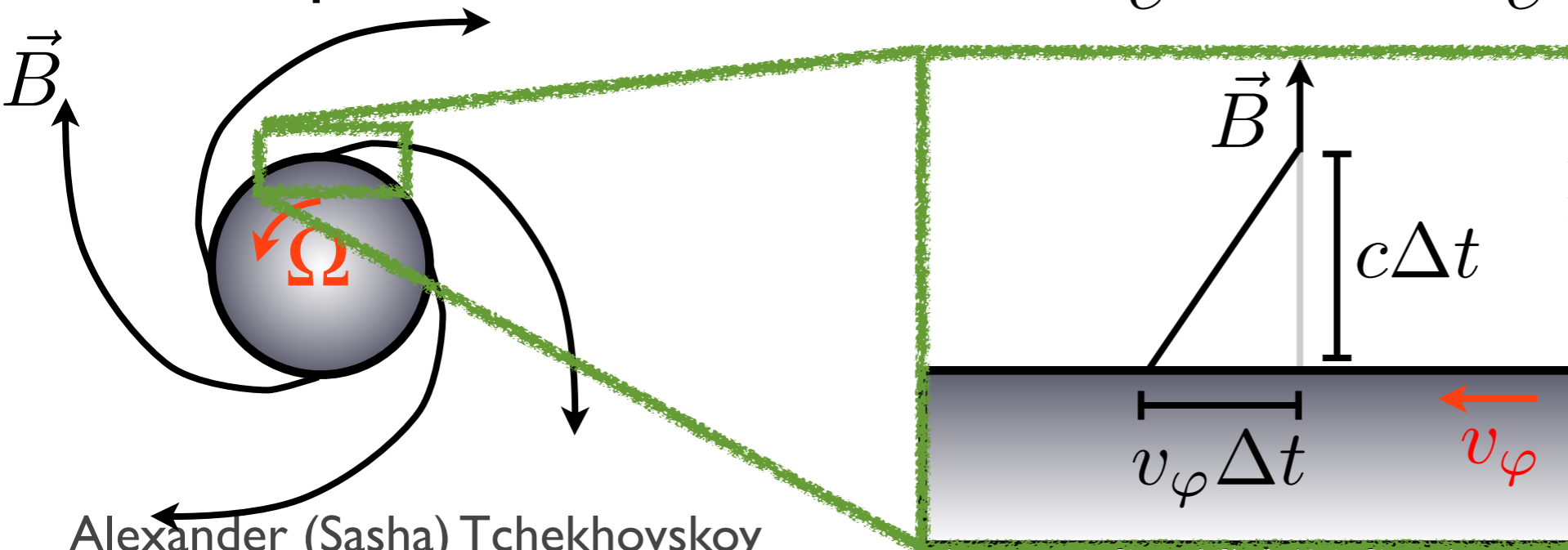
How do ^{force-free} Jets Accelerate?



$$\frac{v}{c} = \left| \frac{\vec{E} \times \vec{B}}{B^2} \right| = \frac{E}{B} \Rightarrow 1 - \frac{v^2}{c^2} = \frac{B^2 - E^2}{B^2}$$

$$E = \left| -\frac{\vec{v}}{c} \times \vec{B} \right| = +\frac{\Omega R}{c} B_r$$

$$B_\varphi = -\frac{v_\varphi}{c} B_r = -\frac{\Omega R}{c} B_r$$



Assume the jets
are **massless**
(**force-free**)
for simplicity

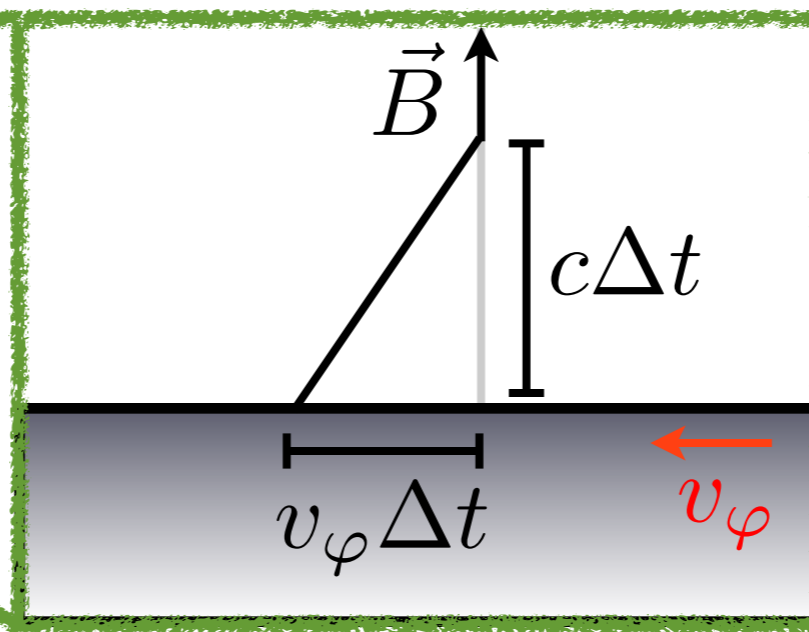
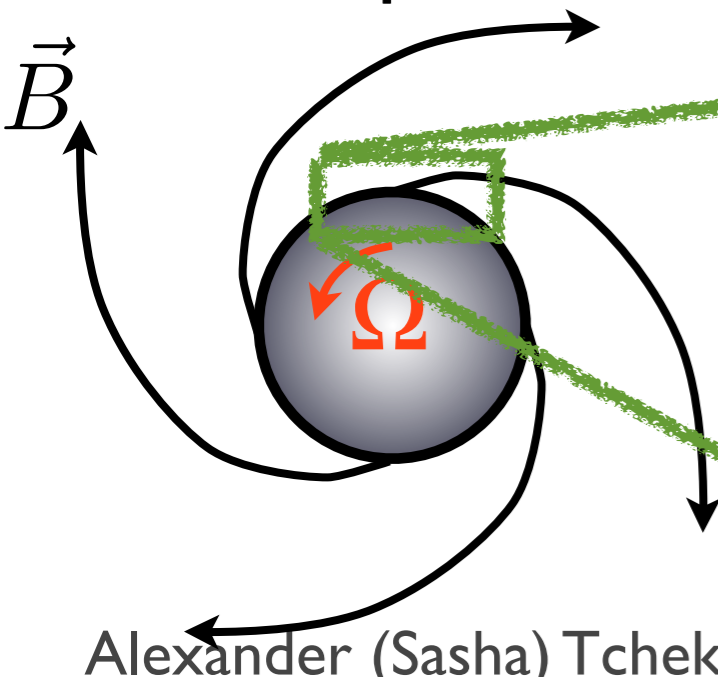
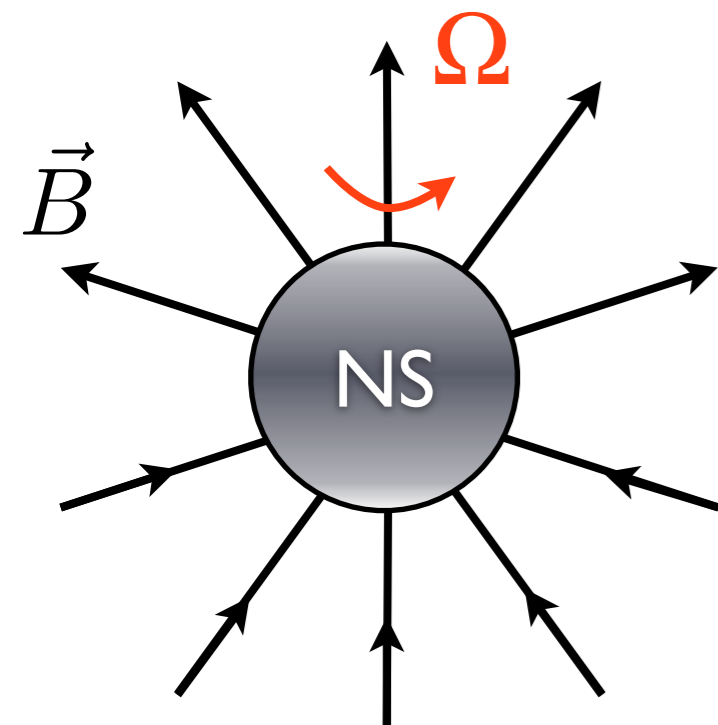
How do ^{force-free} Jets Accelerate?

$$\gamma^2 = \frac{1}{1 - v^2/c^2} = \frac{B_r^2 + B_\phi^2}{B_r^2 + \cancel{B_\phi^2} - E^2} = 1 + \frac{B_\phi^2}{B_r^2} = 1 + (\Omega R/c)^2$$

$$\frac{v}{c} = \left| \frac{\vec{E} \times \vec{B}}{B^2} \right| = \frac{E}{B} \Rightarrow 1 - \frac{v^2}{c^2} = \frac{B^2 - E^2}{B^2}$$

$$E = \left| -\frac{\vec{v}}{c} \times \vec{B} \right| = +\frac{\Omega R}{c} B_r$$

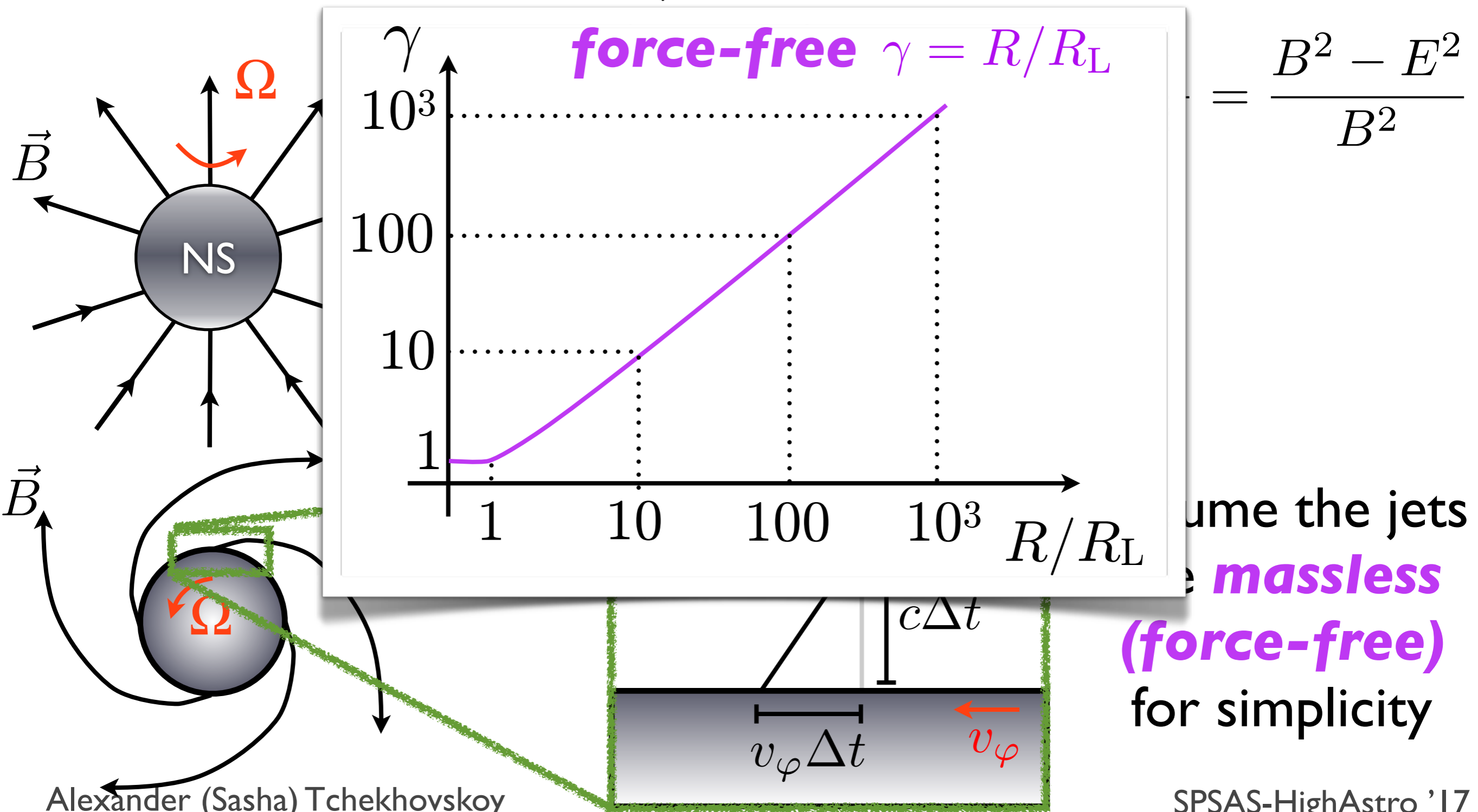
$$B_\phi = -\frac{v_\phi}{c} B_r = -\frac{\Omega R}{c} B_r$$



Assume the jets are **massless (force-free)** for simplicity

How do ^{force-free} Jets Accelerate?

$$\gamma^2 = \frac{1}{1 - v^2/c^2} = \frac{B_r^2 + B_\phi^2}{B_r^2 + \cancel{B_\phi^2} - E^2} = 1 + \frac{B_\phi^2}{B_r^2} = 1 + (\Omega R/c)^2$$



mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$F_B = B_p$$

$$F_M = \gamma \rho v_p$$

mass-loaded How do^vJets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\left. \begin{array}{l} F_B = B_p \\ F_M = \gamma \rho v_p \end{array} \right| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const}$$

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\left. \begin{array}{l} F_B = B_p \\ F_M = \gamma \rho v_p \end{array} \right| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const}$$

$$F_E = F_{EM} + F_{KE}$$

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\left. \begin{array}{l} F_B = B_p \\ F_M = \gamma \rho v_p \end{array} \right| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const}$$

$$F_E = F_{EM} + F_{KE}$$
$$\parallel$$
$$\frac{cEB_\varphi}{4\pi}$$

mass-loaded How do^vJets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\left. \begin{array}{l} F_B = B_p \\ F_M = \gamma \rho v_p \end{array} \right| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const}$$

$$F_E = F_{EM} + F_{KE}$$
$$\begin{array}{cc} \parallel & \parallel \\ \frac{cEB_\varphi}{4\pi} & \gamma F_M \end{array}$$

How do ^{mass-loaded} Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\begin{array}{l}
 F_B = B_p \\
 F_M = \gamma \rho v_p \\
 F_E = \underbrace{F_{EM}}_{\parallel} + \underbrace{F_{KE}}_{\parallel}
 \end{array}
 \left| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const} \right.$$

$$\left. \begin{array}{l}
 \Rightarrow \mu = \frac{F_E}{F_M} = \gamma \frac{F_{EM}}{F_{KE}} + \gamma \\
 \frac{cEB_\varphi}{4\pi} \quad \gamma F_M
 \end{array} \right.$$

How do ^{mass-loaded} Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\begin{array}{l}
 F_B = B_p \\
 F_M = \gamma \rho v_p
 \end{array} \left| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const} \right.$$

$$\sigma = b^2 / 4\pi \rho c^2$$

$$\begin{array}{l}
 F_E = F_{EM} + F_{KE} \\
 \parallel \quad \parallel \\
 \frac{cEB_\varphi}{4\pi} \quad \gamma F_M
 \end{array} \left| \Rightarrow \mu = \frac{F_E}{F_M} = \gamma \frac{F_{EM}}{F_{KE}} + \gamma \right.$$

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\begin{array}{l}
 F_B = B_p \\
 F_M = \gamma \rho v_p
 \end{array} \left| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const} \right.$$

$$\sigma = b^2 / 4\pi \rho c^2$$

$$\begin{array}{l}
 F_E = F_{EM} + F_{KE} \\
 \parallel \qquad \parallel \\
 \frac{cEB_\varphi}{4\pi} \quad \gamma F_M
 \end{array} \left| \Rightarrow \mu = \frac{F_E}{F_M} = \gamma \frac{F_{EM}}{F_{KE}} + \gamma = \gamma(\sigma + 1) \right.$$

mass-loaded How do ^vJets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\left. \begin{array}{l} F_B = B_p \\ F_M = \gamma \rho v_p \end{array} \right| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const}$$
$$\left. \begin{array}{l} F_E = F_{EM} + F_{KE} \end{array} \right| \Rightarrow \mu = \frac{F_E}{F_M} = \gamma \frac{F_{EM}}{F_{KE}} + \gamma = \gamma(\sigma + 1)$$

$\sigma = b^2 / 4\pi \rho c^2$
//

μ sets the max Lorentz factor: $\gamma \leq \mu$

mass-loaded How do ^vJets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\left. \begin{array}{l} F_B = B_p \\ F_M = \gamma \rho v_p \end{array} \right| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const}$$

$$\left. \begin{array}{l} F_E = F_{EM} + F_{KE} \end{array} \right| \Rightarrow \mu = \frac{F_E}{F_M} = \gamma \frac{F_{EM}}{F_{KE}} + \gamma = \gamma(\sigma + 1)$$

$\sigma = b^2 / 4\pi \rho c^2$
//

μ sets the max Lorentz factor: $\gamma \leq \mu$

σ sets the speed of fast waves: $\gamma_F \equiv \sigma^{1/2}$

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\left. \begin{array}{l} F_B = B_p \\ F_M = \gamma \rho v_p \end{array} \right| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const}$$

$$\left. \begin{array}{l} F_E = F_{EM} + F_{KE} \end{array} \right| \Rightarrow \mu = \frac{F_E}{F_M} = \gamma \frac{F_{EM}}{F_{KE}} + \gamma = \gamma(\sigma + 1)$$

$\sigma = b^2 / 4\pi \rho c^2$
//

μ sets the max Lorentz factor: $\gamma \leq \mu$

σ sets the speed of fast waves: $\gamma_F \equiv \sigma^{1/2}$

In *force-free*, $\sigma = \infty$, and fast waves travel at light speed.

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\left. \begin{array}{l} F_B = B_p \\ F_M = \gamma \rho v_p \end{array} \right| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const}$$

$$\left. \begin{array}{l} F_E = F_{EM} + F_{KE} \end{array} \right| \Rightarrow \mu = \frac{F_E}{F_M} = \gamma \frac{F_{EM}}{F_{KE}} + \gamma = \gamma(\sigma + 1)$$

$\sigma = b^2 / 4\pi \rho c^2$
//

μ sets the max Lorentz factor: $\gamma \leq \mu$

σ sets the speed of fast waves: $\gamma_F \equiv \sigma^{1/2}$

In *force-free*, $\sigma = \infty$, and fast waves travel at light speed.

So, *force-free* breaks down when $\gamma = \gamma_F \equiv \sigma^{1/2}$

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\left. \begin{array}{l} F_B = B_p \\ F_M = \gamma \rho v_p \end{array} \right| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const}$$

$$\left. \begin{array}{l} F_E = F_{EM} + F_{KE} \end{array} \right| \Rightarrow \mu = \frac{F_E}{F_M} = \gamma \frac{F_{EM}}{F_{KE}} + \gamma = \gamma(\sigma + 1)$$

$\sigma = b^2 / 4\pi \rho c^2$
//

μ sets the max Lorentz factor: $\gamma \leq \mu$

σ sets the speed of fast waves: $\gamma_F \equiv \sigma^{1/2}$

In *force-free*, $\sigma = \infty$, and fast waves travel at light speed.

So, *force-free* breaks down when $\gamma = \gamma_F \equiv \sigma^{1/2} = (\mu/\gamma)^{1/2}$

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\left. \begin{array}{l} F_B = B_p \\ F_M = \gamma \rho v_p \end{array} \right| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const}$$

$$\sigma = b^2 / 4\pi \rho c^2$$

$$\left. \begin{array}{l} F_E = F_{EM} + F_{KE} \end{array} \right| \Rightarrow \mu = \frac{F_E}{F_M} = \gamma \frac{F_{EM}}{F_{KE}} + \gamma = \gamma(\sigma + 1)$$

μ sets the max Lorentz factor: $\gamma \leq \mu$

σ sets the speed of fast waves: $\gamma_F \equiv \sigma^{1/2}$

In *force-free*, $\sigma = \infty$, and fast waves travel at light speed.

So, *force-free* breaks down when $\gamma = \gamma_F \equiv \sigma^{1/2} = (\mu/\gamma)^{1/2}$
 $\gamma = \mu^{1/3} \ll \mu$

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\left. \begin{array}{l} F_B = B_p \\ F_M = \gamma \rho v_p \end{array} \right| \Rightarrow \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const}$$

$$\sigma = b^2 / 4\pi \rho c^2 //$$

$$\left. \begin{array}{l} F_E = F_{EM} + F_{KE} \end{array} \right| \Rightarrow \mu = \frac{F_E}{F_M} = \gamma \frac{F_{EM}}{F_{KE}} + \gamma = \gamma(\sigma + 1)$$

μ sets the max Lorentz factor: $\gamma \leq \mu$

σ sets the speed of fast waves: $\gamma_F \equiv \sigma^{1/2}$

In *force-free*, $\sigma = \infty$, and fast waves travel at light speed.

So, *force-free* breaks down when $\gamma = \gamma_F \equiv \sigma^{1/2} = (\mu/\gamma)^{1/2}$
 $10 = \gamma = \mu^{1/3} \ll \mu = 10^3$

How do ^{mass-loaded} jets Accelerate?

Conserved

$$F_B = B_p$$

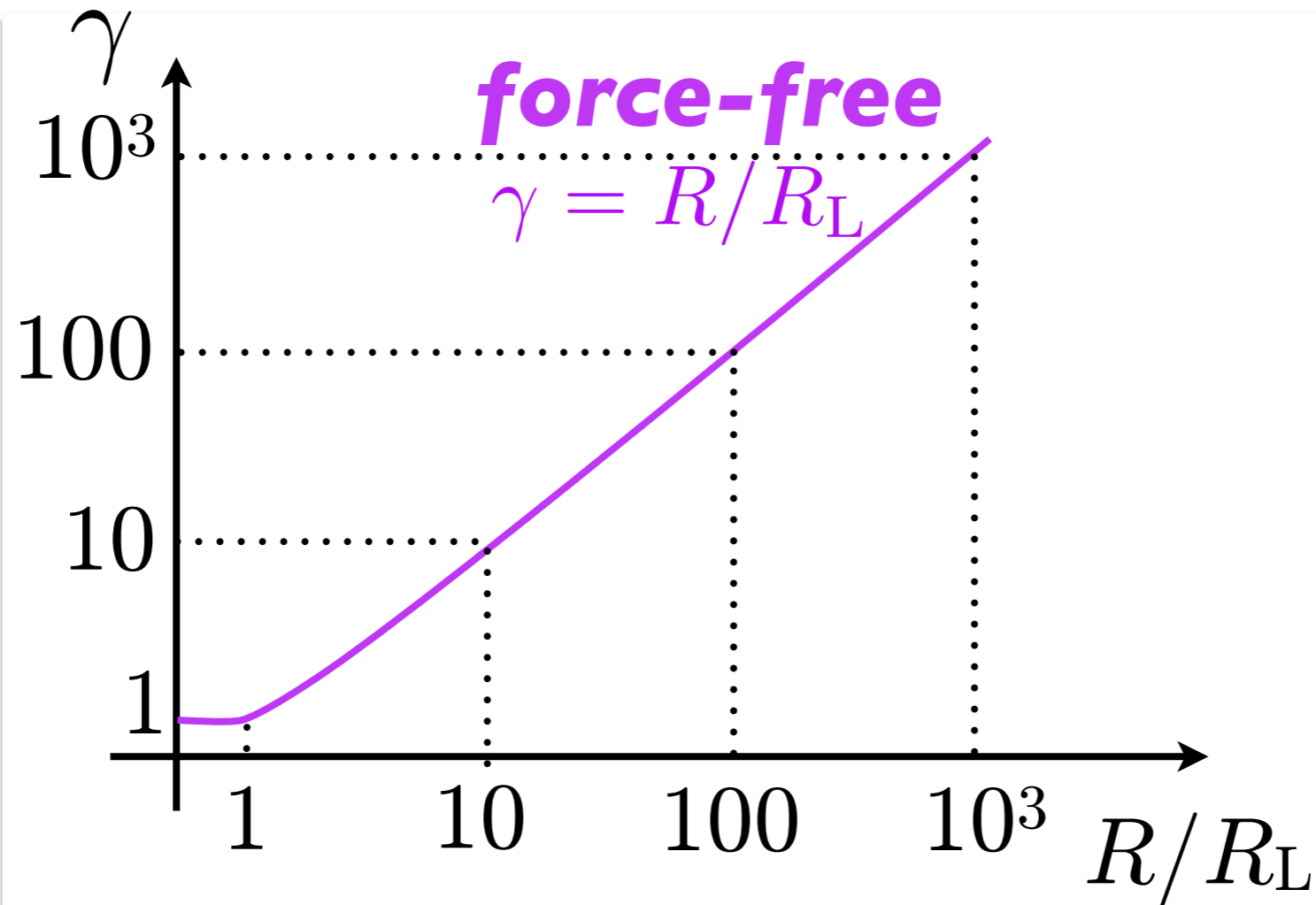
$$F_M = \gamma \rho v_p$$

$$F_E = F_{EM}$$

and fluxes:

$$\rho c^2$$

$$+ 1)$$



σ sets the speed of fast waves: $\gamma_F \equiv \sigma^{1/2}$

In **force-free**, $\sigma = \infty$, and fast waves travel at light speed.

So, **force-free** breaks down when $\gamma = \gamma_F \equiv \sigma^{1/2} = (\mu/\gamma)^{1/2}$
 $10 = \gamma = \mu^{1/3} \ll \mu = 10^3$

How do ^{mass-loaded} jets Accelerate?

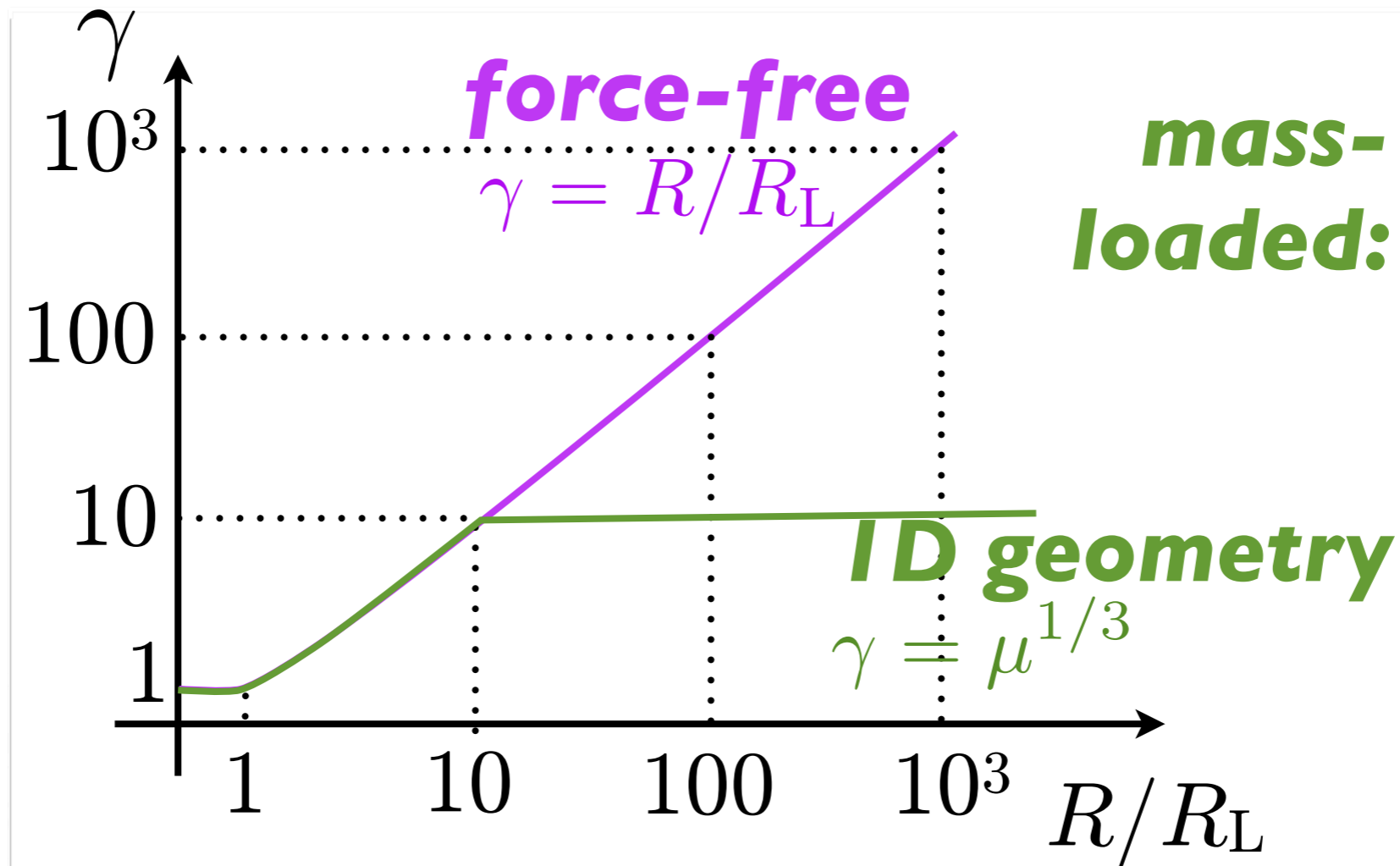
Conserved

$$F_B = B_p$$

$$F_M = \gamma \rho v_p$$

$$F_E = F_{EM}$$

d fluxes:



σ sets the speed of fast waves: $\gamma_F \equiv \sigma^{1/2}$

In **force-free**, $\sigma = \infty$, and fast waves travel at light speed.

So, **force-free** breaks down when $\gamma = \gamma_F \equiv \sigma^{1/2} = (\mu/\gamma)^{1/2}$
 $10 = \gamma = \mu^{1/3} \ll \mu = 10^3$

How do ^{mass-loaded} jets Accelerate?

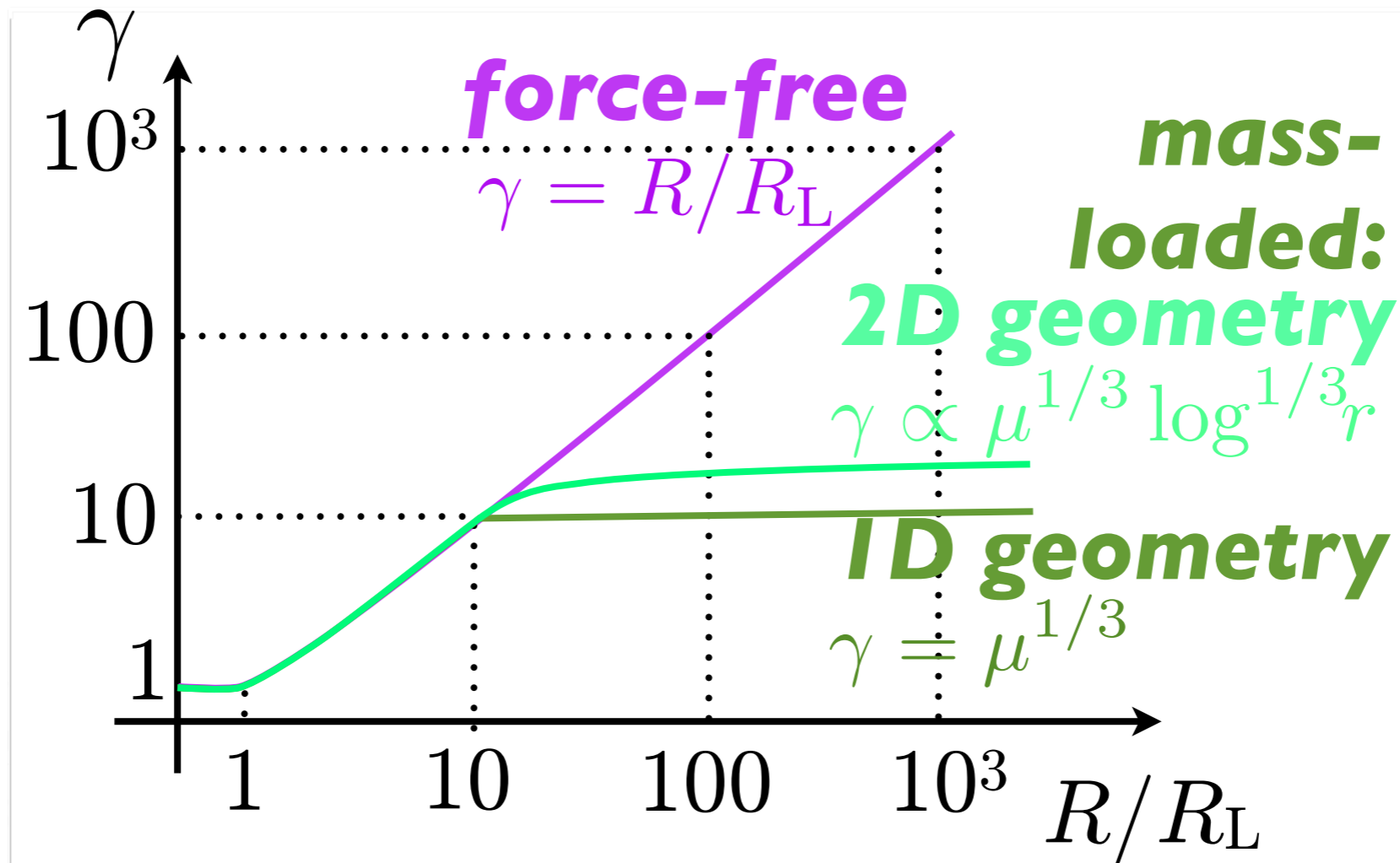
Conserved

$$F_B = B_p$$

$$F_M = \gamma \rho v_p$$

$$F_E = F_{EM}$$

d fluxes:



σ sets the speed of fast waves: $\gamma_F \equiv \sigma^{1/2}$

In **force-free**, $\sigma = \infty$, and fast waves travel at light speed.

So, **force-free** breaks down when $\gamma = \gamma_F \equiv \sigma^{1/2} = (\mu/\gamma)^{1/2}$
 $10 = \gamma = \mu^{1/3} \ll \mu = 10^3$

How do ^{mass-loaded} jets Accelerate? **Badly!**

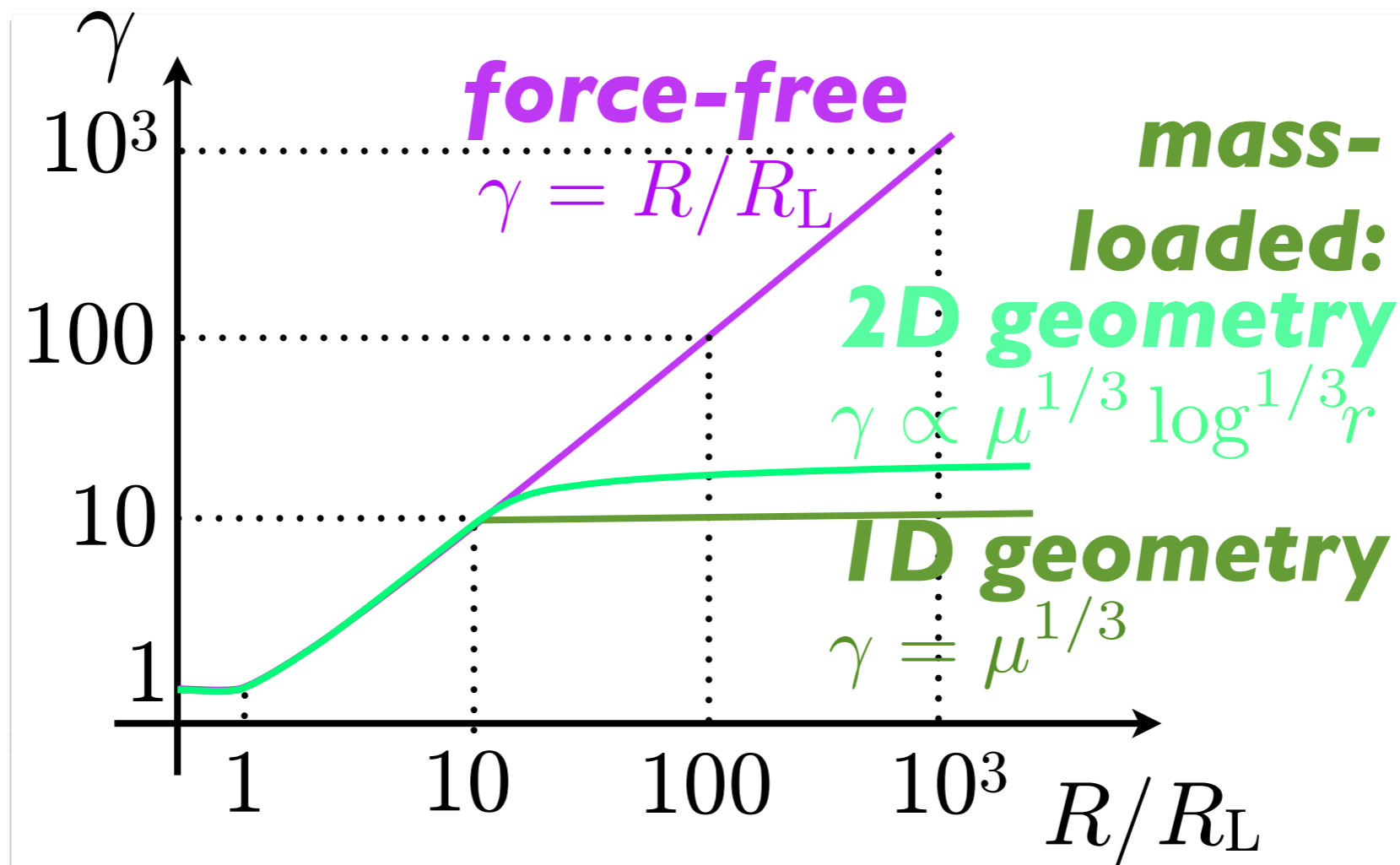
Conserved

$$F_B = B_p$$

$$F_M = \gamma \rho v_p$$

$$F_E = F_{EM}$$

d fluxes:



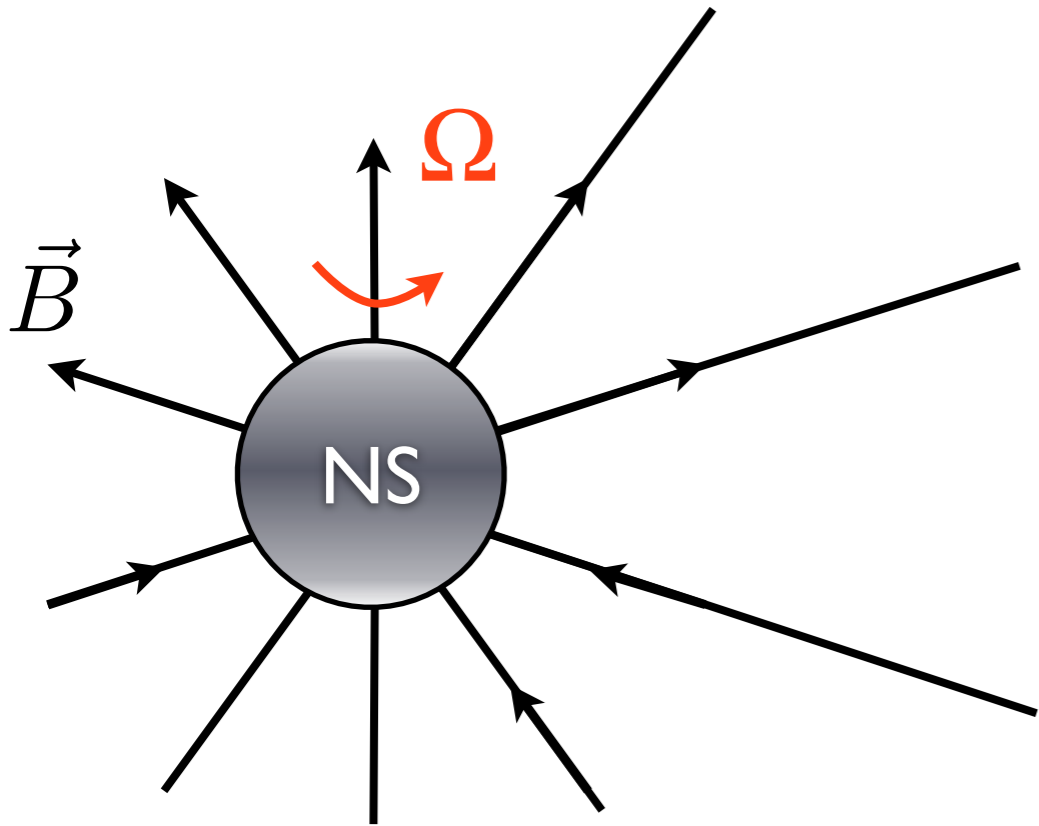
σ sets the speed of fast waves: $\gamma_F \equiv \sigma^{1/2}$

In **force-free**, $\sigma = \infty$, and fast waves travel at light speed.

So, **force-free** breaks down when $\gamma = \gamma_F \equiv \sigma^{1/2} = (\mu/\gamma)^{1/2}$
 $10 = \gamma = \mu^{1/3} \ll \mu = 10^3$

Why So **Bad** (1/2)?

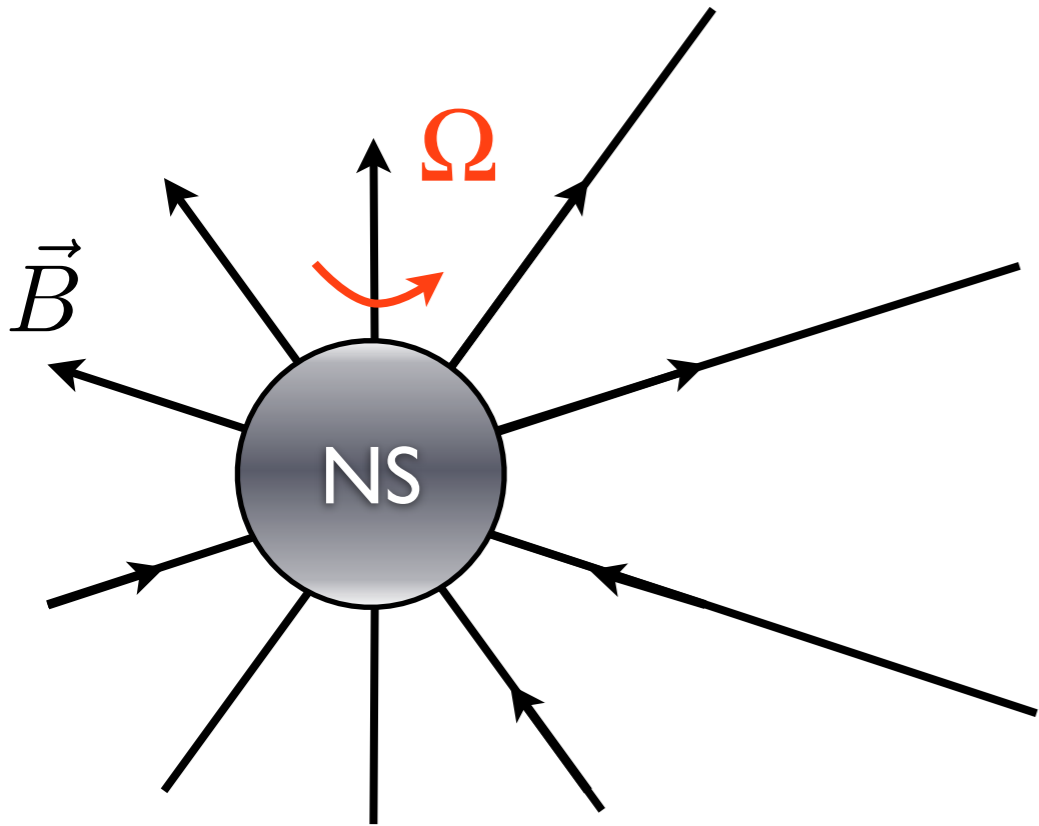
$$\gamma^2 = \frac{1}{1 - v^2/c^2} = \frac{B_r^2 + B_\varphi^2}{B_r^2 + \cancel{B_\varphi^2} - E^2} = 1 + \frac{B_\varphi^2}{B_r^2} = 1 + (\Omega R/c)^2$$



Why So **Bad** (1/2)?

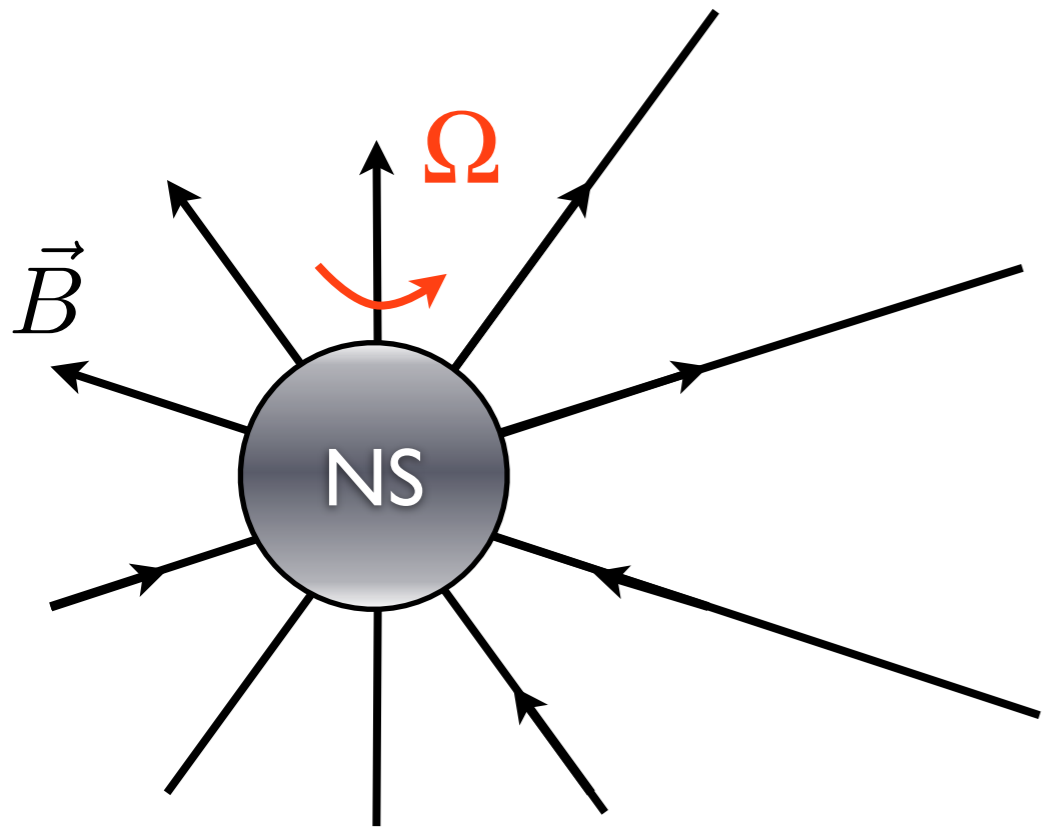
$$\gamma^2 = \frac{1}{1 - v^2/c^2} = \frac{B_r^2 + B_\varphi^2}{B_r^2 + \cancel{B_\varphi^2 - E^2}} = 1 + \frac{B_\varphi^2}{B_r^2} = 1 + (\Omega R/c)^2$$

Our key assumption: $B_\varphi^2 - E^2 \ll B_r^2$



Why So **Bad** (1/2)?

$$\gamma^2 = \frac{1}{1 - v^2/c^2} = \frac{B_r^2 + B_\varphi^2}{B_r^2 + \cancel{B_\varphi^2 - E^2}} = 1 + \frac{B_\varphi^2}{B_r^2} = 1 + (\Omega R/c)^2$$

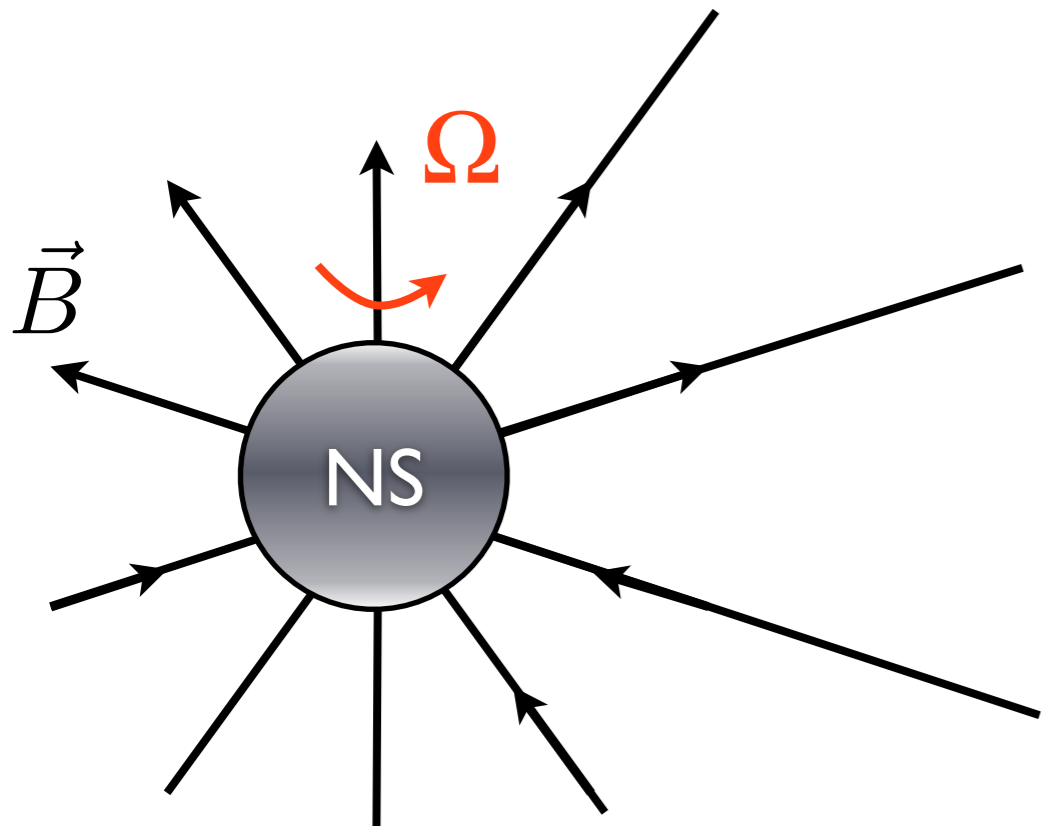


Our key assumption: $B_\varphi^2 - E^2 \ll B_r^2$

This means $B_r(\theta) = \text{constant}$ is the solution, $p_m = B_r^2 + \cancel{B_\varphi^2 - E^2} = B_r^2$

Why So **Bad** (1/2)?

$$\gamma^2 = \frac{1}{1 - v^2/c^2} = \frac{B_r^2 + B_\varphi^2}{B_r^2 + \cancel{B_\varphi^2 - E^2}} = 1 + \frac{B_\varphi^2}{B_r^2} = 1 + (\Omega R/c)^2$$



Our key assumption: $B_\varphi^2 - E^2 \ll B_r^2$

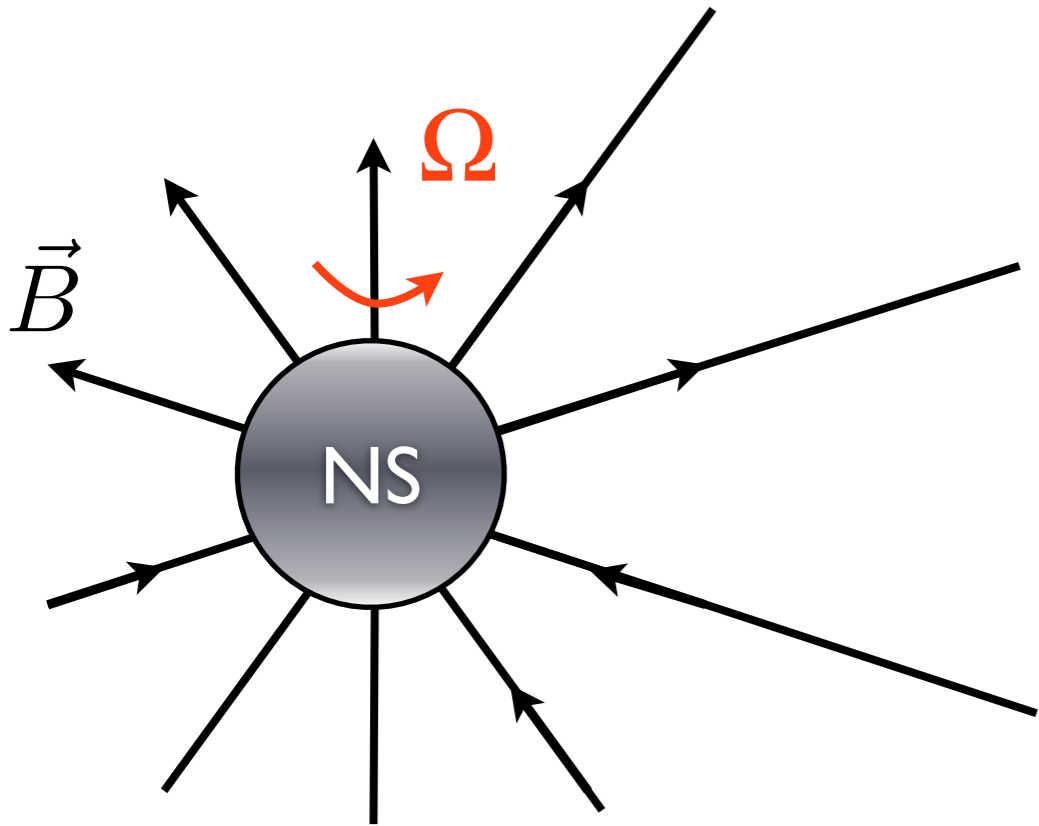
This means $B_r(\theta) = \text{constant}$ is the solution, $p_m = B_r^2 + \cancel{B_\varphi^2 - E^2} = B_r^2$

The opposite limit is what we need:

$$\gamma^2 = \frac{\cancel{B_r^2} + B_\varphi^2}{\cancel{B_r^2} + B_\varphi^2 - E^2} = \frac{B_\varphi^2}{B_\varphi^2 - E^2}$$

Why So **Bad** (1/2)?

$$\gamma^2 = \frac{1}{1 - v^2/c^2} = \frac{B_r^2 + B_\varphi^2}{B_r^2 + \cancel{B_\varphi^2 - E^2}} = 1 + \frac{B_\varphi^2}{B_r^2} = 1 + (\Omega R/c)^2$$



Our key assumption: $B_\varphi^2 - E^2 \ll B_r^2$

This means $B_r(\theta) = \text{constant}$ is the solution, $p_m = B_r^2 + \cancel{B_\varphi^2 - E^2} = B_r^2$

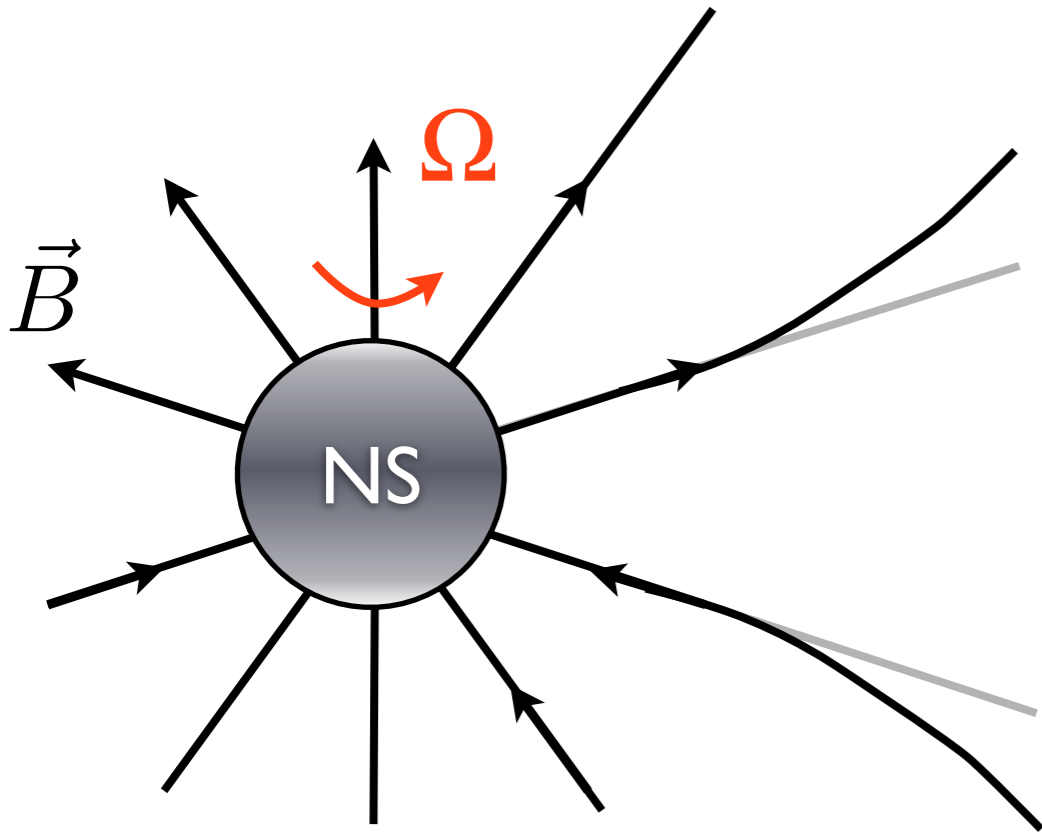
The opposite limit is what we need:

$$\gamma^2 = \frac{\cancel{B_r^2} + B_\varphi^2}{\cancel{B_r^2} + B_\varphi^2 - E^2} = \frac{B_\varphi^2}{B_\varphi^2 - E^2}$$

Because $B_\varphi^2 - E^2 \gg B_r^2$, toroidal magnetic pressure contribution breaks force balance, and magnetic field lines get **bent**. *Fast jets cannot make sharp turns.*

Why So **Bad** (1/2)?

$$\gamma^2 = \frac{1}{1 - v^2/c^2} = \frac{B_r^2 + B_\varphi^2}{B_r^2 + \cancel{B_\varphi^2} - E^2} = 1 + \frac{B_\varphi^2}{B_r^2} = 1 + (\Omega R/c)^2$$



Our key assumption: $B_\varphi^2 - E^2 \ll B_r^2$

This means $B_r(\theta) = \text{constant}$ is the solution, $p_m = B_r^2 + \cancel{B_\varphi^2} - E^2 = B_r^2$

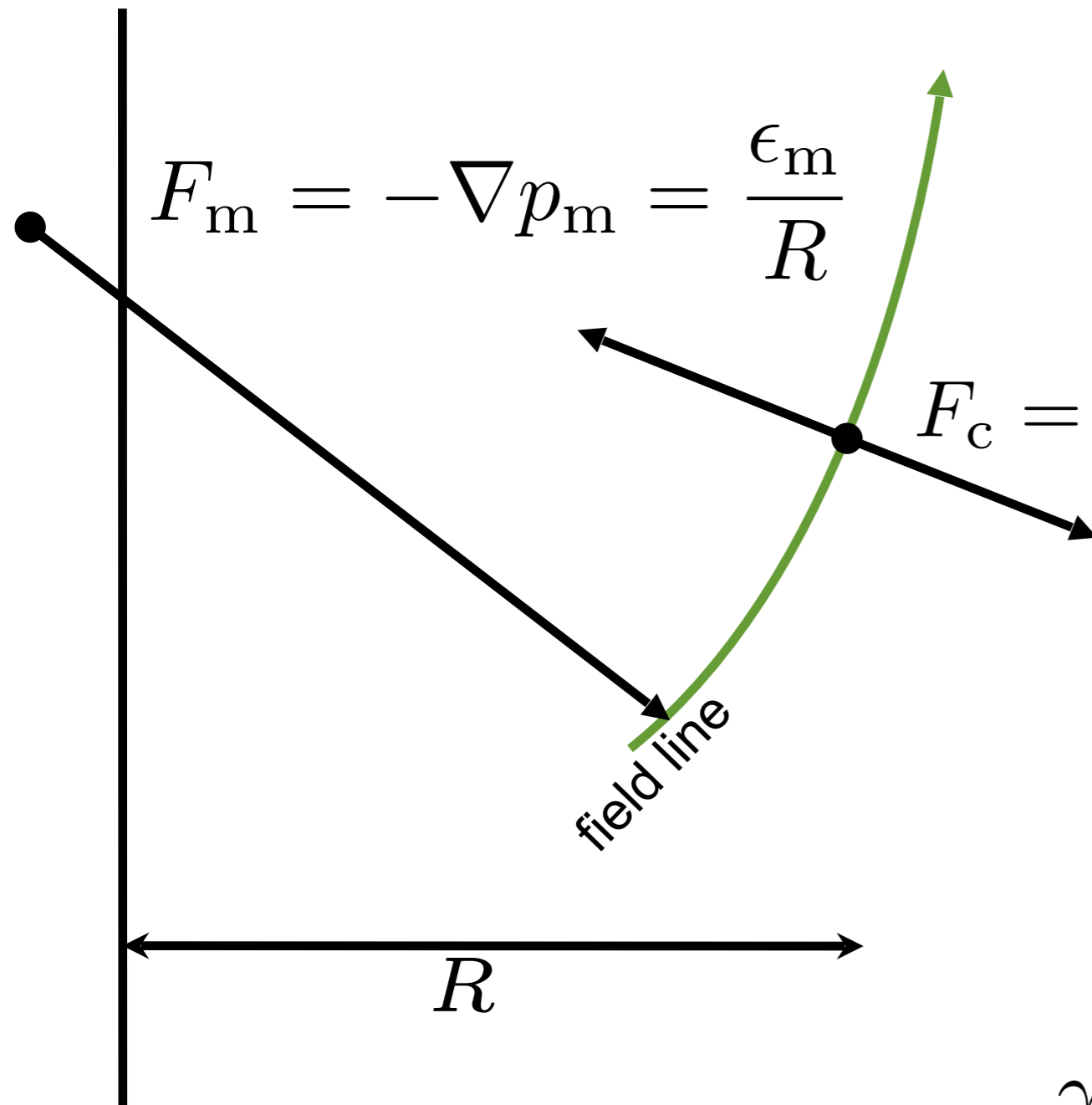
The opposite limit is what we need:

$$\gamma^2 = \frac{\cancel{B_r^2} + B_\varphi^2}{\cancel{B_r^2} + B_\varphi^2 - E^2} = \frac{B_\varphi^2}{B_\varphi^2 - E^2}$$

Because $B_\varphi^2 - E^2 \gg B_r^2$, toroidal magnetic pressure contribution breaks force balance, and magnetic field lines get ***bent***. *Fast jets cannot make sharp turns.*

Why So **Bad** (2/2)?

Force-balance across **bent** magnetic field lines, $B_\varphi^2 - E^2 \gg B_r^2$



$$F_m = -\nabla p_m = \frac{\epsilon_m}{R}$$

$$F_c = \frac{\epsilon_m \gamma^2}{R_c}$$

$$F_m = F_c$$

$$\frac{\epsilon_m}{R} = \frac{\epsilon_m \gamma^2}{R_c}$$

$$\gamma = \left(\frac{R_c}{R} \right)^{1/2}$$

Can combine both limits:

$$\frac{1}{\gamma^2} = \frac{1}{\gamma_1^2} + \frac{1}{\gamma_2^2}$$

$$\gamma_1 \approx \frac{\Omega R}{c}$$

$$\gamma_2 \approx \left(\frac{R_c}{R} \right)^{1/2}$$

$$B_\varphi^2 - E^2 \ll B_r^2$$

$$B_\varphi^2 - E^2 \gg B_r^2$$

How do ^{mass-loaded} Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\begin{array}{l}
 F_B = B_p \\
 F_M = \gamma \rho v_p = \eta B_p \\
 F_E = F_{EM} + F_{KE} \\
 \quad \parallel \quad \parallel \\
 \quad \frac{cEB_\varphi}{4\pi} \quad \gamma F_M
 \end{array}
 \left| \Rightarrow \begin{array}{l}
 \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const} \\
 \mu = \frac{F_E}{F_M} = \gamma \frac{F_{EM}}{F_{KE}} + \gamma = \gamma(\sigma + 1)
 \end{array} \right.$$

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\left. \begin{aligned} F_B &= B_p \\ F_M &= \gamma \rho v_p = \eta B_p \\ F_E &= F_{EM} + F_{KE} \\ &\quad \parallel \quad \parallel \\ &\quad \frac{cEB_\varphi}{4\pi} \quad \gamma F_M \end{aligned} \right| \Rightarrow \begin{aligned} \eta &= \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const} \\ \mu &= \frac{F_E}{F_M} \end{aligned}$$

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\begin{array}{l}
 F_B = B_p \\
 F_M = \gamma \rho v_p = \eta B_p \\
 F_E = F_{EM} + F_{KE} \\
 \quad \parallel \quad \parallel \\
 \quad \frac{cEB_\varphi}{4\pi} \quad \gamma F_M \\
 \quad \parallel E = B_\varphi = \Omega R B_p / c \\
 \quad \frac{\Omega^2 R^2 B_p^2}{4\pi c}
 \end{array}
 \left| \Rightarrow \begin{array}{l}
 \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const} \\
 \mu = \frac{F_E}{F_M}
 \end{array} \right.$$

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\begin{array}{l}
 F_B = B_p \\
 F_M = \gamma \rho v_p = \eta B_p \\
 F_E = F_{EM} + F_{KE} \\
 \quad \parallel \quad \parallel \\
 \quad \frac{cEB_\varphi}{4\pi} \quad \gamma F_M \\
 \quad \parallel E = B_\varphi = \Omega R B_p / c \\
 \quad \frac{\Omega^2 R^2 B_p^2}{4\pi c}
 \end{array}
 \left| \Rightarrow \right.
 \begin{array}{l}
 \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const} \\
 \mu = \frac{F_E}{F_M} \\
 \parallel \\
 \frac{\Omega^2}{4\pi^2 \eta c} \pi B_p R^2 + \gamma
 \end{array}$$

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\begin{array}{l}
 F_B = B_p \\
 F_M = \gamma \rho v_p = \eta B_p \\
 F_E = F_{EM} + F_{KE} \\
 \quad \parallel \quad \parallel \\
 \quad \frac{cEB_\varphi}{4\pi} \quad \gamma F_M \\
 \quad \parallel E = B_\varphi = \Omega R B_p / c \\
 \quad \frac{\Omega^2 R^2 B_p^2}{4\pi c}
 \end{array}
 \Rightarrow
 \begin{array}{l}
 \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const} \\
 \mu = \frac{F_E}{F_M} \\
 \parallel \\
 \frac{\Omega^2}{4\pi^2 \eta c} \pi B_p R^2 + \gamma = \frac{\mu}{\Phi} \pi B_p R^2 + \gamma
 \end{array}$$

$\Phi \equiv (\pi B_p R^2)_F$

mass-loaded How do^v Jets Accelerate?

Conserved quantities along jets = ratios of conserved fluxes:

$$\begin{array}{l}
 F_B = B_p \\
 F_M = \gamma \rho v_p = \eta B_p \\
 F_E = F_{EM} + F_{KE} \\
 \quad \parallel \quad \parallel \\
 \quad \frac{cEB_\varphi}{4\pi} \quad \gamma F_M
 \end{array}
 \Rightarrow
 \begin{array}{l}
 \eta = \frac{F_M}{F_B} = \frac{\gamma \rho v_p}{B_p} = \text{const} \\
 \mu = \frac{F_E}{F_M} \\
 \parallel \\
 \frac{\Omega^2}{4\pi^2 \eta c} \pi B_p R^2 + \gamma = \frac{\mu}{\Phi} \pi B_p R^2 + \gamma
 \end{array}$$

$$\Phi \equiv (\pi B_p R^2)_F$$

$$\frac{\gamma}{\mu} = 1 - \frac{\pi B_p R^2}{\Phi}$$

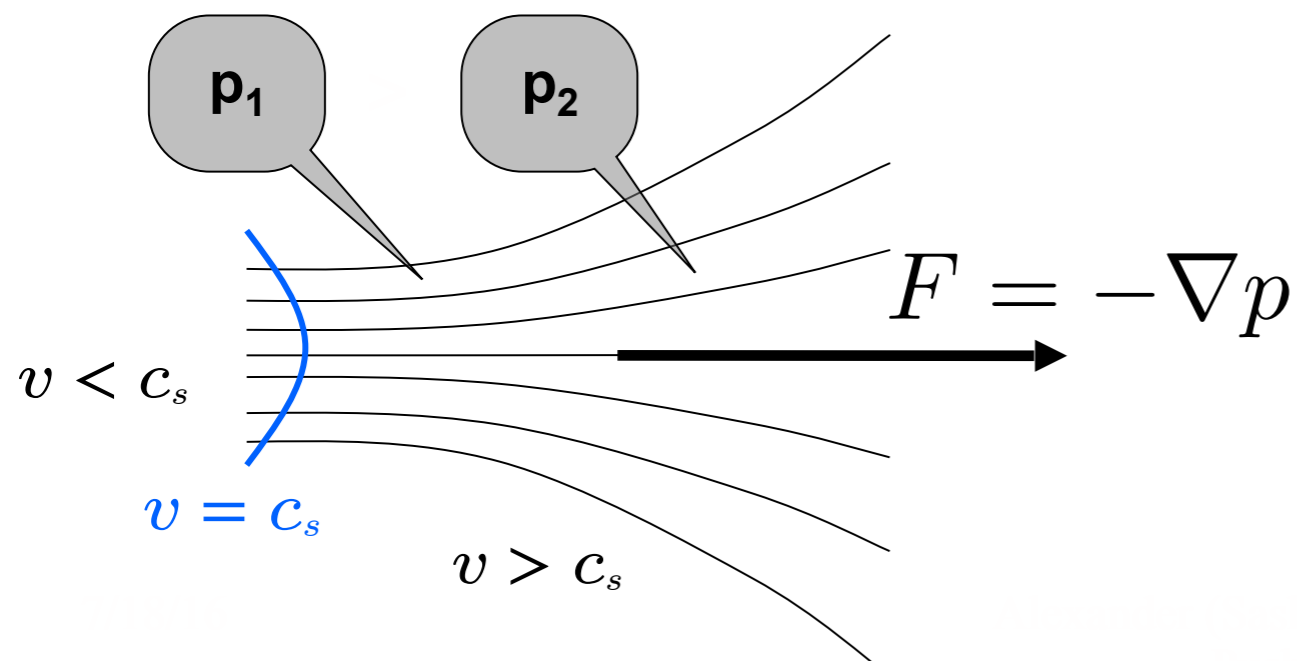
In order to accelerate efficiently, need reduction in local field line density (Komissarov+09, AT+09)

Acceleration in a magnetic nozzle

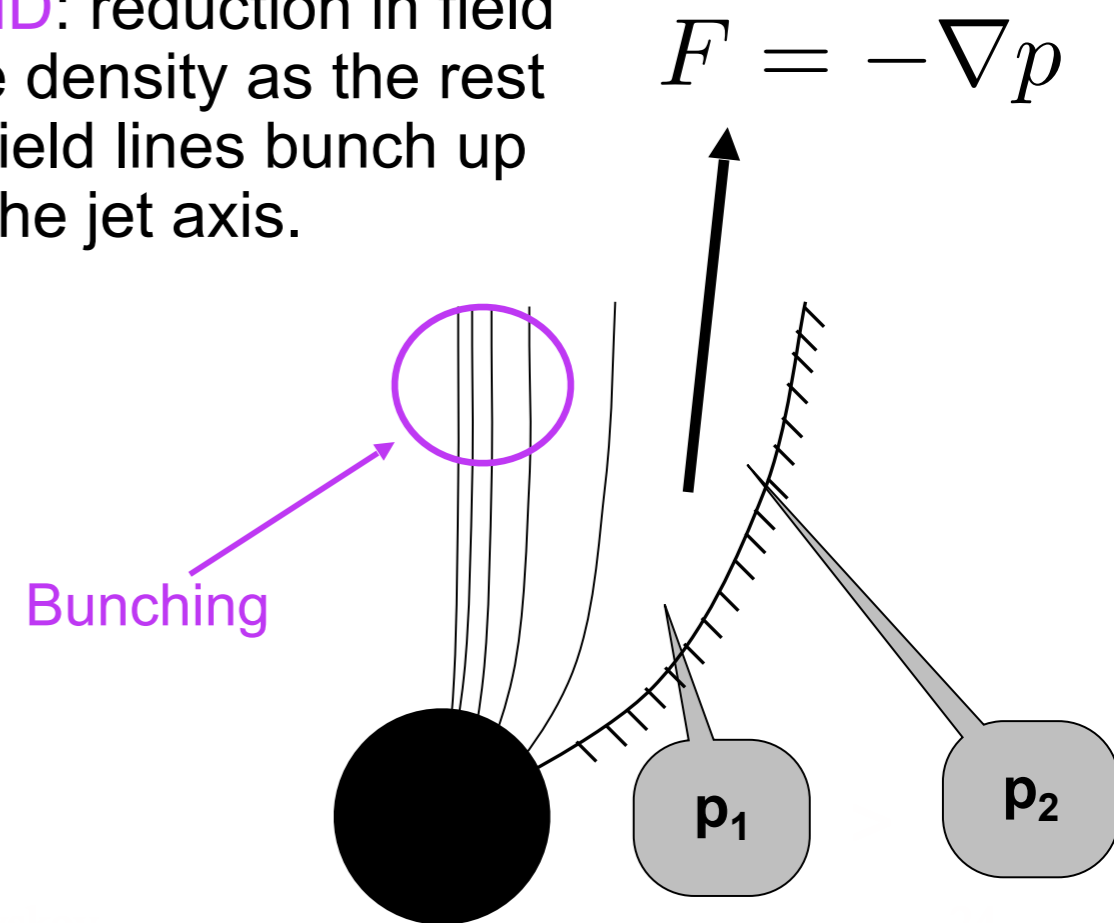
$$\frac{\gamma}{\mu} = 1 - \frac{\pi B_p R^2}{\Phi}$$

If $B_P(R) = \text{const}$, no acceleration.
Need magnetic flux bunching toward jet axis.

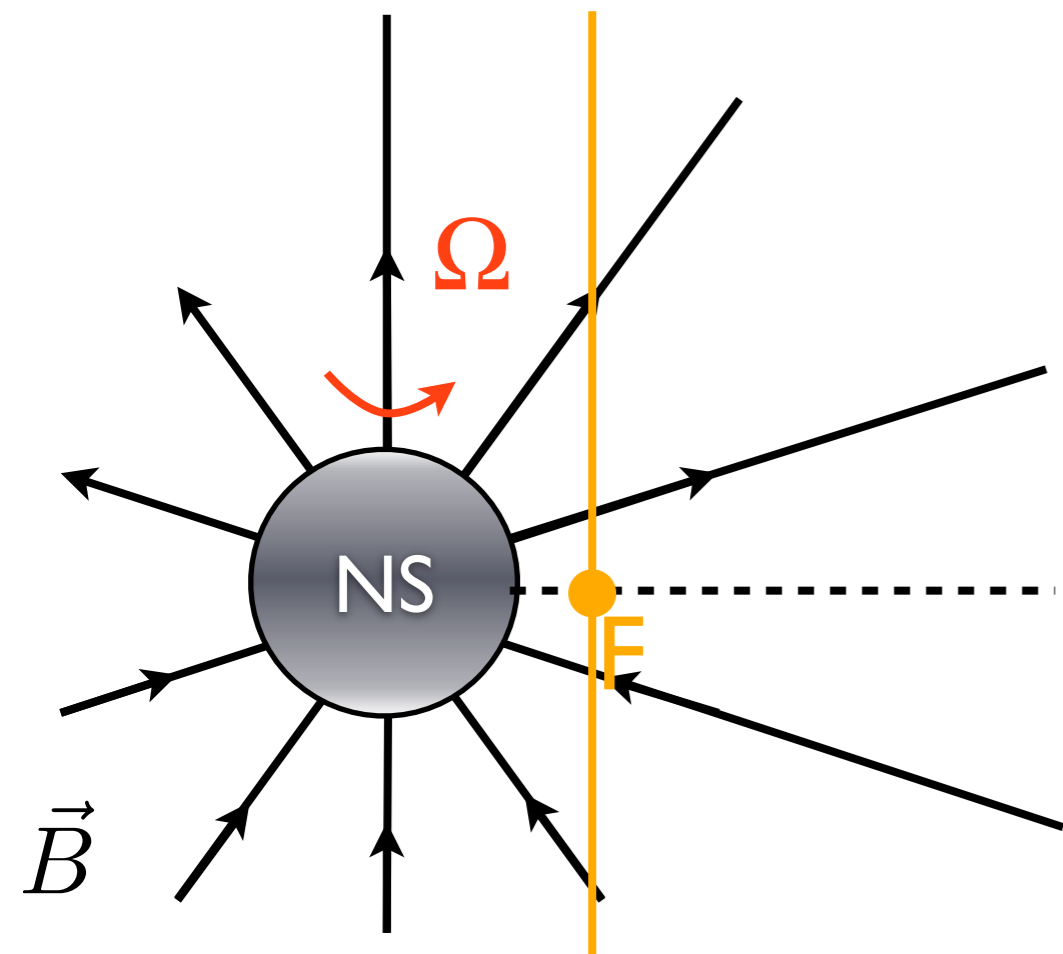
Hydro: de Laval nozzle: flow opens up after sonic surface \rightarrow pressure drops $\rightarrow \nabla p$ accelerates flow:



MHD: reduction in field line density as the rest of field lines bunch up at the jet axis.

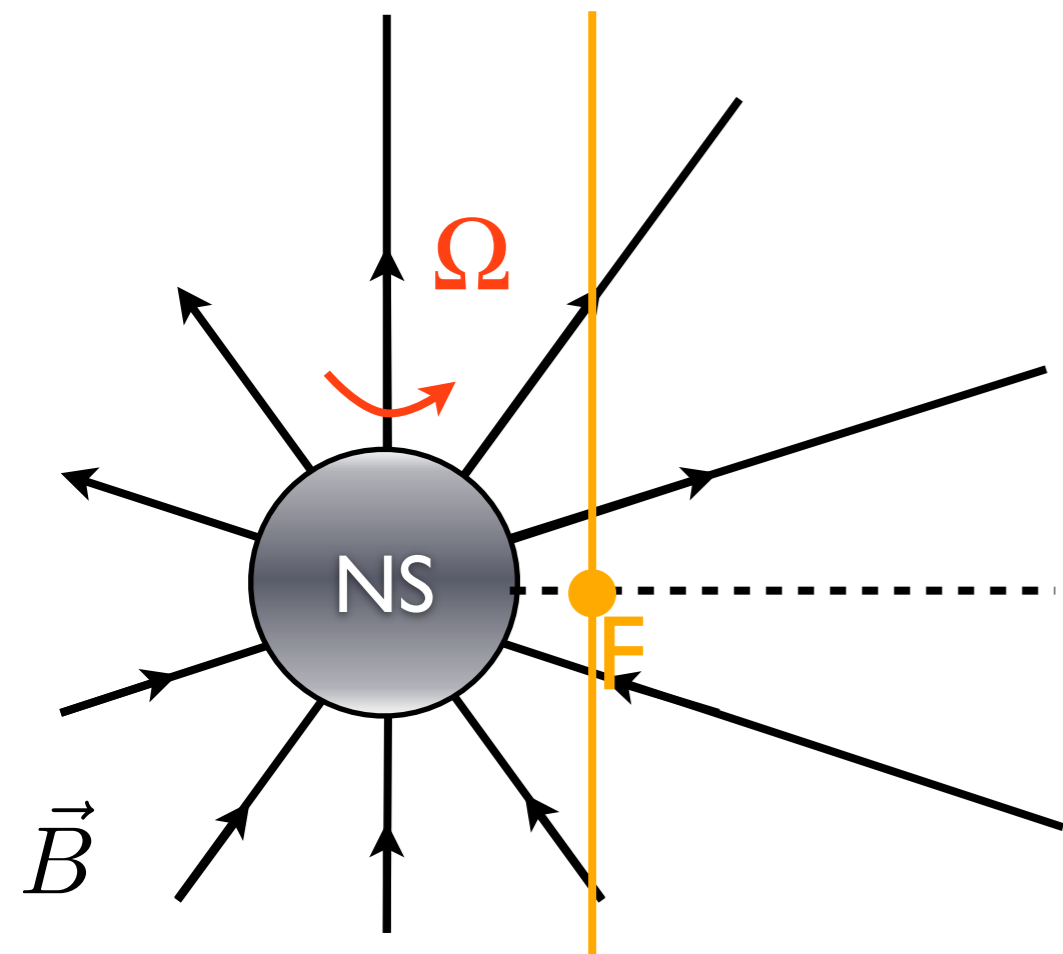


When Can Jets Accelerate?



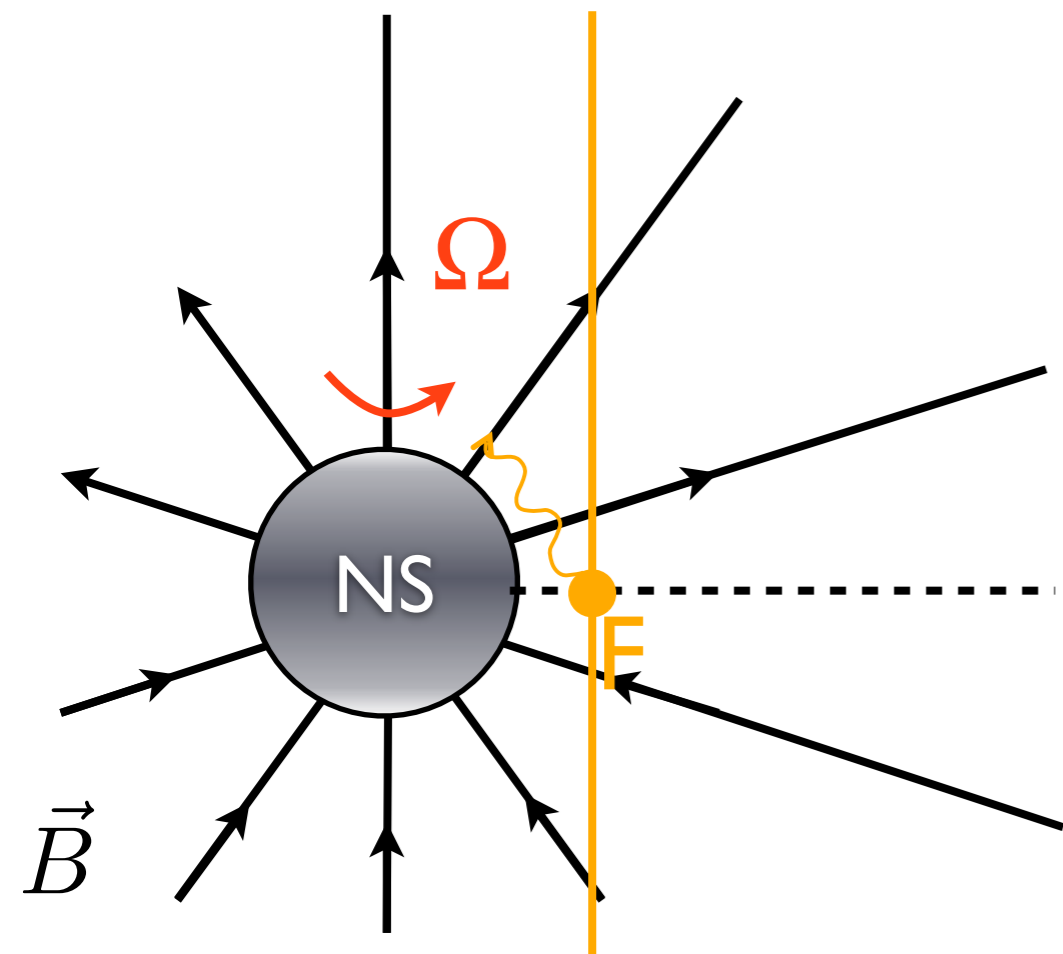
When Can Jets Accelerate?

- Communication is essential

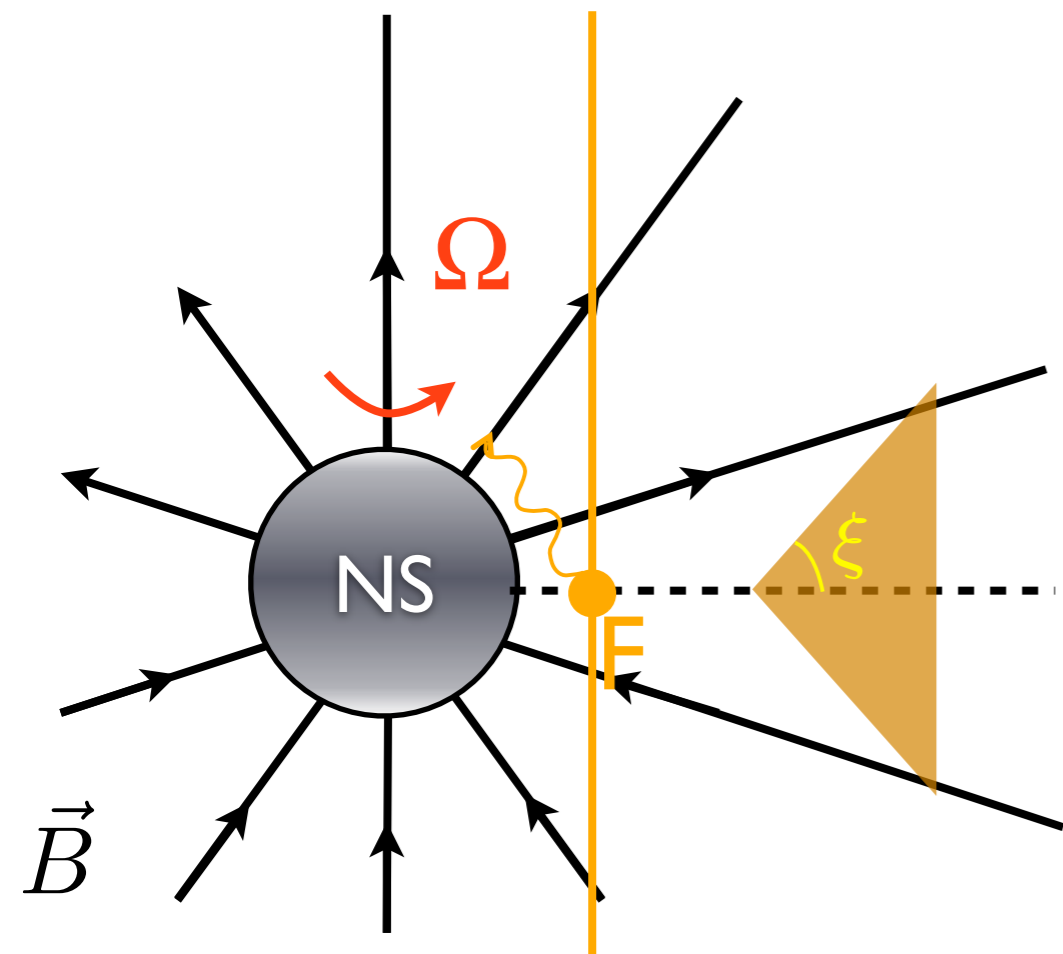


When Can Jets Accelerate?

- Communication is essential



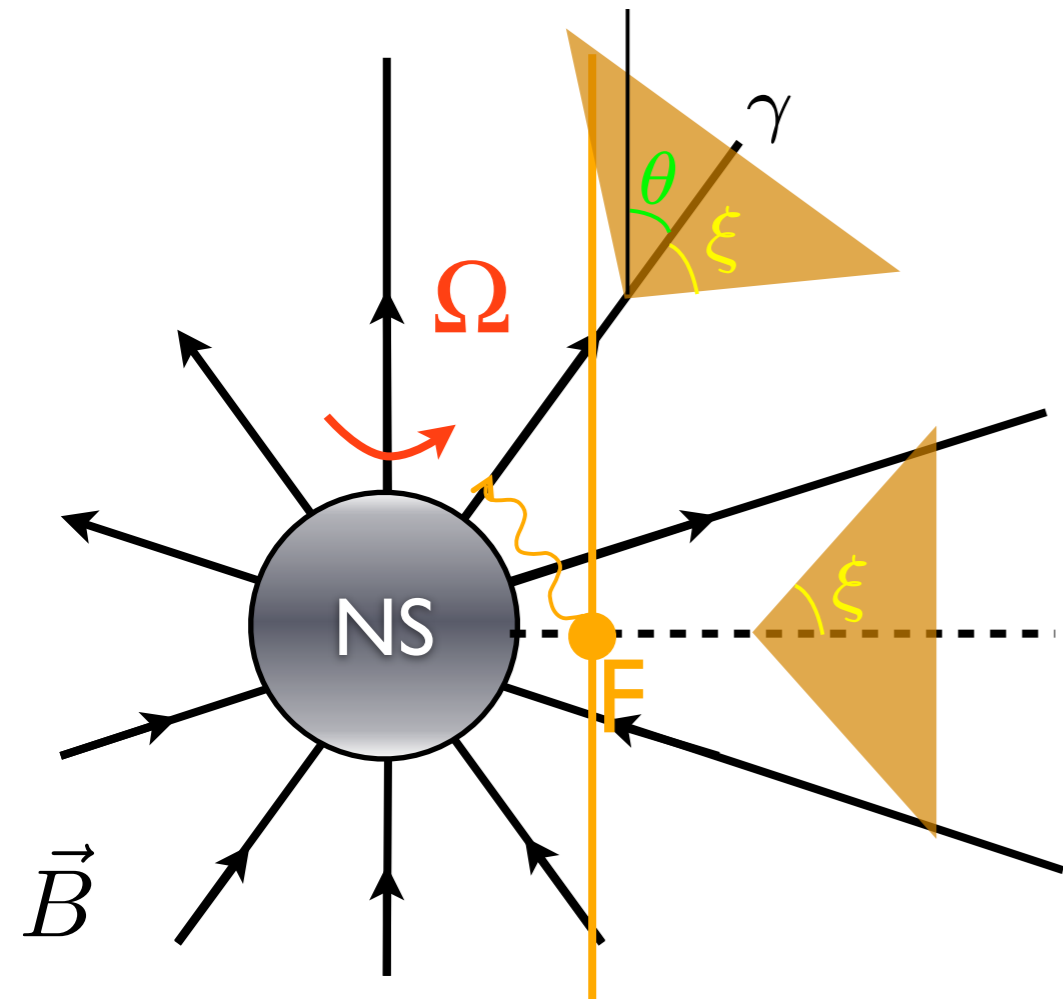
When Can Jets Accelerate?



- Communication is essential
- All signals travel inside the Mach cone ξ :

$$\xi = \frac{\gamma_F}{\gamma} \approx \frac{\sigma^{1/2}}{\gamma}$$

When Can Jets Accelerate?



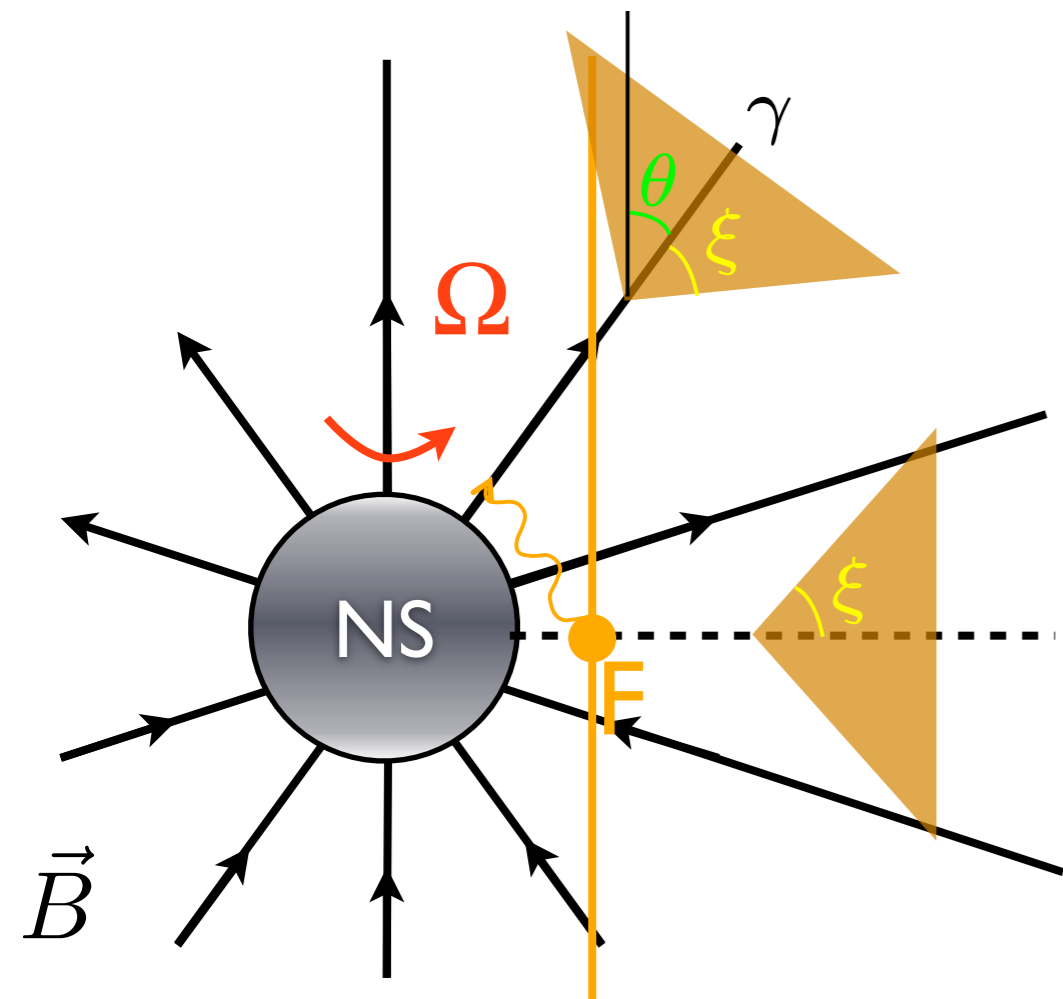
- Communication is essential
- All signals travel inside the Mach cone ξ :

$$\xi = \frac{\gamma_F}{\gamma} \approx \frac{\sigma^{1/2}}{\gamma}$$

- For communication across jet need $\theta \lesssim \xi$, so

$$\gamma\theta \lesssim \sigma^{1/2} = \left(\frac{\mu}{\gamma}\right)^{1/2}$$

When Can Jets Accelerate?



- Communication is essential
- All signals travel inside the Mach cone ξ :

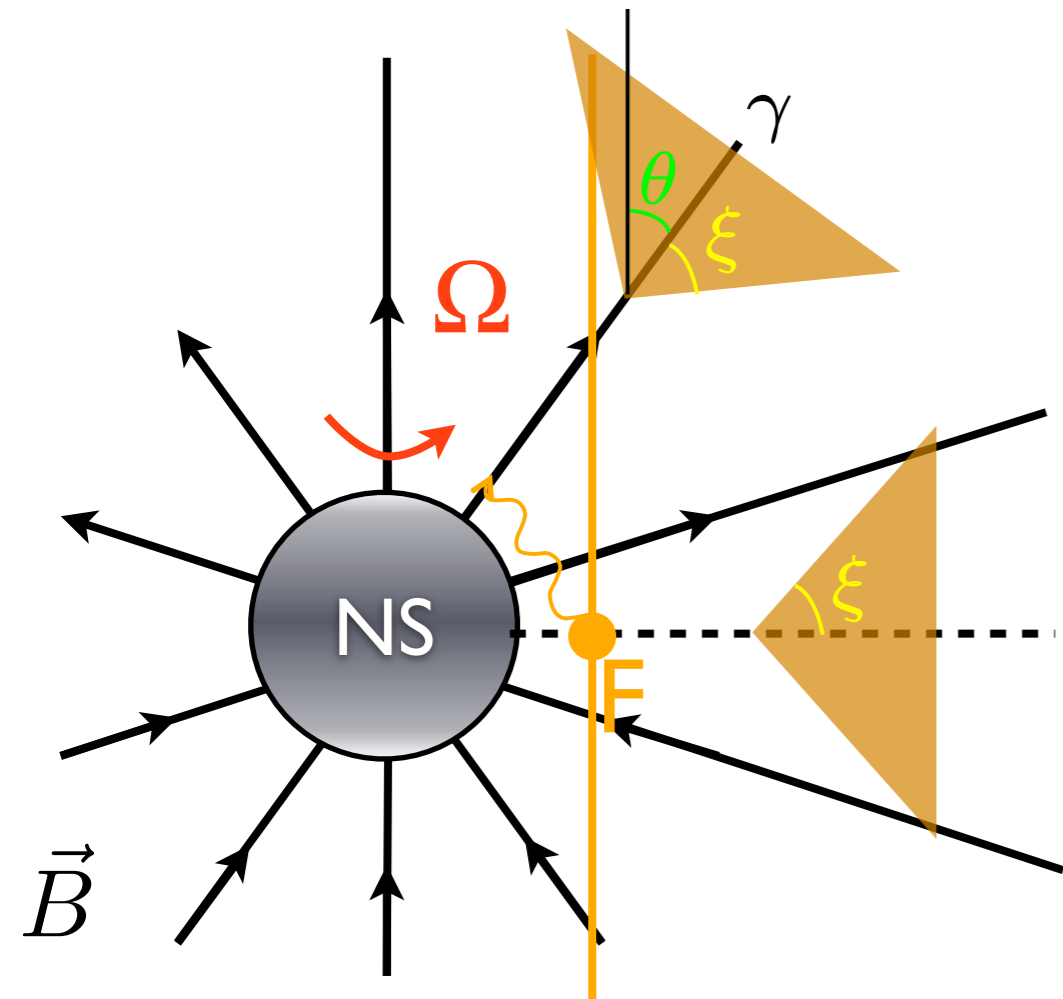
$$\xi = \frac{\gamma_F}{\gamma} \approx \frac{\sigma^{1/2}}{\gamma}$$

- For communication across jet need $\theta \lesssim \xi$, so

$$\gamma\theta \lesssim \sigma^{1/2} = \left(\frac{\mu}{\gamma}\right)^{1/2}$$

- Thus:
- $$\gamma \lesssim \frac{\mu^{1/3}}{\theta^{2/3}}$$

When Can Jets Accelerate?



- Communication is essential
- All signals travel inside the Mach cone ξ :

$$\xi = \frac{\gamma_F}{\gamma} \approx \frac{\sigma^{1/2}}{\gamma}$$

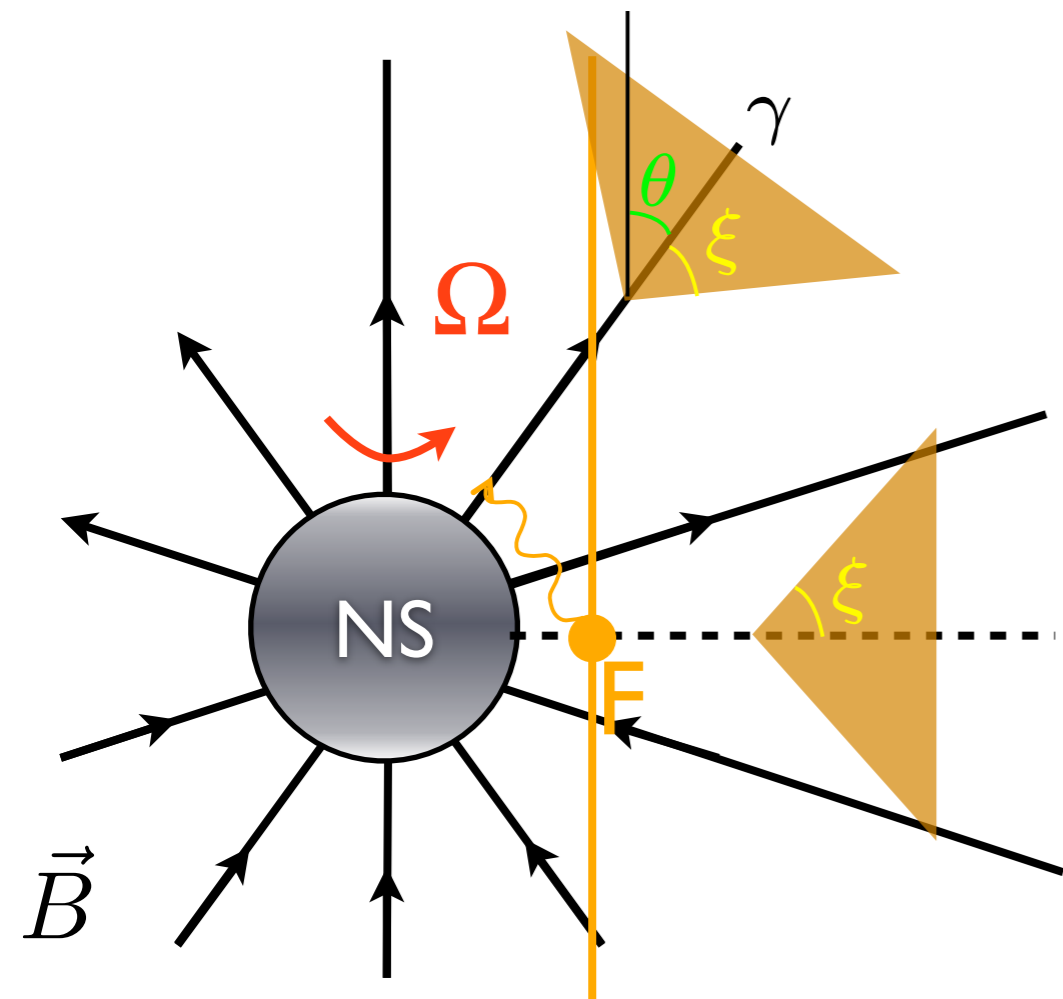
- For communication across jet need $\theta \lesssim \xi$, so

$$\gamma\theta \lesssim \sigma^{1/2} = \left(\frac{\mu}{\gamma}\right)^{1/2}$$

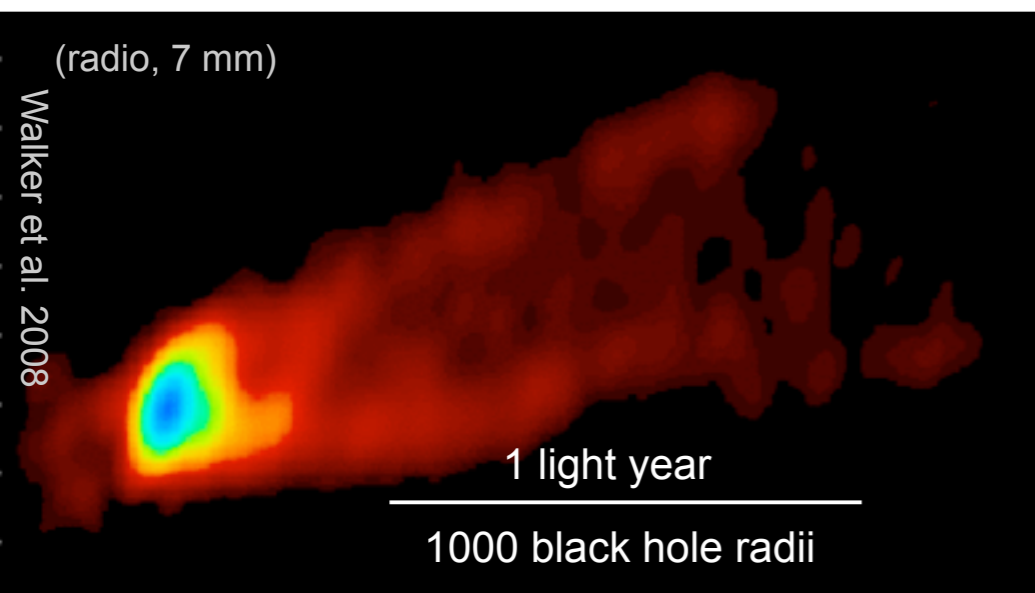
- Thus:
- $$\gamma \lesssim \frac{\mu^{1/3}}{\theta^{2/3}}$$

- Jets accelerate better near the axis

When Can Jets Accelerate?



but, most jets are collimated:



- Communication is essential
- All signals travel inside the Mach cone ξ :

$$\xi = \frac{\gamma_F}{\gamma} \approx \frac{\sigma^{1/2}}{\gamma}$$

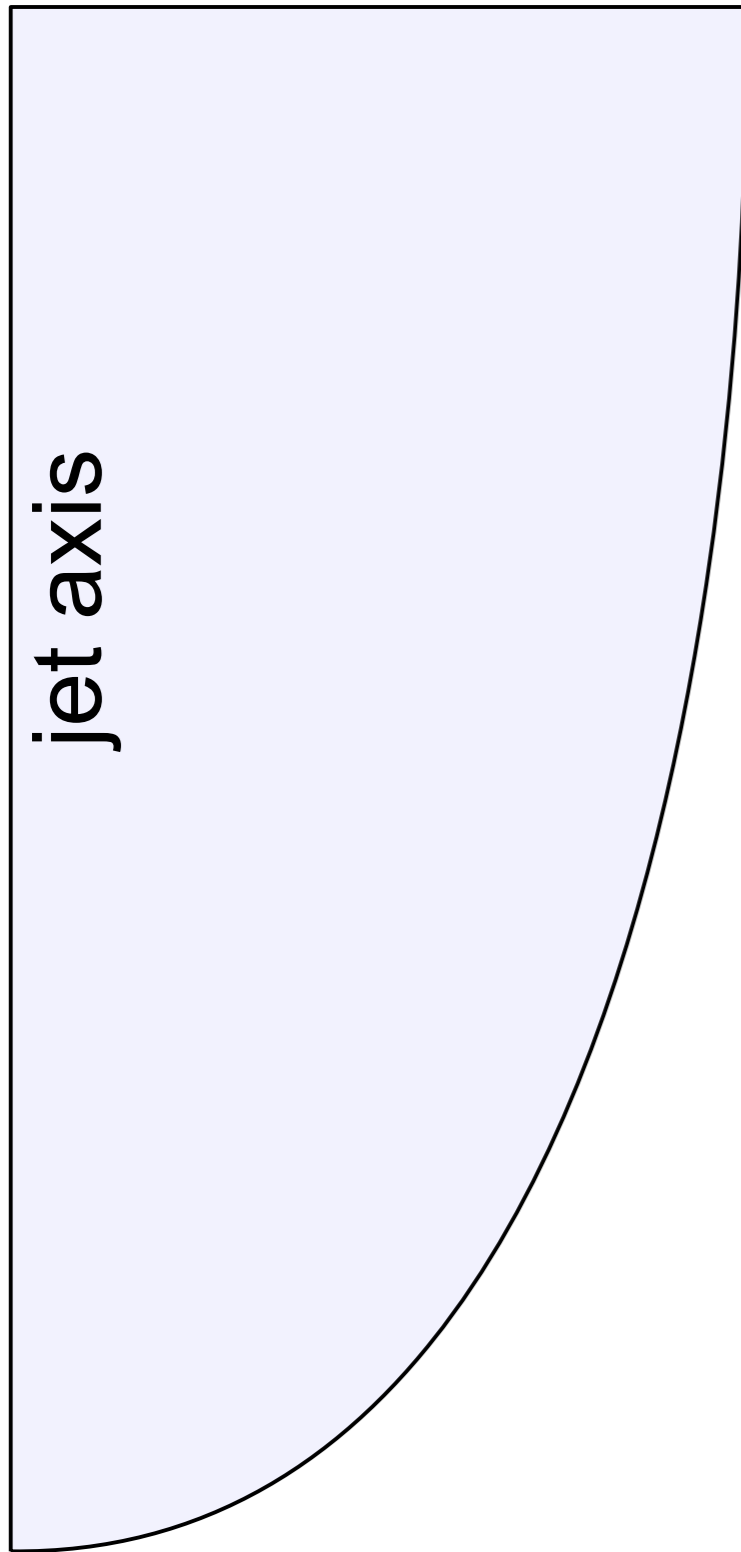
- For communication across jet need $\theta \lesssim \xi$, so

$$\gamma\theta \lesssim \sigma^{1/2} = \left(\frac{\mu}{\gamma}\right)^{1/2}$$

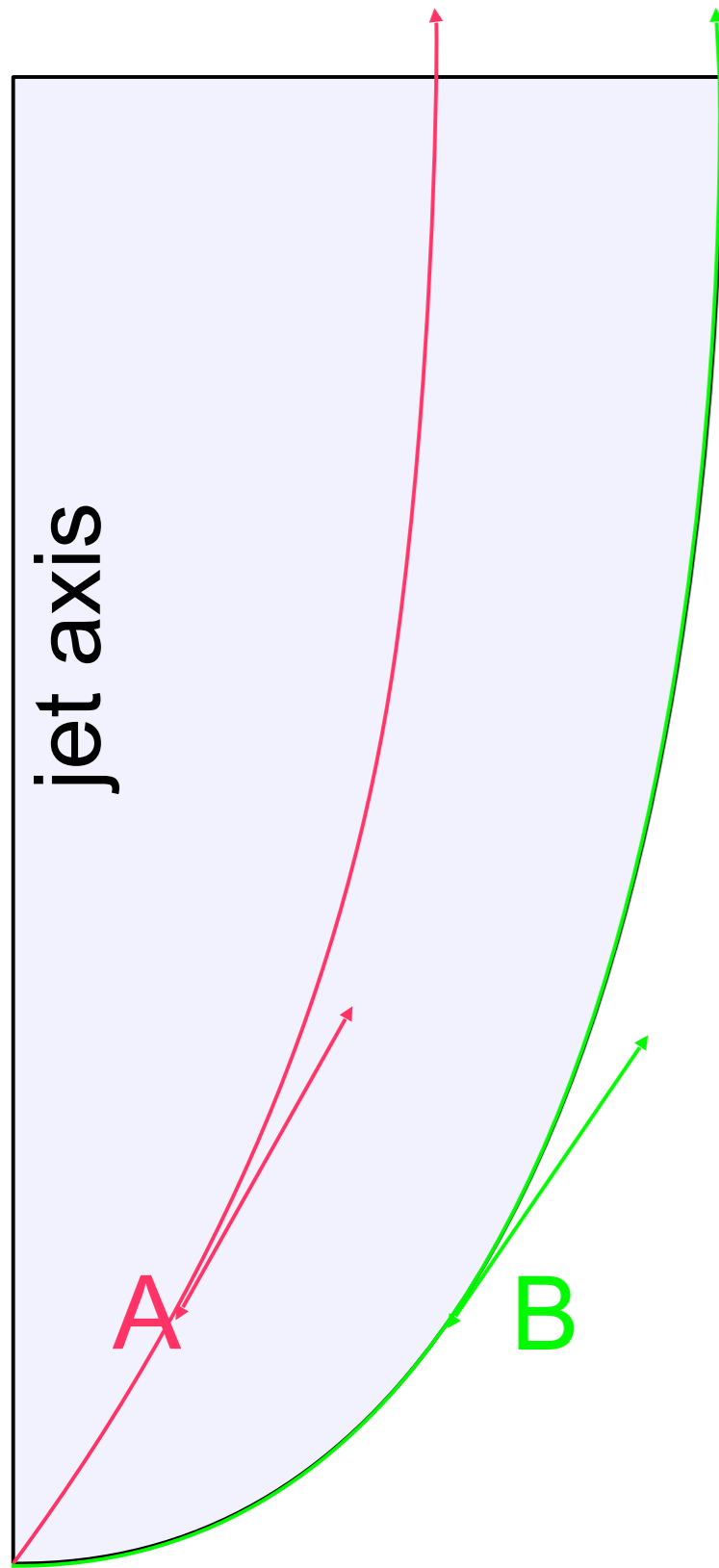
- Thus:
$$\gamma \lesssim \frac{\mu^{1/3}}{\theta^{2/3}}$$

- Jets accelerate better near the axis

How Do Collimated Jets Accelerate?

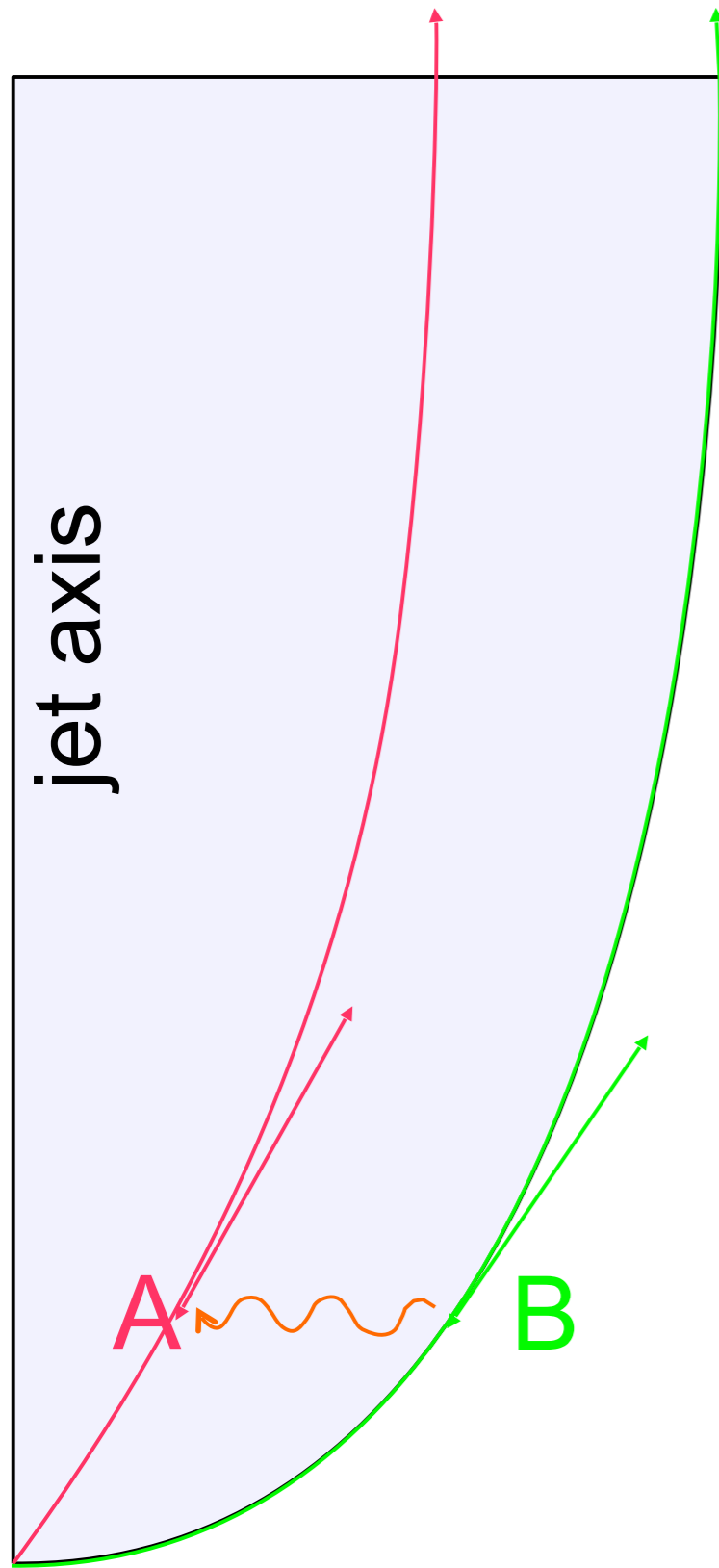


How Do Collimated Jets Accelerate?



- **Communication** is essential

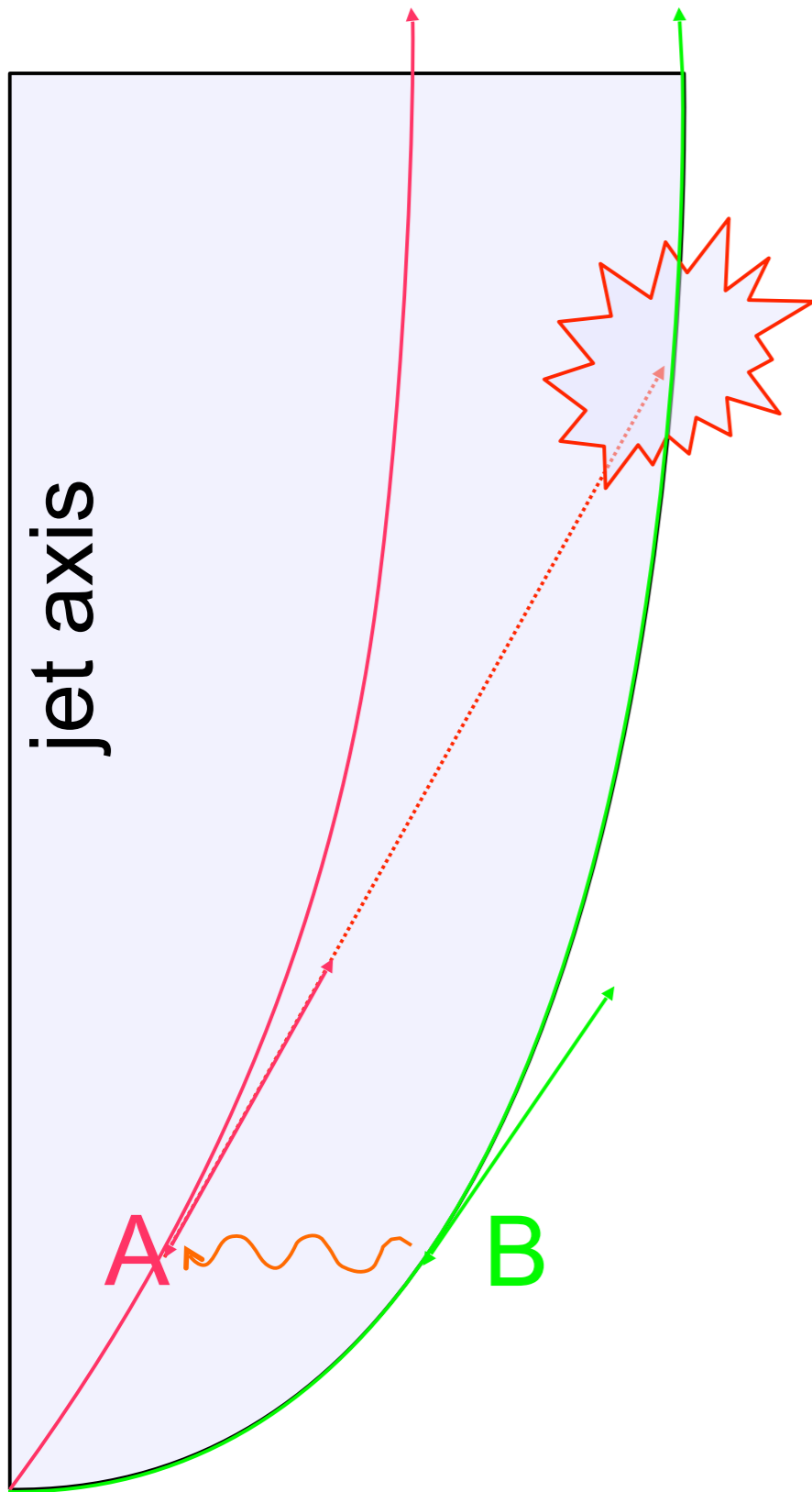
How Do Collimated Jets Accelerate?



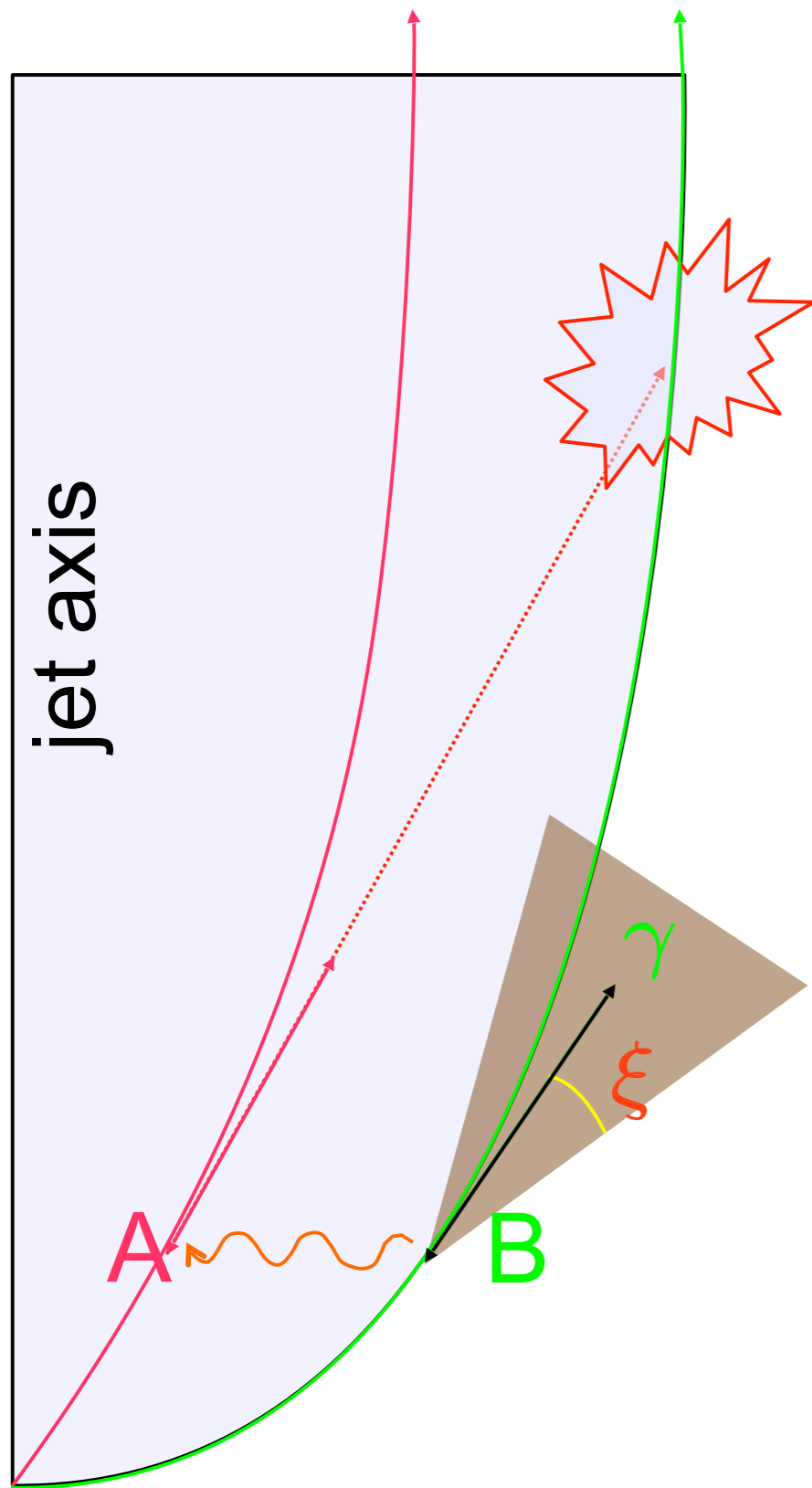
- Communication is essential

How Do Collimated Jets Accelerate?

- Communication is essential to avoid collisions

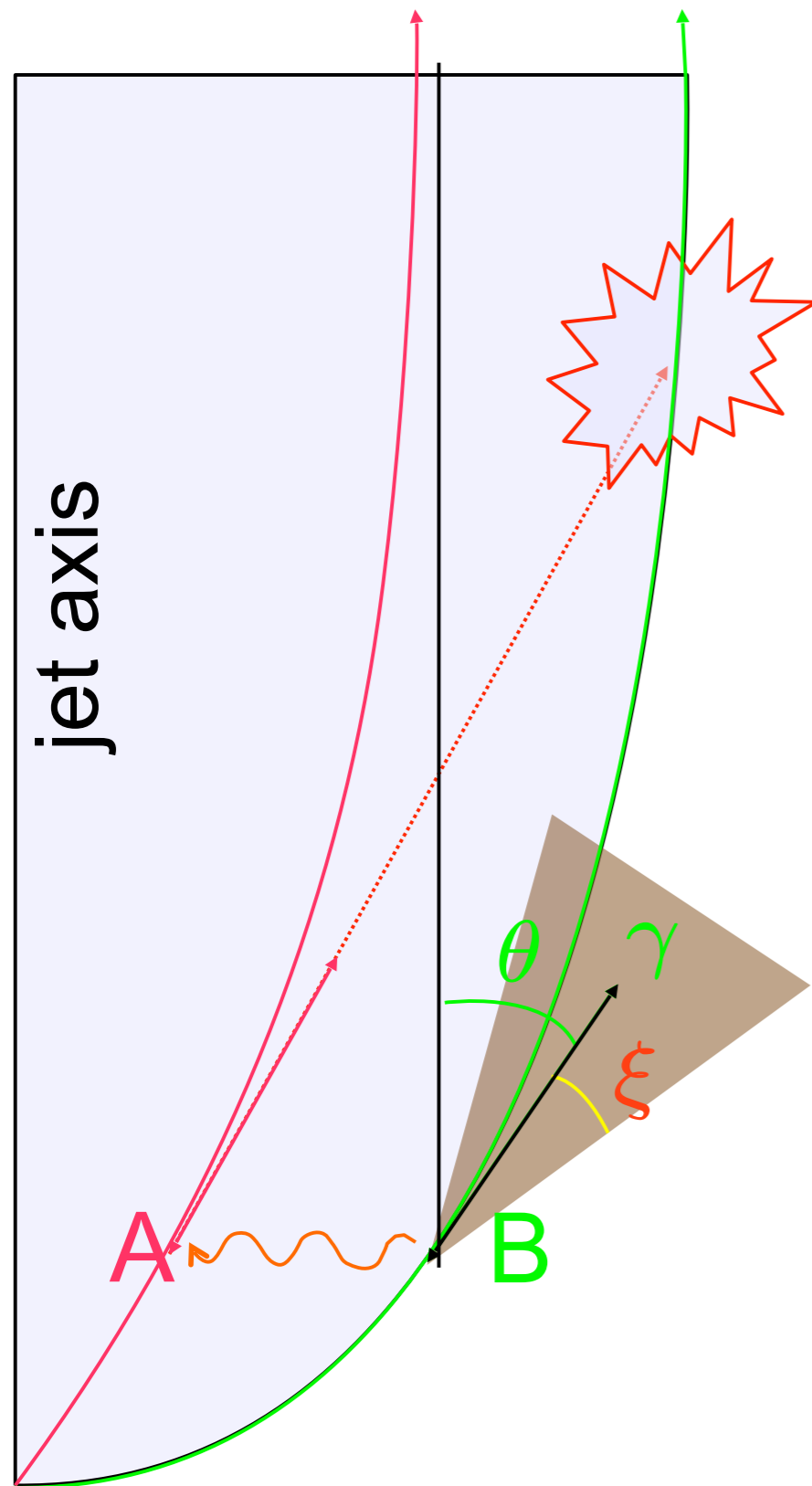


How Do Collimated Jets Accelerate?



- Communication is essential to avoid collisions
- Jet boundary B needs to keep announcing its trajectory to the rest of the jet

How Do Collimated Jets Accelerate?



- **Communication** is essential **to avoid collisions**
- **Jet boundary B** needs to keep **announcing** its trajectory to the **rest** of the jet
- All signals travel inside **Mach cone ξ**:

$$\xi \leq \frac{\gamma_F}{\gamma} = \frac{\sigma^{1/2}}{\gamma}$$
- For **communication** across jet need
 $\theta \lesssim \xi$, so $\theta \lesssim \sigma^{1/2} / \gamma$
- Robust conclusion: $\gamma\theta \lesssim \sigma^{1/2}$
- Collimated jets accelerate efficiently

What Do We Observe?

- *Expect* in collimated jets: $\gamma\theta \lesssim \sigma^{1/2} \lesssim 1$
- Observe:
 - Active Galactic Nuclei: $\gamma\theta \sim 0.1-0.2$
 - Gamma-ray bursts (GRBs): $\gamma\theta \sim 10-100$
- Does it mean that GRB jets are unmagnetized?

GRB Jets: Problem Setup

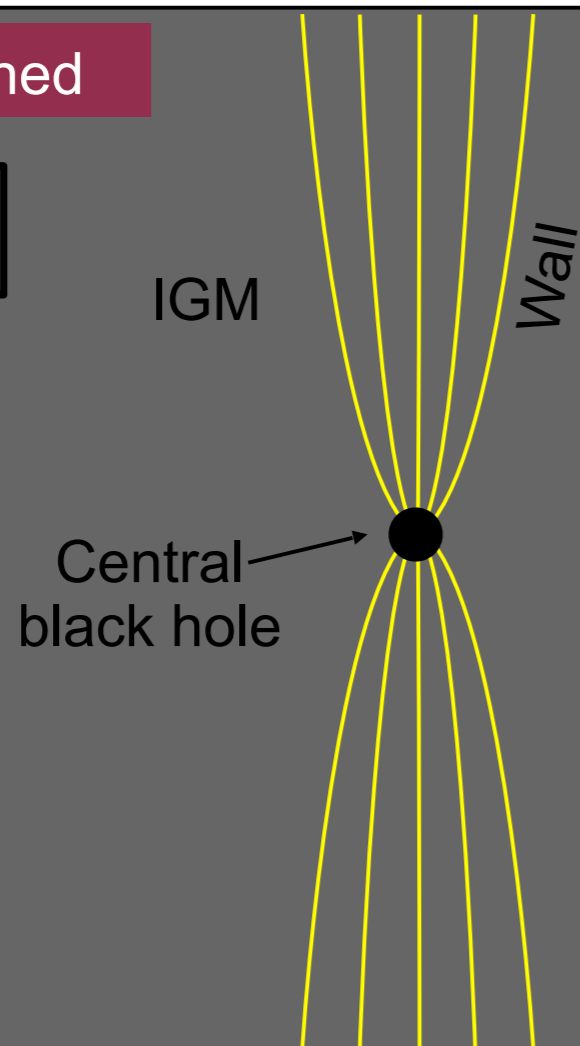
Simulations of magnetized
confined jets:

$$\gamma\theta \lesssim 1$$

(Komissarov et al., MNRAS, 2009)

Confined

$$\gamma\theta = 2$$



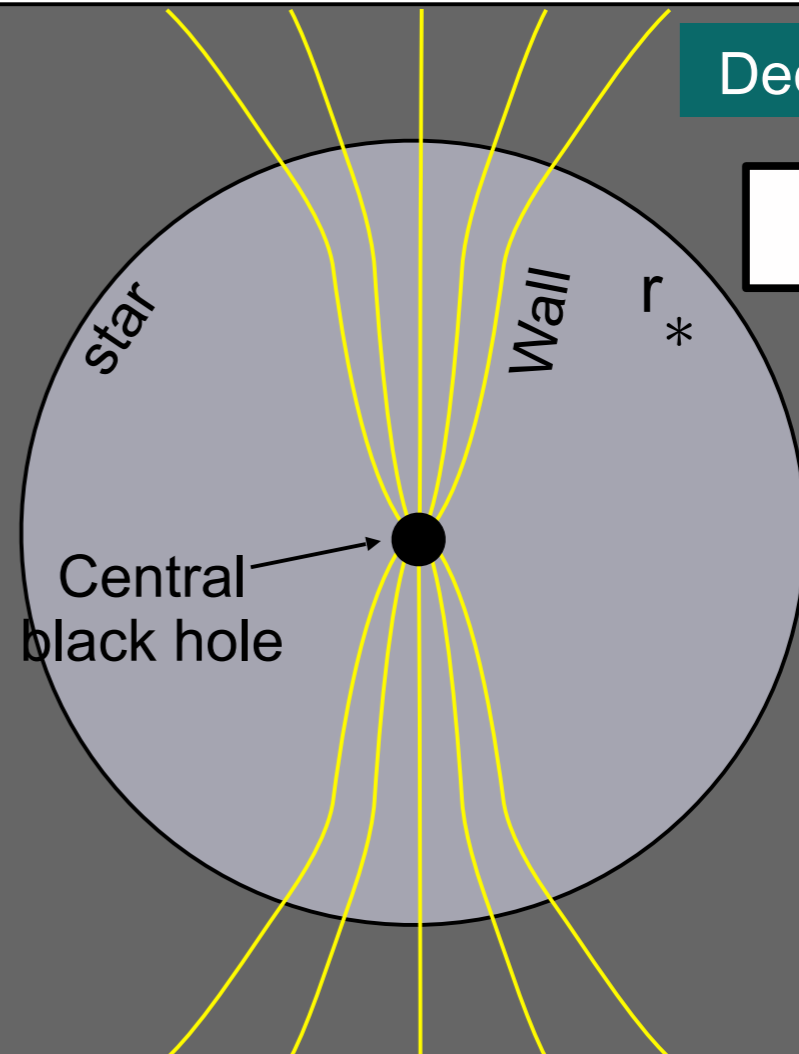
GRB jets are DEconfined:

$$\gamma\theta \gtrsim 10$$

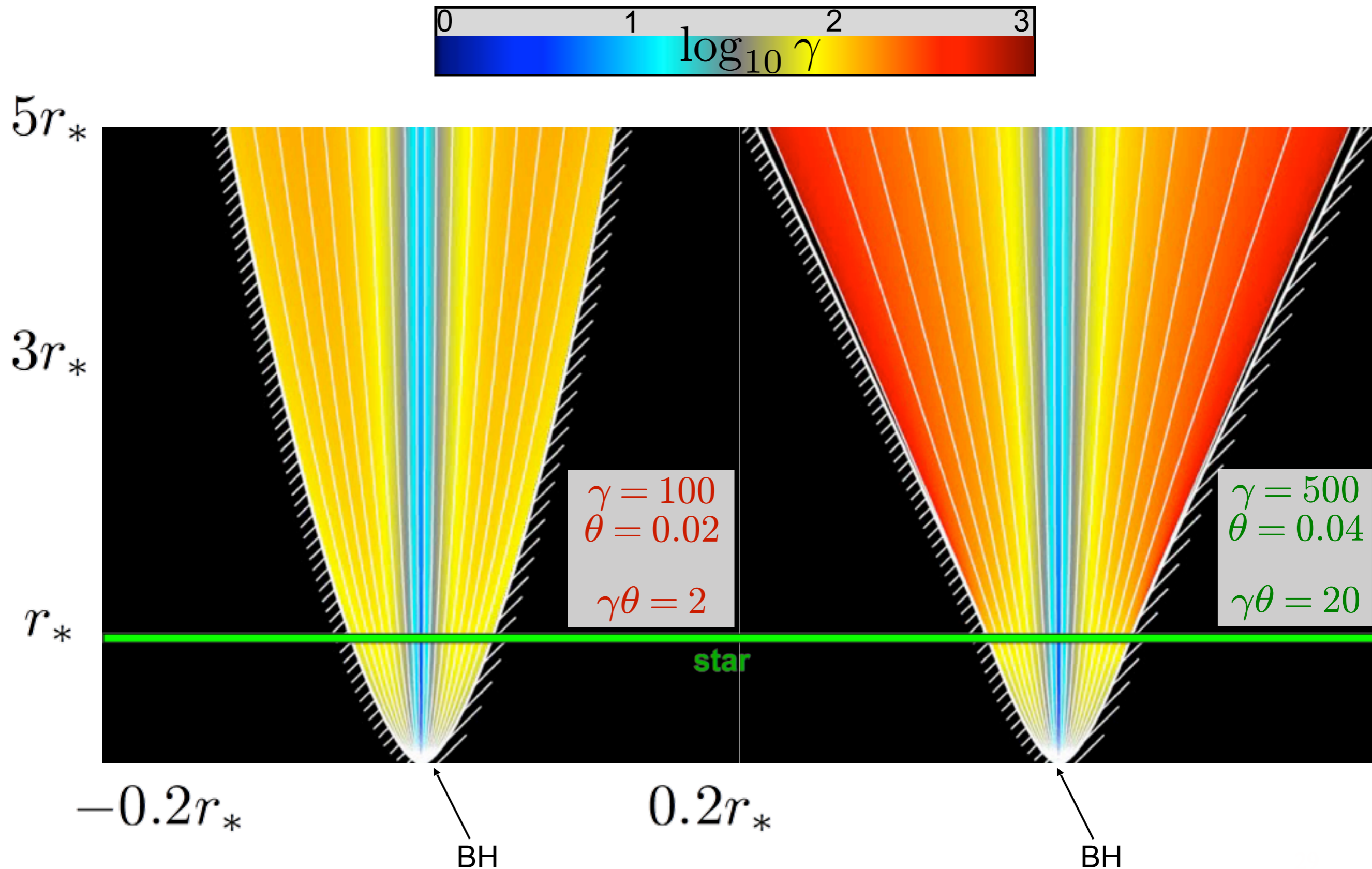
(Tchekhovskoy, Narayan, McKinney, New Astronomy,
2010)

Deconfined

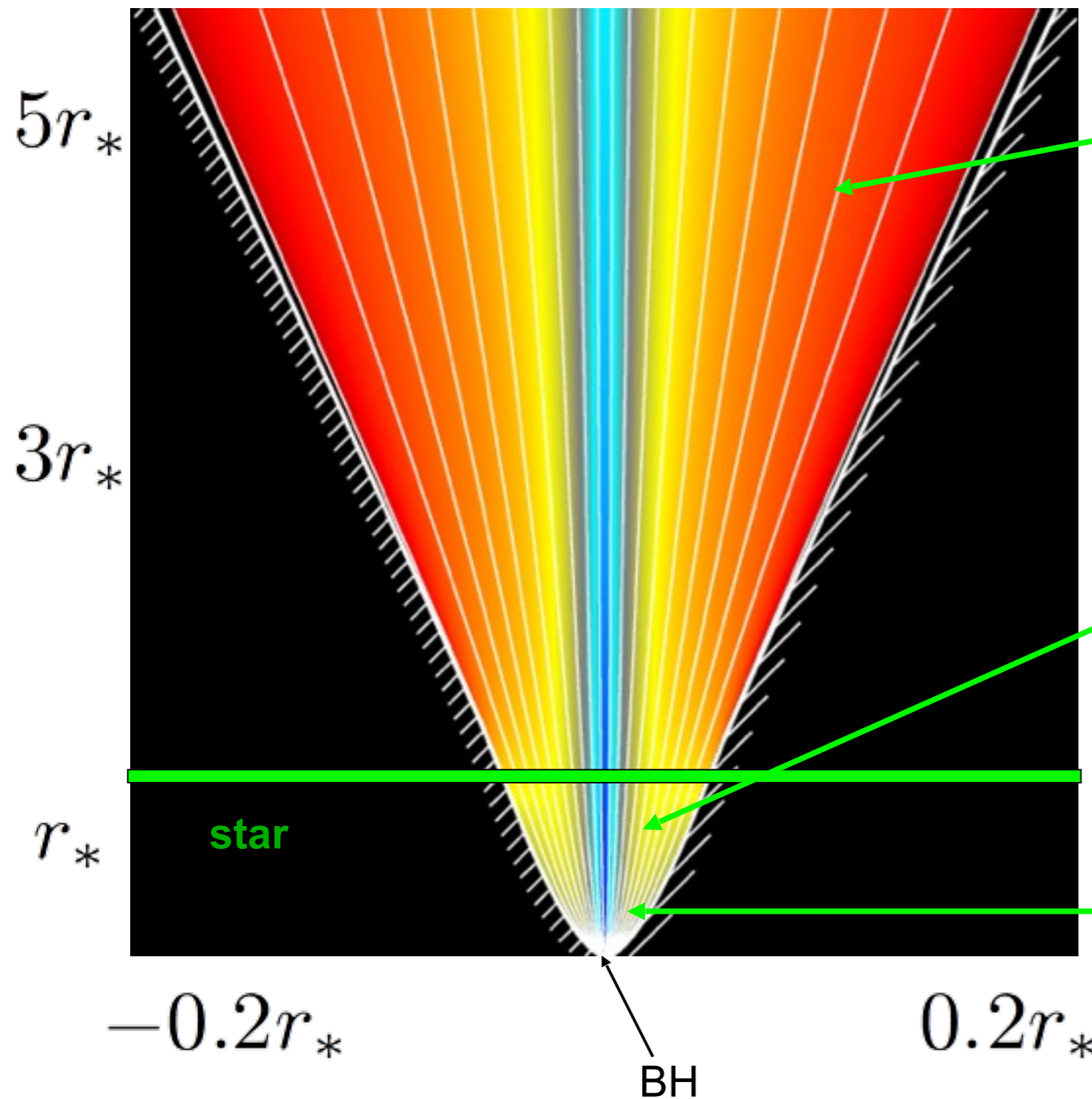
$$\gamma\theta = 20 \checkmark$$



Confined vs. Deconfined



Jet Structure Recap



Fully unconfined jet:

$$\gamma\theta \simeq 20\sigma^{1/2} \quad (\text{AT+ 2010})$$

Fully confined jet, large distance. Centrifugal force limits jet velocity (AT+ 2008):

$$\gamma \approx \left(\frac{R_c}{R} \right)^{1/2}$$

Fully confined jet, small distance. Linear increase:

$$\gamma \approx \Omega R / c \quad (\text{Michel 1969})$$

Magnetic Summary

- Rotation + large-scale magnetic flux \rightarrow jets
- Black holes do not have their own magnetic flux, and rely on accretion disks for flux supply
- Jet power increases with rotational frequency squared and magnetic flux squared
- Jets naturally accelerate magnetically, but only collimating jets do so well
- Many jets are consistent with being powered magnetically, but other processes such as radiative driving can also be at play

Homework

- Exercises with HARMPI code: fully parallel, 3D general relativistic MHD code
 - MONOPOLE_PROBLEM_1D
 - MONOPOLE_PROBLEM_2D
- Documentation and download at:
<https://github.com/atchekho/harmpi>