#### Introduction to interferometry and VLBI

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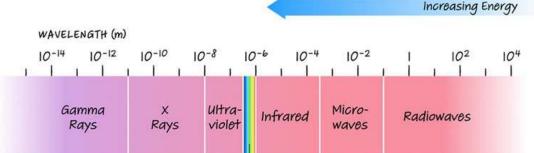
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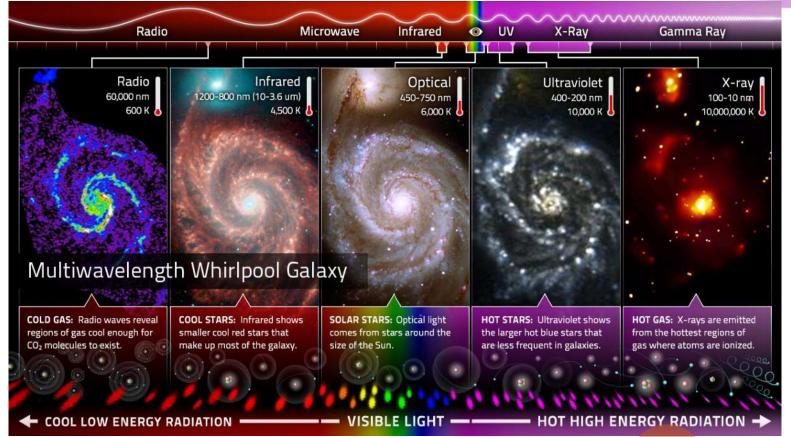
HARVARD UNIVERSITY

# Part 1: Radioastronomy and Interferometry in a Nutshell

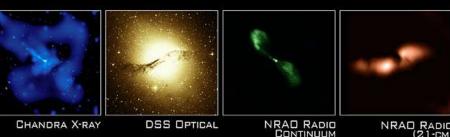
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#### Multi-wavelength Astronomy



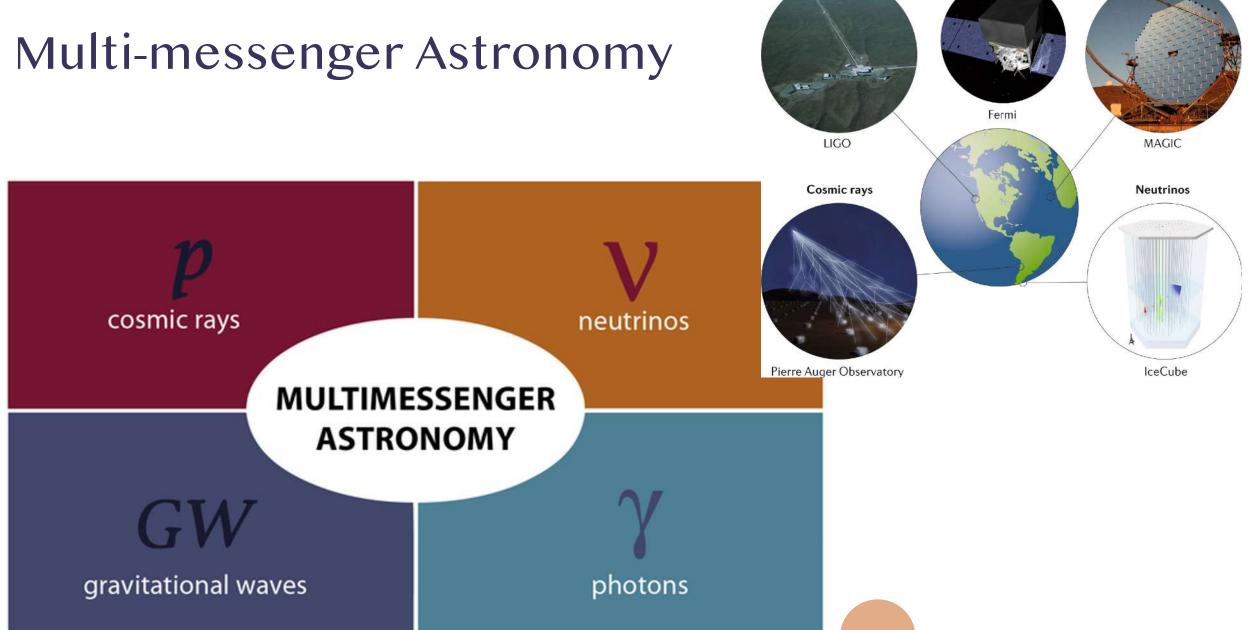






Spiral Galaxy M51

Active galaxy Centaurus A



Gravitational waves

γ-rays

γ-rays

## Astronomy Nobel Prizes



1974: M. Ryle & A. Hewish (J. Bell)
Imaging synthesis; pulsar



2011: S. Perlmutter, B.Schmidt & A. Riess Dark energy (Cosmic expansion)



1978: A. Penzias & R. Wilson

Cosmic Microwave Background



2017: R. Weiss, B. Barrish & K. Thorne Gravitational wave detection



1993: R. Hulse & J. Taylor

Double pulsar



2019: M. Mayor & D. Queloz Extrasolar planets



2002: R. Giacconi, M. Koshiba & R. Davis Jr. Cosmic X-ray and neutrino sources



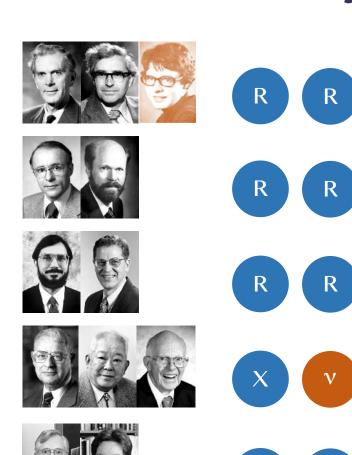
2020: R. Genzel & A. Ghez Black hole @ Galactic center



2006: J. Mather & G. Smoot

Cosmic Microwave Background

## Astronomy Nobel Prizes









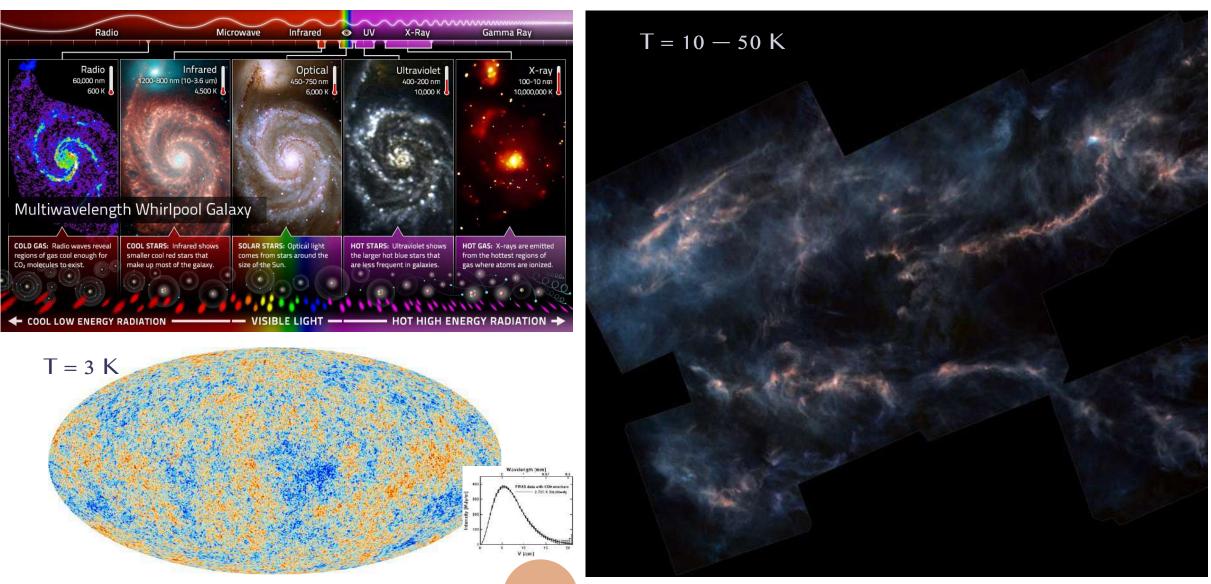








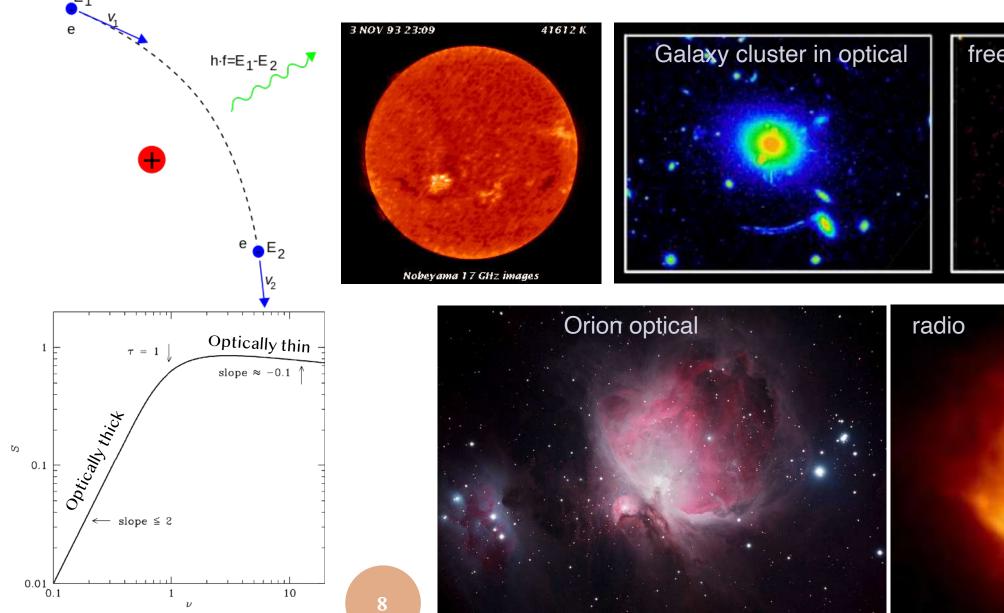
#### Emission processes: thermal (blackbody)

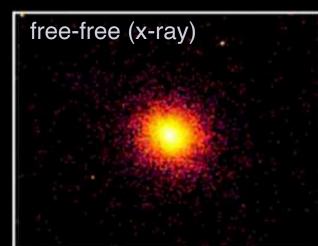


The cosmological microwave background

Taurus clouds with the Herschel satellite (150, 250, 350, 500  $\mu$ m)

#### Emission processes: free-free (thermal)

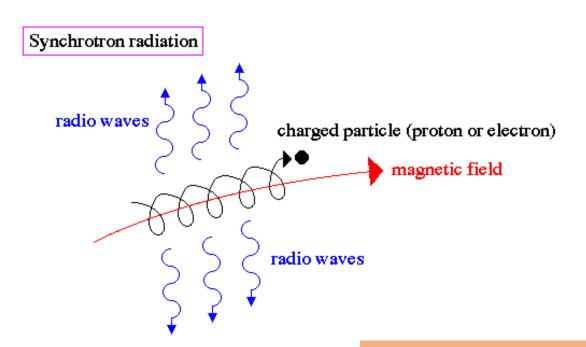


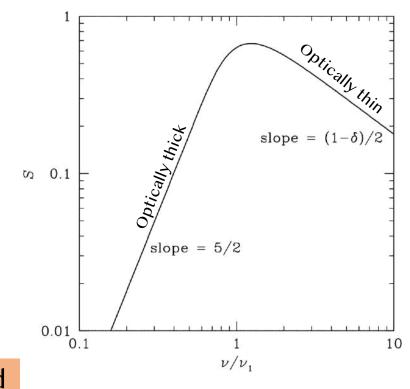




#### Emission processes: synchrotron (non-thermal)

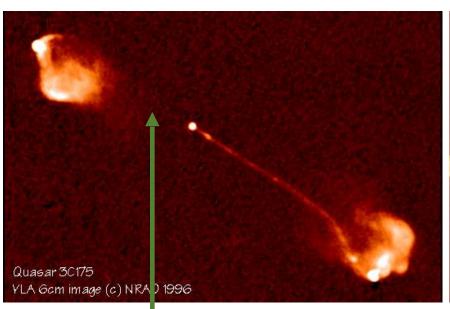
Radioastronomy (and, as we will see, particularly VLBI) is the realm of non-thermal processes

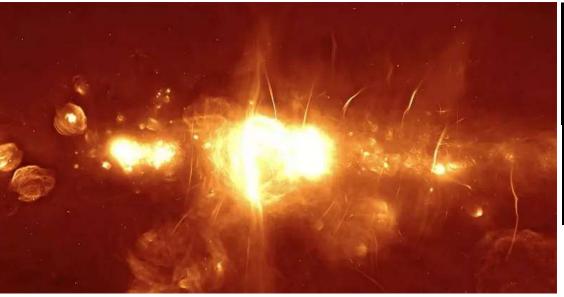


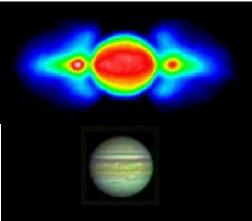


Relativistic electrons needed (acceleration mechanism)

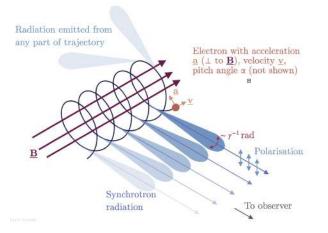
#### Emission processes: synchrotron (non-thermal)



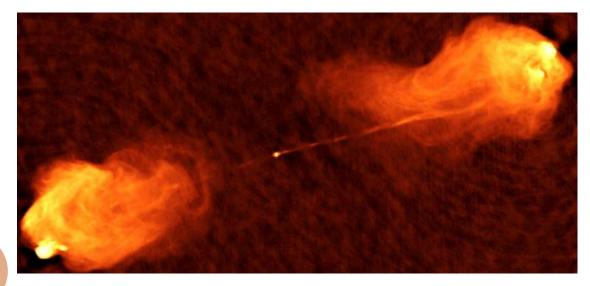




Quiz: why isn't there a counterjet?

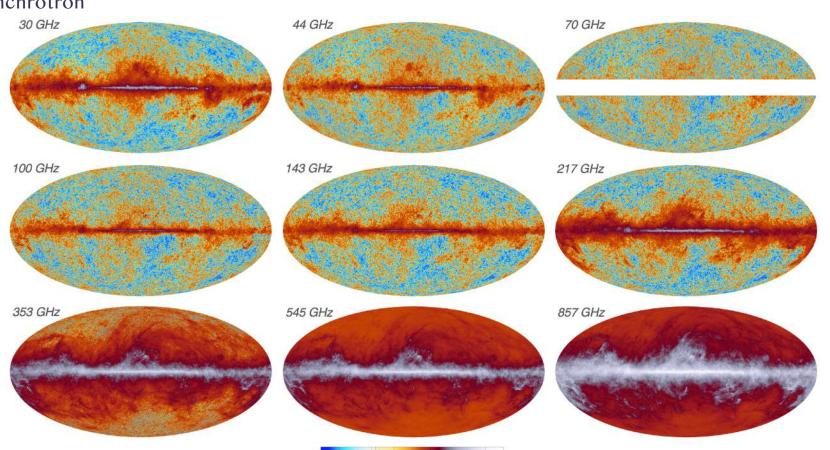


The Galactic center in radio waves



# All-sky Planck microwave maps

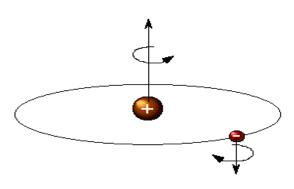
At low frequencies, strong contamination by foreground freefree and synchrotron The sweet spot for cosmological microwave background

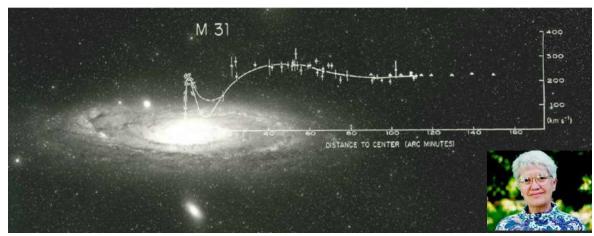


30-353 GHz: 8T [uK....]: 545 and 857 GHz: surface brightness [kJv/sr

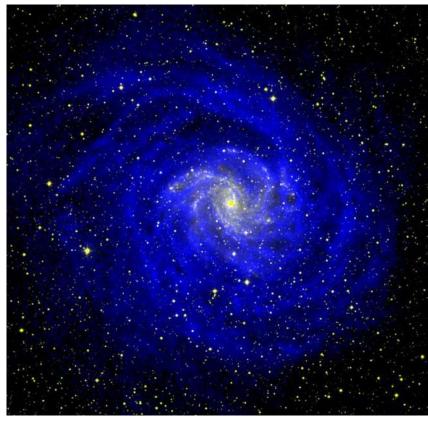
At high frequencies, strong contamination by foreground dust

## Line emission: 21-cm hydrogen line





Vera Rubin

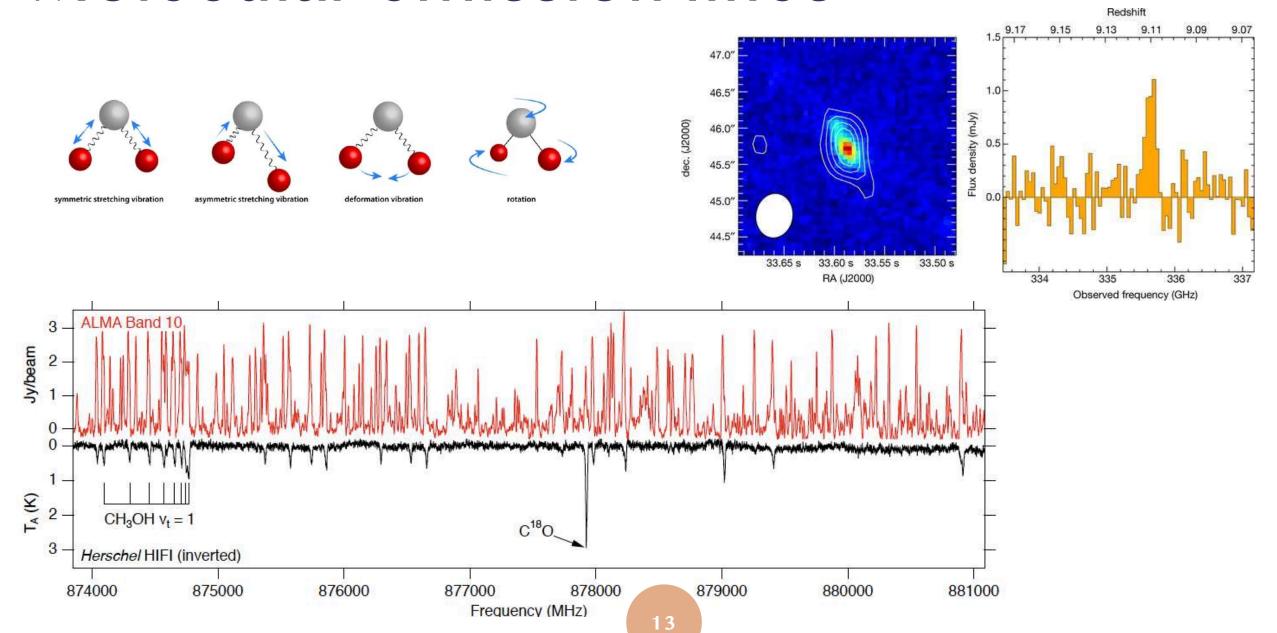


Optical image

21-cm map

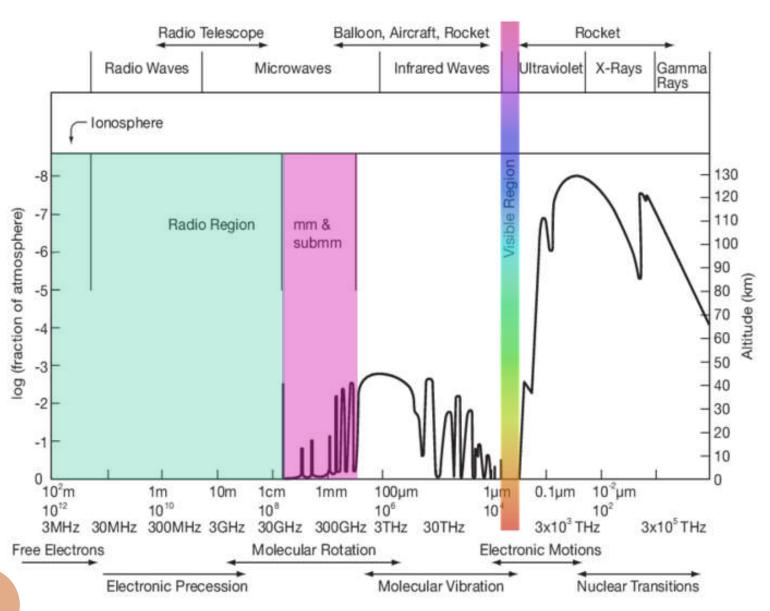
#### Molecular emission lines

#### Oxygen at z = 9 with ALMA

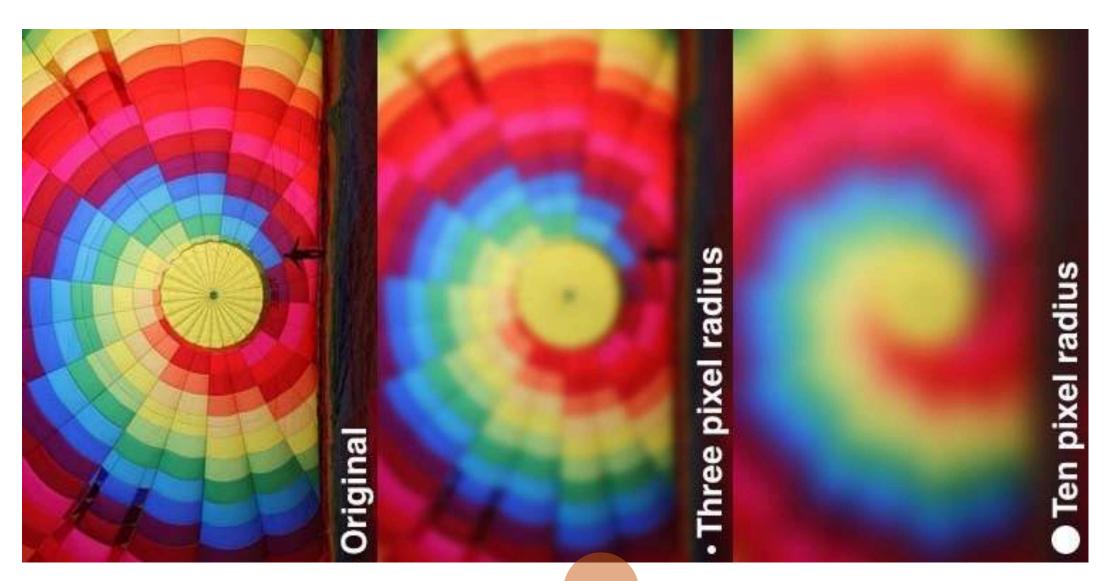


## Atmospheric transparency/opacity





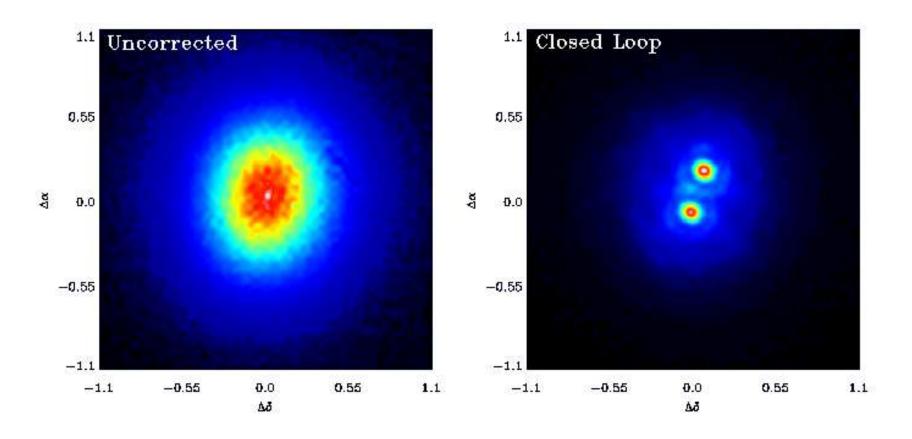
#### Angular resolution (image sharpness)



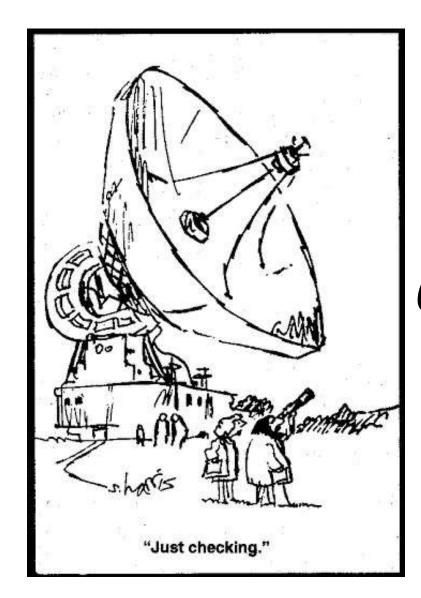
#### Angular resolution definition

The angular resolution is the minimum separation that an optical instrument can resolve:

$$hetapproxrac{\lambda}{D}$$



#### The angular resolution **problem** in radioastronomy



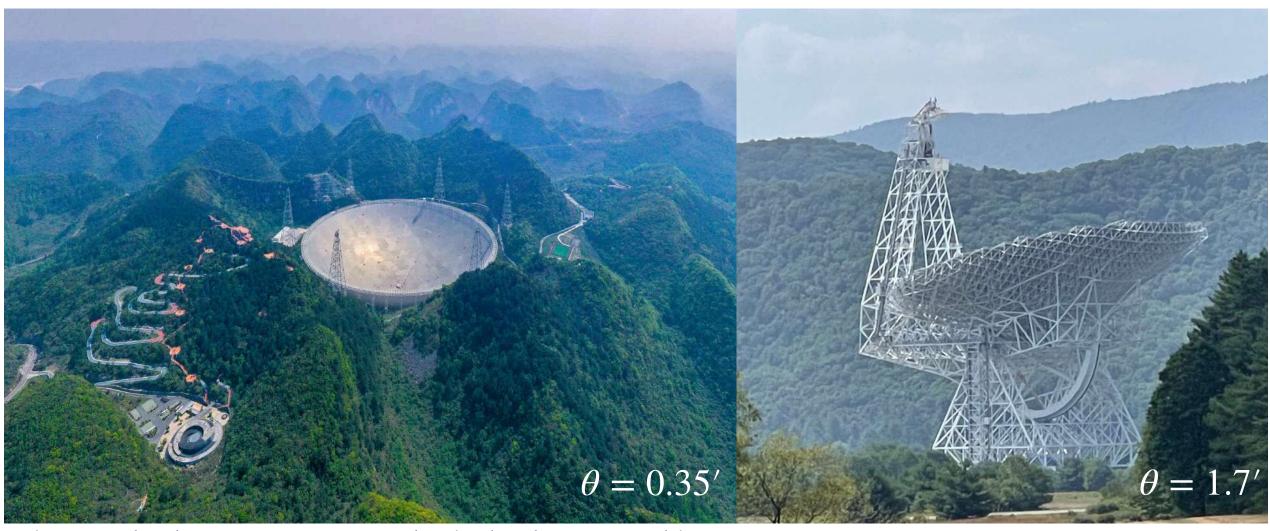
$$\lambda = 5 \text{ cm } \& D = 160 \text{ m} \longrightarrow \theta = 1'$$

$$\theta pprox \frac{\lambda}{D}$$

Radio-astronomers need to build **BIG** telescopes

$$\lambda = 0.5 \ \mu \text{m} \& D = 1.6 \ \text{mm} \longrightarrow \theta = 1'$$

#### That's hard and expensive...



Five Hundred Meter Aperture Spherical Telescope (China)

Green Bank Telescope (100m, USA)

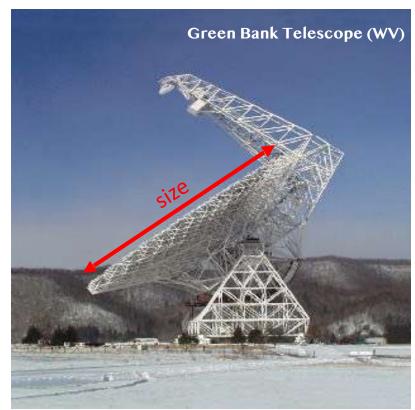
## ... and potentially risky



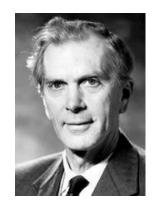
300 foot telescope in Green Bank (November 15, 1988)

300 foot telescope in Green Bank (November 16, 1988)

#### Two kinds of radio telescopes







Sir Martin Ryle (Nobel 1974) Imaging synthesis

Single-dish

$$hetapprox rac{\lambda}{D}$$

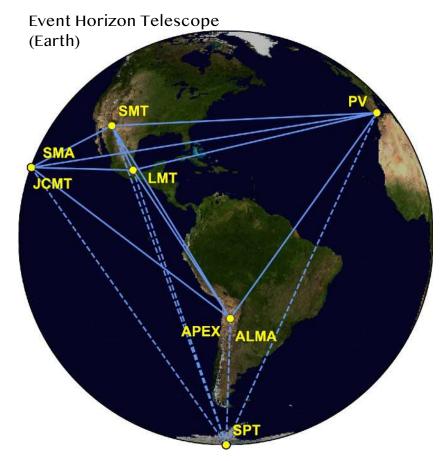
Interferometers

$$heta pprox rac{\lambda}{B_{max}}$$
 Maximum "baseline"

#### Very Long Baseline Interferometry (VLBI)



1 milli-arcsecond resolution



25 microarcsecond resolution

#### The angular resolution solution in radioastronomy







Biggest single-dish

Same resolution as human eye (~ one arcminute)

"conventional" interferometers

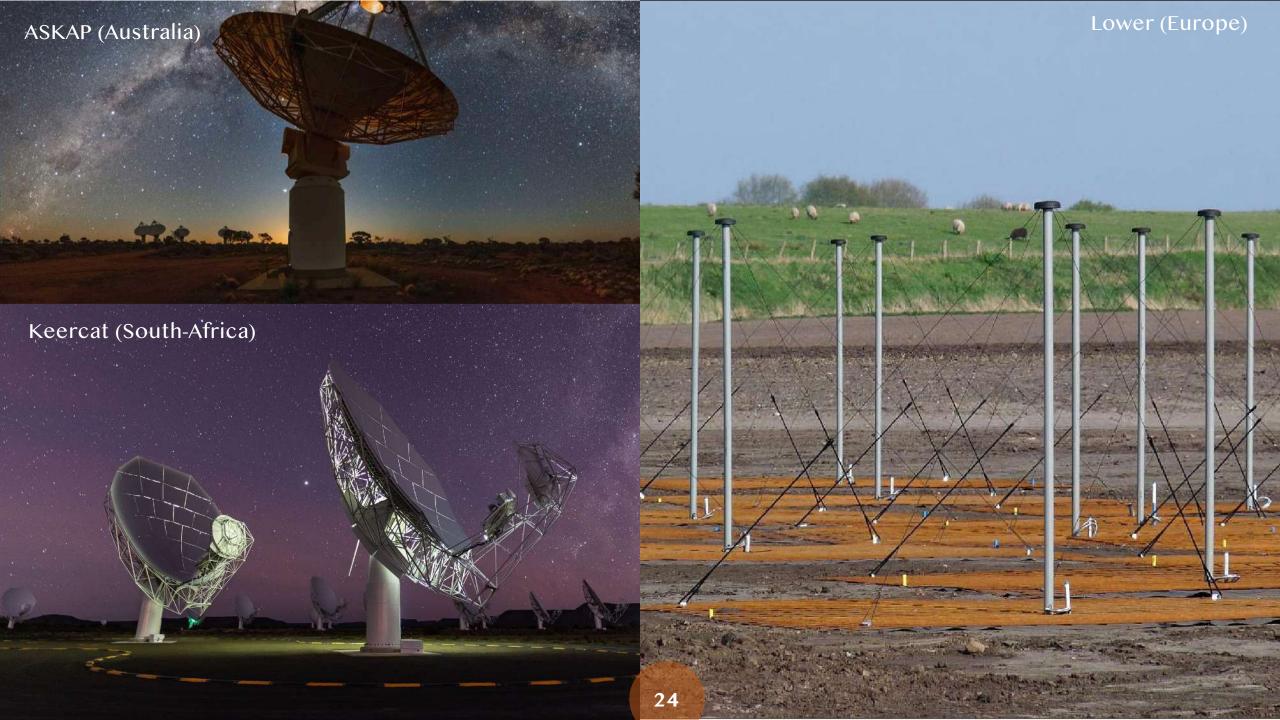
Same resolution as large optical telescope (~ 0.1 arcsecond)

#### **VLBI** arrays

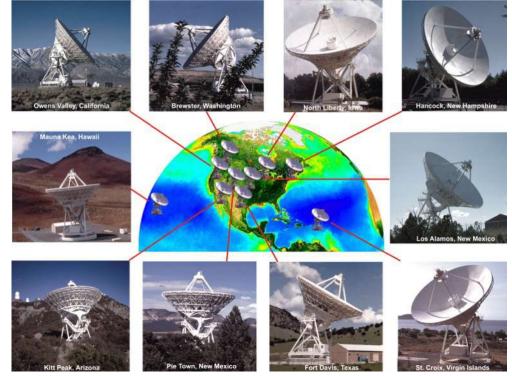
Highest resolution in astronomy
(~ 1 milli-arcsecond down to 10 micro-arcsecond)

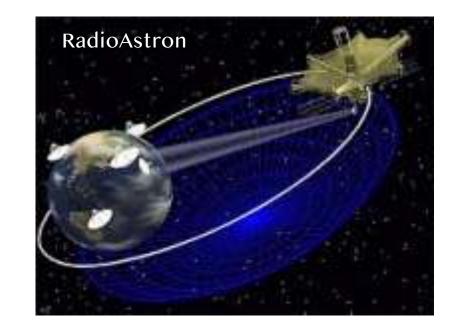
This is the highest angular resolution achievable in all of astronomy



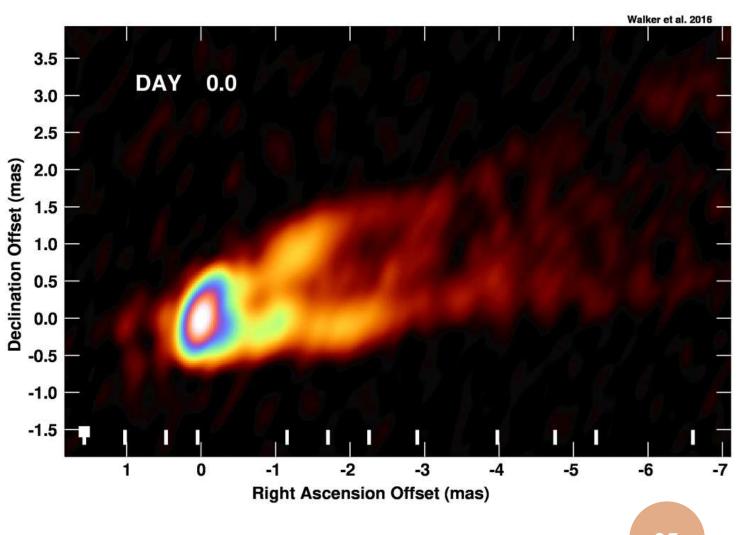


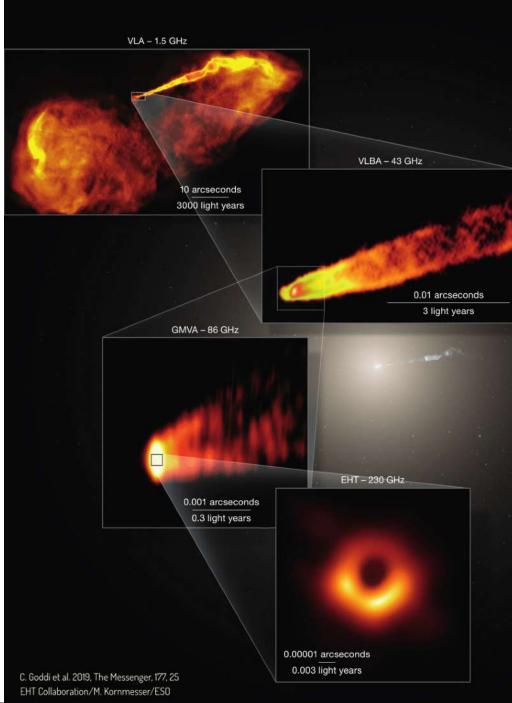






# VLBI highlights



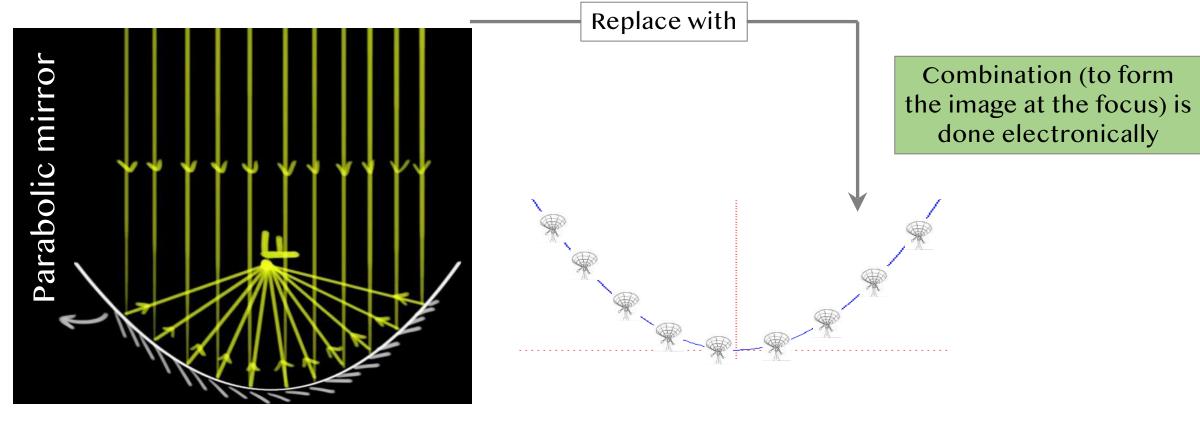




#### Fair warning: interferometry is more work than single-dish...

- Observing with N  $\gg$  1 antennas, is N times more work than observing with one antenna (or maybe even N<sup>2</sup> times more work...)
- As we will see in details later, interferometers do not provide directly an image. Rather they deliver complex quantities ( $\in \mathbb{C}$ ) that need to be manipulated mathematically to computationally reconstruct an image
- The resulting images contain artifacts caused by the geometry of the array
- Sensitivity is limited
- Calibration is more work than for single-dish telescopes

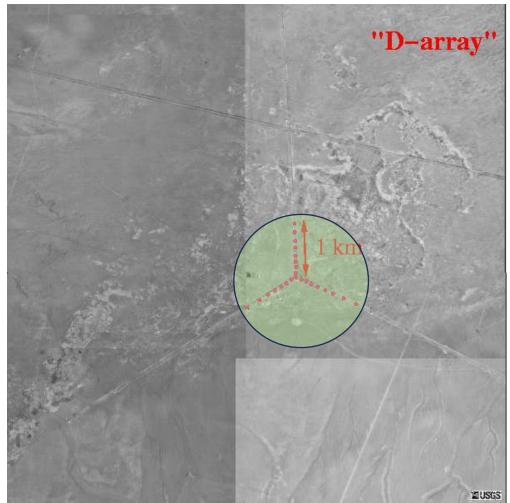
## Pictorial principle of interferometry

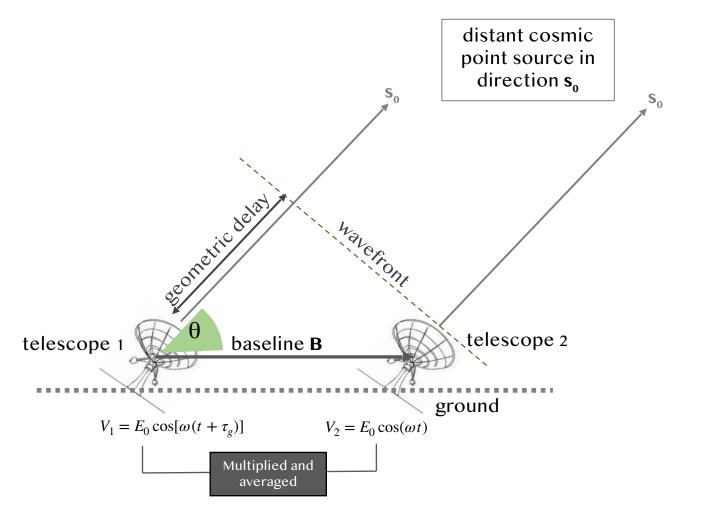


Reflector telescope

#### Interferometers as fragmented mirrors







- Each telescope of the array acts as a **coherent sensor** of the incoming electric field **E**
- $\hfill \square$  The geometric delay ,  $\tau_g$  , is:

$$\tau_g = \frac{\mathbf{B} \cdot \hat{\mathbf{s}}_0}{c}$$

☐ The output of the interferometer is:

$$R_c = P_0 \cos(\omega \tau_g)$$

 $P_0 = E_0^2/2$  is the received power in the EM wave

☐ If the source is not point-like,

$$R_c = \iint I(\hat{\mathbf{s}})\cos(\omega \tau_g) d\Omega$$

 $I(\hat{\mathbf{s}})$  is the sky brightness

By adding a  $\pi/2$  phase delay to the output of telescope 2 before multiplying, we can get a complementary interferometer response:

$$R_s = \iint I(\hat{\mathbf{s}}) \sin(\omega \tau_g) d\Omega$$

 $\Box$  The specialized hardware that produces R<sub>c</sub> and R<sub>s</sub> from the output of the telescopes is called a "correlator"

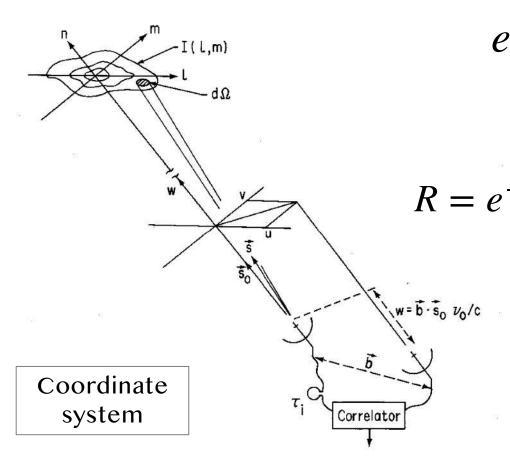
 $\square$  We now **define** a new **complex** function R, by combining R<sub>c</sub> and R<sub>s</sub>:

$$R = R_c - iR_s$$

 $\square$  From the formulae of R<sub>c</sub> and R<sub>s</sub>, we see that:

$$R = \iint I(\hat{\mathbf{s}})e^{-i\omega\tau_g}d\Omega = \iint I(\hat{\mathbf{s}})e^{-2\pi i\mathbf{B}\cdot\hat{\mathbf{s}}/\lambda}d\Omega$$

 $\square$  Because R is a complex number, the specialized hardware that produces R (i.e. R<sub>c</sub> and R<sub>s</sub>) from the output of the telescopes is called a "complex correlator"



$$e^{-2\pi i \mathbf{B} \cdot \hat{\mathbf{s}}/\lambda} = e^{-2\pi i \mathbf{B} \cdot \hat{\mathbf{s}}_0/\lambda} \cdot e^{-2\pi i (ul + vm)}$$

$$R = e^{-2\pi i \mathbf{B} \cdot \hat{\mathbf{s}}_0 / \lambda} \iint I(l, m) e^{-2\pi i (ul + vm)} \frac{dldm}{\sqrt{1 - l^2 - m^2}}$$

Complex visibility function

#### Fundamental result

$$R = e^{-2\pi i \mathbf{B} \cdot \hat{\mathbf{s}}_0 / \lambda} \iint I(l, m) e^{-2\pi i (ul + vm)} \frac{dldm}{\sqrt{1 - l^2 - m^2}}$$

Response of the interferometer

$$V(u,v) = \iint I(l,m)e^{-2\pi i(ul+vm)} \frac{dldm}{\sqrt{1-l^2-m^2}}$$
 Complex visibility function

Van Cittert-Zernike Theorem

$$I(l, m) = \sqrt{1 - l^2 - m^2} \iint V(u, v) e^{2\pi i(lu + mv)} du dv$$

Sky brightness

Fourier Transform

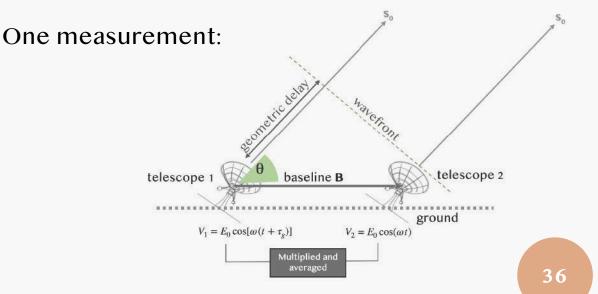
#### But, careful...

$$V(u, v) = \iint I(l, m)e^{-2\pi i(ul + vm)} \frac{dldm}{\sqrt{1 - l^2 - m^2}}$$

This is a **function** of (u, v)

$$I(l, m) = \sqrt{1 - l^2 - m^2} \iint V(u, v) e^{2\pi i(lu + mv)} du dv$$

This is a **function** of (l, m)



corresponds to one baseline: B,

and therefore also to one value of  $(u_i, v_i)$ 

(actually two points: also  $(-u_i, -v_i)$ Including both baselines from telescope 1 to telescope 2 and from telescope 2 to telescope 1)

#### Summary

- The complex visibility function, V(u, v), and the sky brightness, I(l, m), are Fourier conjugates (van Cittert-Zernike Theorem).
- Measuring the visibility function and taking the inverse Fourier transform enables us to obtain the sky brightness.
- This is the (electronic and computational) process by which radio interferometers re-construct an image "in the focal plane" of the telescope they simulate.
- For each observation with two antennas, one only gets one the values of the complex visibility at two points  $(u_i, v_i)$  and  $(-u_i, -v_i)$ .