

# GMTNIRS and the Promise of High Resolution Spectroscopy in the Infrared: Part 1

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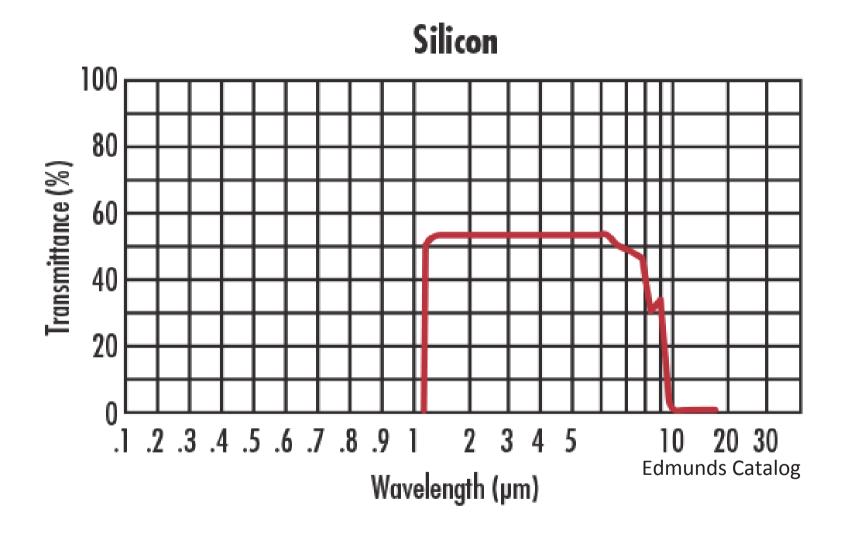


Today:

What is the "infrared"? State of the art instruments today Requirements and where they come from GMTNIRS Spectrograph basics Immersion gratings- the "secret sauce"



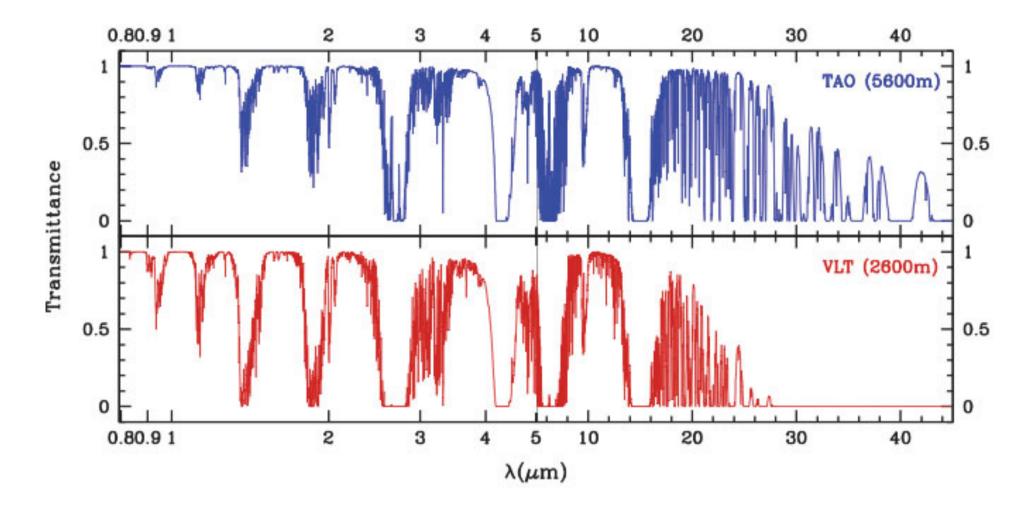
Why are silicon detectors so dominant?



Dominant IR options are HgCdTe or InSb, both with Si multiplexers.



#### The infrared divides into windows



University of Tokyo Observatory



#### Background from Sky

 $N = B_{v} \Delta v A \Omega_{pix} \varepsilon / (hv)$ 

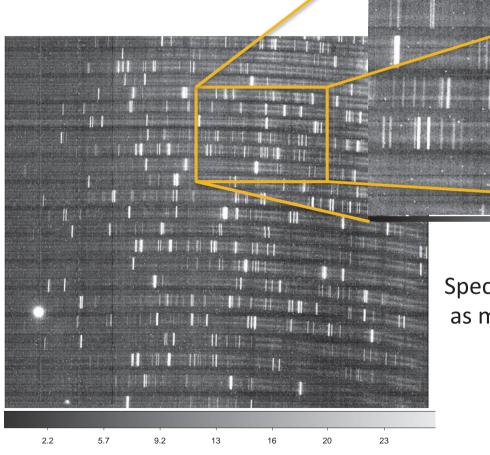
#### GMTNIRS Background Counts e/s/pixel J H K L 5e-10 6e-6 1e-2 8

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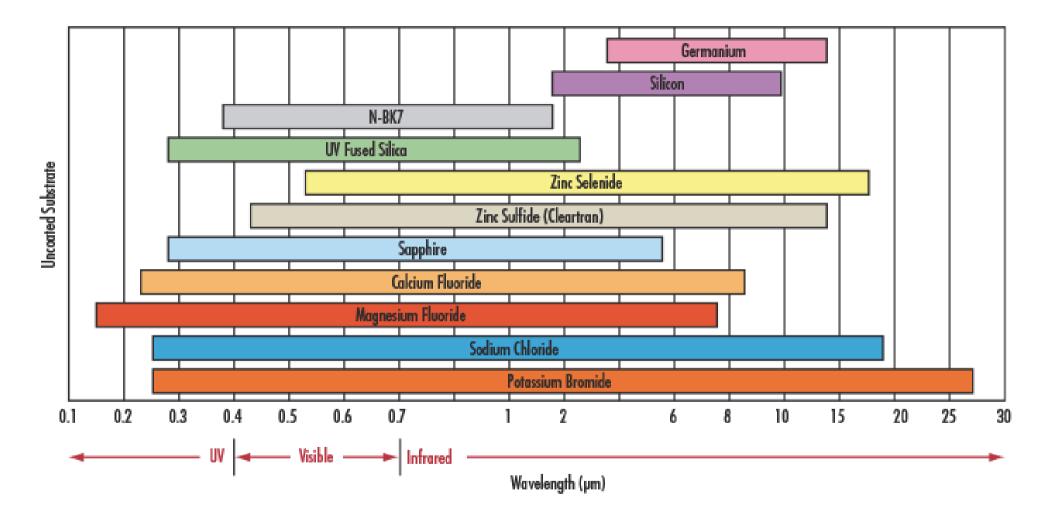


# OH Night sky in the H band (Chan Park et al.)



Spectrum of the entire H band as measured with IGRINS





Edmunds.com



State of the art instruments today Resolving power and resolution  $R=\lambda/\Delta\lambda$  Resolution =  $\Delta\lambda$  or in some cases  $\Delta V$ Coverage angstroms, nm,  $\mu$ m –take your pick. Note Window widths: H and K are each ~400 nm or about  $\lambda/5$ 

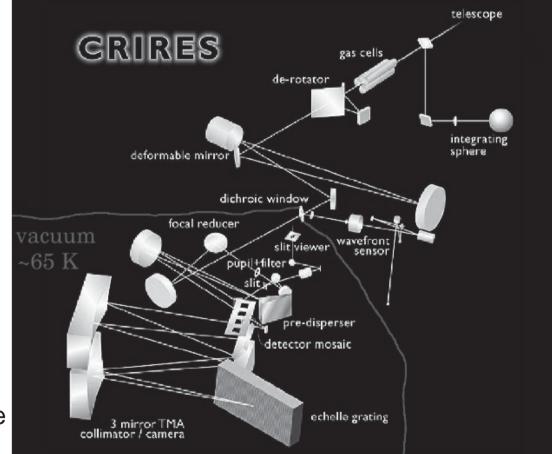


First-generation high resolution spectrographs for the IR

Small wavelength coverage: CSHELL:  $\lambda/200$ Phoenix:  $\lambda/100$ CRIRES:  $\lambda/70$ 

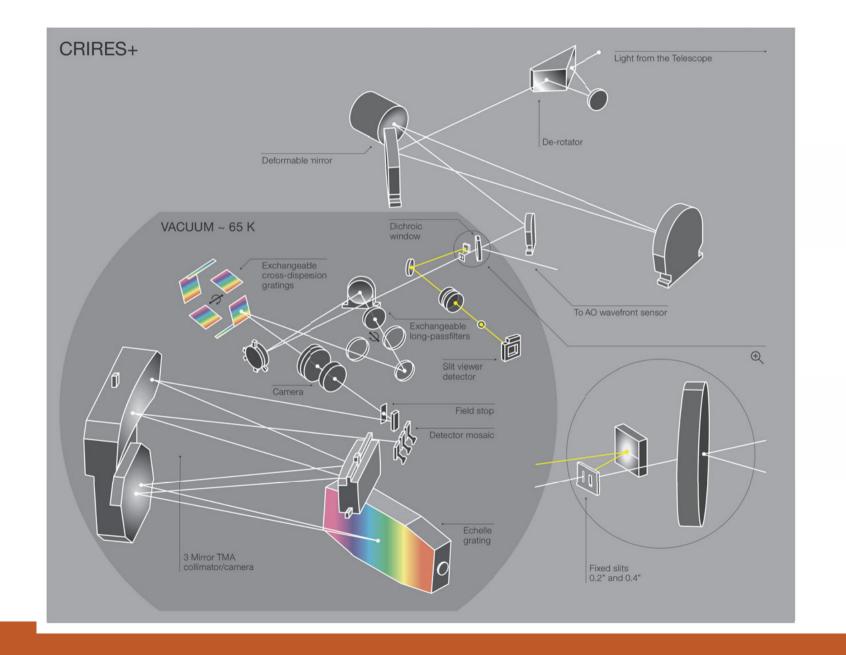
Small coverage means lots of moving parts.

Many surfaces. Hard to get decent slit sizes at high resolution. Instruments already quite large and so not too stiff.

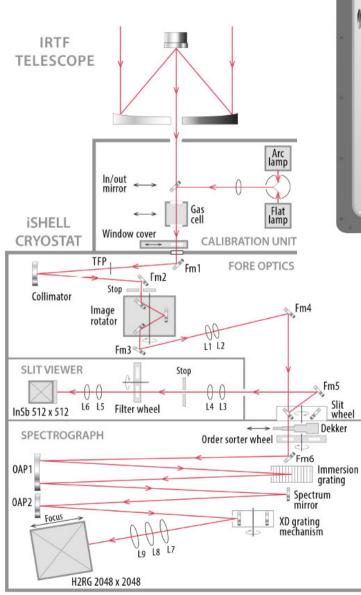


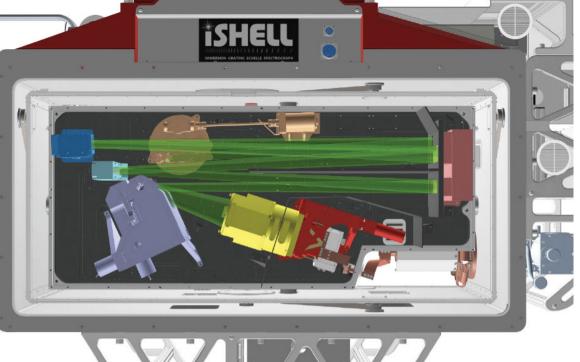


## VLT CRIRES+ - Full coverage of K









#### iSHELL on the NASA IRTF

#### R=70,000 1-5 $\mu$ m in snippets



**Resolving power** Spectral coverage Detector, Background Throughput Sampling Slit length Spatial resolution, Spectral stability Stability overall



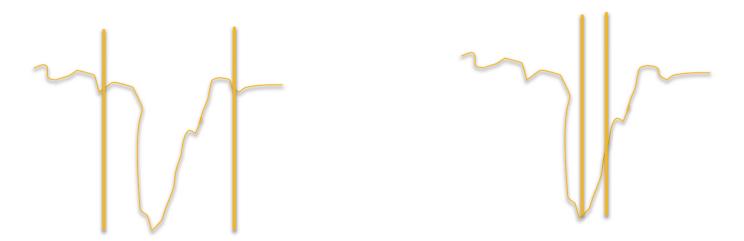
Resolving power

Science driven and technology limited.

 $R_{\text{Diff}}$ = 2Wtan $\beta/\lambda$ 

You are limited by the size of the grating and the need to let light into your system. Let's look at the S/N aspects.

Let's fix the number of pixels per resolution element:





#### How does S/N scale with R?

		Short l (low bknd)	Long I (high bknd)
Low S/N (=>detector noise limit when background is low)	Continuum	1/R	1/SQRT(R)
	Unresolved	Independent of R	SQRT(R)
	Resolved	1/R	1/SQRT(R)
	Emission unresolved	Independent of R	SQRT(R)
High S/N	Continuum	1/SQRT(R)	1/SQRT(R)
	Unresolved	SQRT(R)	SQRT(R)
	Resolved	1/SQRT(R)	1/SQRT(R)
	Emission unresolved	Independent of R	SQRT(R)



#### Which regimes matter?

		Short l (low bknd)	Long I (high bknd)
Low S/N (=>detector noise limit when background is low)	Continuum	1/R	1/SQRT(R)
	Unresolved	Independent of R	SQRT(R)
	Resolved	1/R	1/SQRT(R)
	Emission unresolved	Independent of R	SQRT(R)
High S/N	Continuum	1/SQRT(R)	1/SQRT(R)
	Unresolved	SQRT(R)	SQRT(R)
	Resolved	1/SQRT(R)	1/SQRT(R)
	Emission unresolved	Independent of R	SQRT(R)

#### What if you bin?



#### Spectral coverage: Which wavelengths?

Short IR Thermal IR Mid-IR

#### Spectral coverage: Grasp

Enough to cover line of interest Enough to cover group of lines of interest Enough to cover band or bands

Pan-spectral coverage enables a new mode of science for the infrared



Detector

- Quantum efficiency: As high as possible
- Read noise: Below reasonable backgrounds or as low as possible Dark current: Below reasonable backgrounds or less than read noise in reasonable exposure times.
- Wavelength coverage: to as long as necessary but not longer Readout speed: Not too slow (Fowler sampling)
- Size: Big
- Pixel Pitch: People may disagree but not small at long  $\lambda$ Memory and cross-talk: As little as possible Well depth: Big enough to allow high S/N per exposure Linearity: Good over full range or calibratable



Internal background: Basically a requirement on how cold your internal optics need to be.

Has to be less than dark current plus external background at the shortest wavelength you observe with a given channel.

Why do I phrase it that way?

#### **External Background**

Good site As few external surfaces as possible Cut off where atmosphere or detector does Cold stop- Must rotate The trade for AO- more surfaces vs. smaller  $A\Omega$ Slit losses vs. Strehl



Throughput: Obviously more is better. This is particularly true at short wavelengths where you will be source noise or detector noise limited. It matters somewhat less at longer wavelengths where cold throughput affects signal and background both-This tells you which channels to concentrate on.



Sampling: A combined real estate and noise question. You want to sample enough that your slit or diffraction are limiting R but oversampling takes up pixels you could use for slit length or spectral coverage and more pixels/resel means higher read noise to overcome.

For high resolution IR spectroscopy, the sweet spot lies between 2.5 and 4 pixels per resolution element.



Slit length: Needs to be long enough to allow sky subtraction by switching along the slit while always having the object in it. =>Depends on seeing or AO.

Planetary astronomy and ISM work want to drive you to longer slits.



#### Spatial resolution: Why?

ISM, Planets, disks, binaries, confused regions, exoplanets

Limiting Factors: Seeing- go to AO Diffraction Pixel size Image wander



Spectral stability: Depends on application. Achievable by mechanics, calibration, or both.

Stability overall: Helps data reduction, calibration. Allows for higher limiting S/N

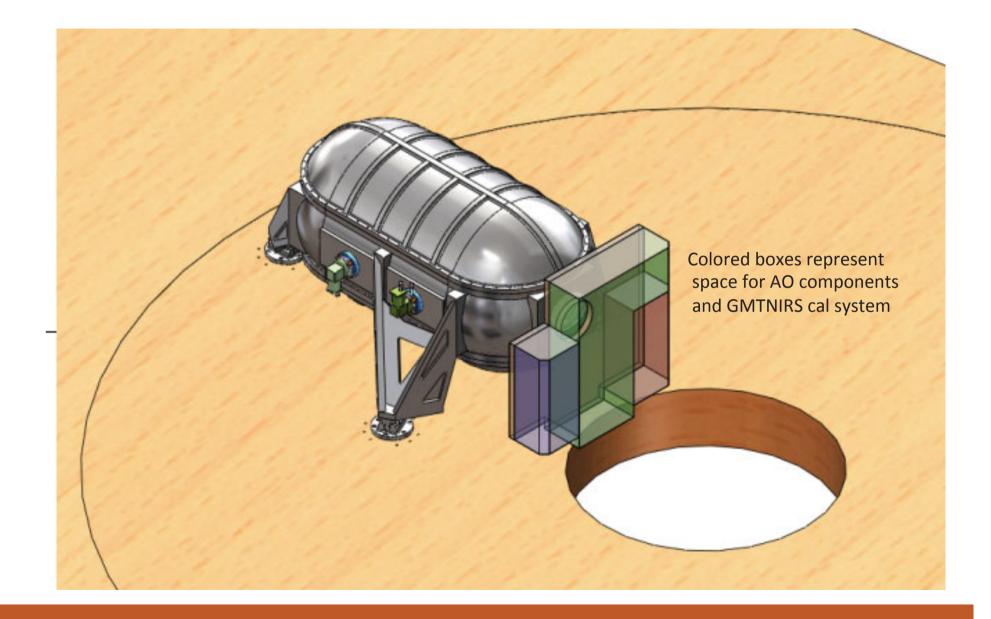


-Covers all spectrum transmitted by the atmosphere from 1.07  $\mu m$  to 5.3  $\mu m$  in a single exposure.

- -Resolving power 65,000 at  $\lambda$ < 2.5  $\mu$ m, 85,000 at  $\lambda$ >2.9  $\mu$ m.
- 65 mas slit, 1.3" long
- No moving parts after slit
- High throughput leads to high sensitivity
- K-band slit viewing camera
- Automated, near publication quality data pipeline



## GMTNIRS on rotating instrument platform



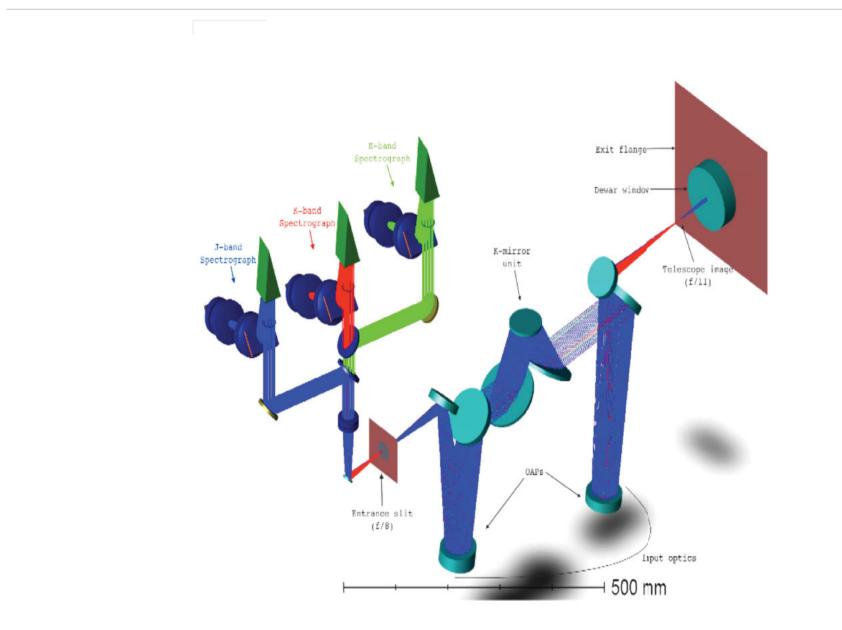




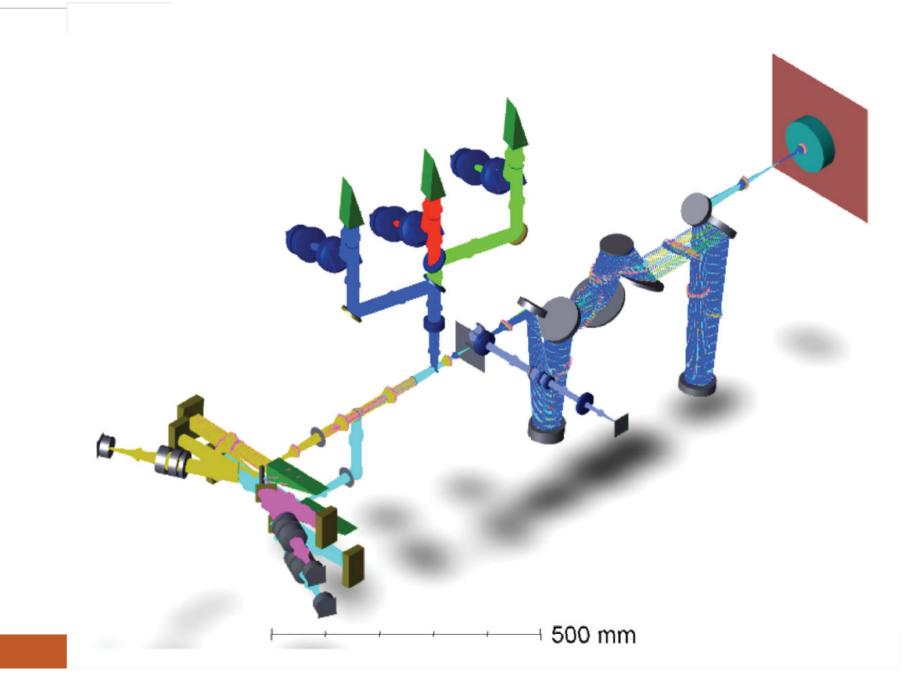
GMTNIRS sensitivity calculated conservatively: Average position along the blaze of both the echelle and the cross-disperser, 6x600 second exposures, image size = slit size.

GMTNIRS Sensitivity (JHK: R=50,000, LM: R=100,000) 1,000 In J,H,K, GMTNIRS can take a S/N> 100, R=50,000 \_ M 100 spectrum of basically anything S/N in the 2MASS catalog in under 10 an hour 8 9 10 11 12 13 14 15 16 17 18 19 20 Target Vega Magnitude











Spectrograph basics: Spectrograph design without the grating equation:

 $R_{diff}$  = Nm #of diffraction limited resels/order = # of grooves Diffraction limited spot size at focal plane = f/ •  $\lambda$ Diffraction limited spatial and spectral spot sizes the same. Real resolving power:

$$R = R_{diff} \bullet (\theta_{diff}/\theta_{slit}) = Nm (\lambda/D) (\theta_{slit})^{-1}$$

(where in the right-hand version  $\theta$  is in radians) slit



Let's design a spectrograph that can cover the entire K band (2.0-2.5 $\mu$ m) in a single exposure at R=50,000 for use on GMT.

Givens:

Detector is 2048x2048 with 18  $\mu$ m pixels (standard Teledyne product) Diffraction limit of telescope = 0.08\*2.2 = 18 mas.

We need a slit of 70 mas (this is about  $4\lambda/D$ )

This protects against phase errors and image wander, allows long wavelength use of same slit.

Let's also decide that we want to sample with 4 pixels per slit width.



The 4 $\lambda$ /D slit implies that R<sub>diff</sub>=200,000 to get R=50,000

With 4 pixel sampling  $R_{perpix} = 4x50,000 = 200,000$ 

The two together mean that  $R_{perpixel}$  = 200,000 or  $\lambda/D$  i.e. one  $R_{diff}$  for each pixel.

This also means that each pixel is one diffraction limited spatial element, or 18 mas.



Because of the length of a free spectral range varies across the band, we cannot use the full detector width at 2.2 microns. Better to use about 80% or 1600 pixels.

So, N=1600 grooves which will make the 2.2 $\mu$ m order free spectral range cover 1600 pixels.

How long does the grating have to be?  $R_{diff}$  = 200,000. L= $\lambda R_{diff}$ /(2sin $\delta$ ) For reasonable values:

 $L = 2.2 \times 10^{-4} \cdot 2 \times 10^{5} / (1.8) = 24 \text{ cm}$ 



What is the groove pitch we need?

2.4x10<sup>-1</sup>/1600= 1.6x10<sup>-4</sup>m= 160  $\mu$ m => 6 grooves/mm THIS IS IMPOSSIBLY COARSE

What camera f/ number do we need? f/•  $\lambda$  needs to equal the pixel size of 18  $\mu$ m, so we need f/8. THIS IS PRETTY EASY.



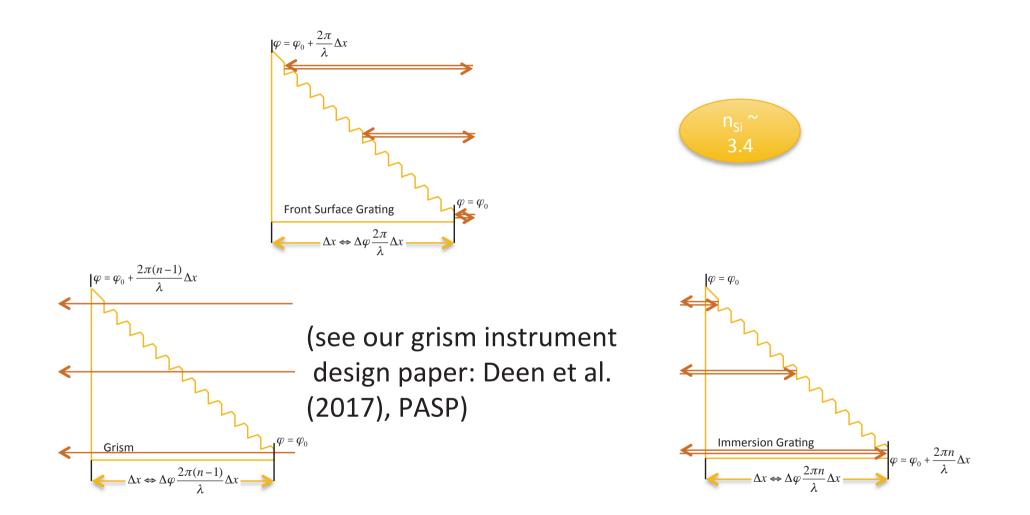
Now, the only question is the cross-dispersion.

The window is about  $\Delta\lambda/\lambda = 0.22$ , so we need 0.22\*50,000= 11,000individual resolution elements. Our central order has 400 elements, so we need 28 orders. If we spread these across 2000 pixels, we have 71 pixels per order for a mean separation of

0.018\*71= 1.3".

In practice, we will need to leave some space between orders so the slit length will be closer to 1". Good thing this is an AO instrument.







Resolving power is equal to phase delay across the beam measured in wavelengths divided by the ratio of the angular size of the slit to the diffraction limited angular scale.

 $\mathsf{R} = \lambda / \Delta \lambda = (2\mathsf{Lsin}\delta) / \lambda \bullet \theta_{\mathsf{diff}} / \theta_{\mathsf{slit}}$ 

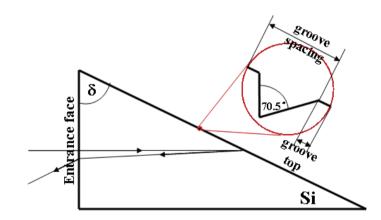
Where L is the grating length and  $\delta$  is the grating angle.

In immersion, the internal wavelength shrinks by a factor of n, the refractive index

$$m\lambda = n_G \sigma(\sin \alpha + \sin \beta)$$
$$R_{\max} = \frac{2n_G L \sin \delta}{\lambda}$$
$$\frac{d\beta}{d\lambda} = \frac{2n_G \tan \beta}{\lambda}$$



1. The "eureka moment"- Spring 1992 If you illuminate a grating from inside a transparent medium, the density of the medium leads the light to see the grating as "longer", that is, bigger by a factor of the refractive index, n, which is 3.4 for silicon



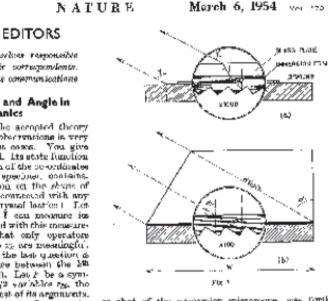


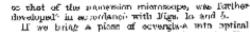
### 2.The "oh" moment – discovering a US patent from 1984

United States Patent [19]		atent Number:	4,475,792
Sica, Jr.	[45] Date of Patent:	Oct. 9, 1984	
54] HIGH RESOLUTION DIFFRACTION GRATING		OTHER PUBLICAT	
75] Inventor: Louis Sica, Jr., Alexandria, Va.	Aoyagi et al., Optics Communications, vol. 29, No. 3, Jun., 1979, pp. 253-255. Hugh L. Garvin, E. Garmire, S. Somekh, H. Stoll and A. Yariv, "Ion Beam Micromachining of Integrated Optics Components" Applied Optics, vol. 12, No. 3; Mar., 1973; pp. 455-459. E. G. Loewen, M. Neviere and D. Maystre "Grating Efficiency Theory as it Applies to Blazed and Holo-		
73] Assignce: The United States of America as represented by the Secretary of the Navy, Washington, D.C.			
[21] Appl. No.: 402,403			
[22] Filed: Jul. 27, 1982	graphic Gratings" Applied Optics, vol. 16, No. 10; Oct., 1977 pp. 2711-2721.		
[51]         Int. Cl. <sup>3</sup> G02B 5/18           [52]         U.S. Cl.         350/162.17; 350/162.22; 350/162.23; 356/305	Primary Examiner—F. L. Evans Attorney, Agent, or Firm—Robert F. Beers; William T. Ellis; Charles E. Krueger		
[58] Field of Search	[57]	ABSTRACT	
[56]         References Cited           U.S. PATENT DOCUMENTS         2,339,053           2,339,053         1/1944         Coleman           3,034,338         5/1962         Barnes et al.           3,054,338         5/1962         Barnes et al.           3,663,8424         4/1972         Elliott .           3,883,221         5/1957         Birgrod .           FOREIGN PATENT DOCUMENTS         FOREIGN PATENT DOCUMENTS	An optical prism, with an input surface, a reflecting surface and a base, and with a diffraction grating formed on the reflecting surface. The prism is designed so that if an input beam is incident normal to the input surface, then the efficiency of the order diffracted in the direction opposite the input beam is greater than a pre- determined value. If the prism is fabricated of a material with index of refraction n, then the resolving power of the grating is increased by n.		
312534 5/1929 United Kingdom . 706711 12/1979 U.S.S.R		8 Claims, 6 Drawing 1	Figures
	β <sub>1</sub> = θ <sub>0</sub>	10 16 (	



#### 3. The "hah" moment- finding an article from Nature, 1954





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#### and Angle in anics

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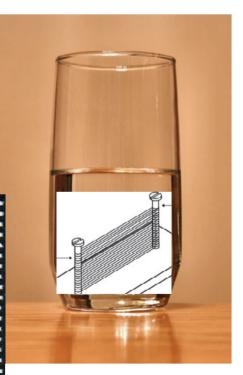
4. The "duh" moment: Gegenseitige Einwirkung der im Wasser und anderen brechenden Mitteln gebeugten Strahlen Joseph Fraunhofer 1823

The interaction of diffracted radiation in water and other refractive media.

Fraunhofer immersed his gratings in water, linseed oil, and turpentine. We immerse ours in solid silicon but the physics is the same.

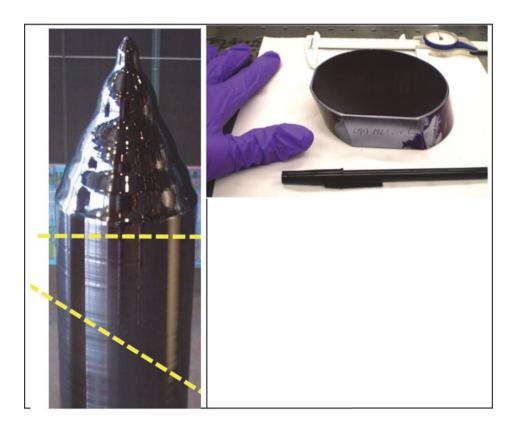








The thick substrates require a completely independent process chain from what wafers require.





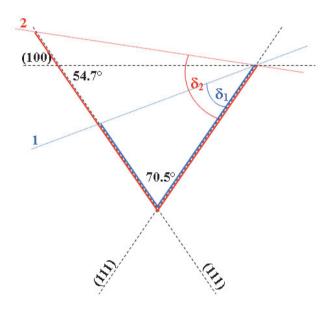
Material preparation requires coordination of:

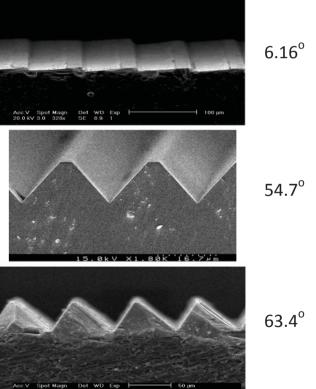
- 1. Float-zone silicon boule, resistivity of >10,000 ohm-cm.
- 2. X-ray crystallography to orient to 0.04°. Tilt and slice boule.
- 3. Chemically mechanically polished to  $\lambda/10$
- 4. LP-CVD silicon nitride 600 Å  $\pm$  5%



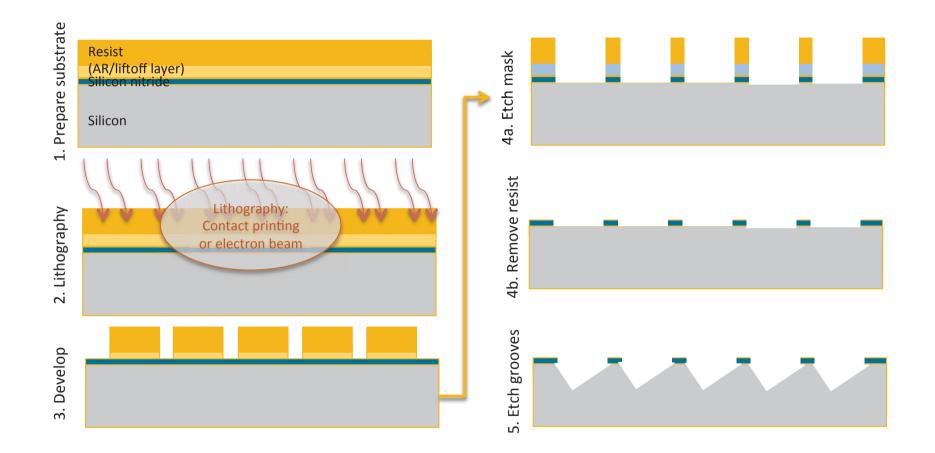


By correctly orienting the disk surface with respect to the silicon crystal planes, we can produce a grating with any blaze angle.

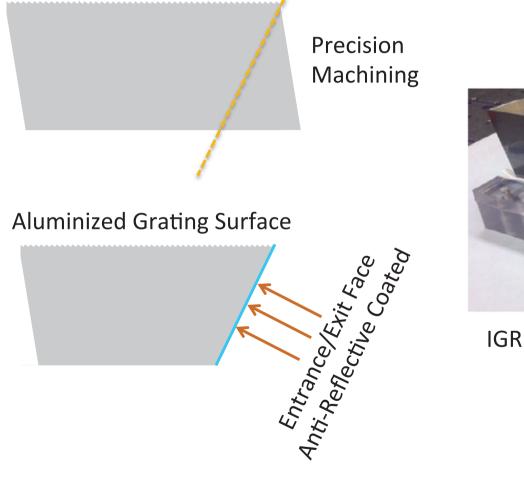


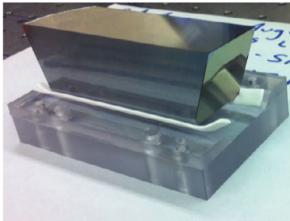












**IGRINS** immersion grating



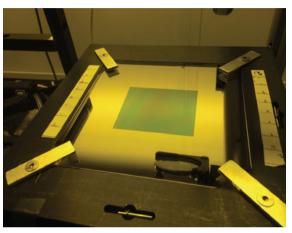
Test Front surface interferogram in Littrow at 633 nm (Equivalent to~2.15 μm	Specification $\lambda/4$ in immersion at operating wavelength		
in immersion) Laser spectrum	Ghosts < 2 x 10 <sup>-3</sup>	31 <u>18</u> m	
Imaging with scanning electron microscope (SEM)	Smooth surface, <<1% defect area	505012 mag B WD of 1901 HV - 100 um - 112312AM 503 102 mm ETD 30 1000 kV T1 - 11	CA1o 1.0 0.8 0.6 0.4 -
Efficiency in immersion over wavelength range	>= 75%	14 14	0.4 0.2 1400 1600 1800 2000 2200 2400

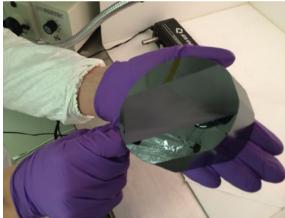


### **Contact Lithography**

### UTexas Grating Contact Aligner

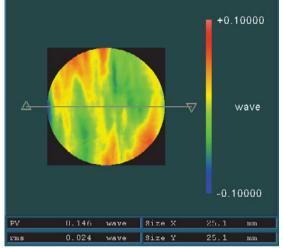




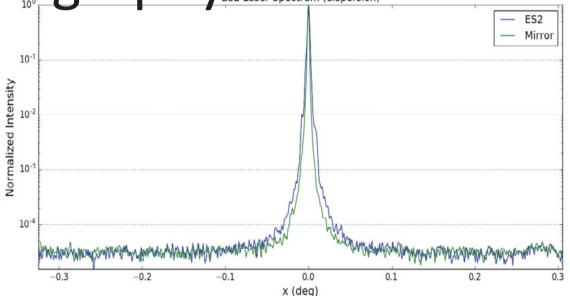




## Contact Lithography Results



- Front surface interferogram in Littrow at 71.6° for a 25 mm beam.
- Phase PV is 0.146 waves at 633 nm, corresponding to ~2.0 μm in immersion.

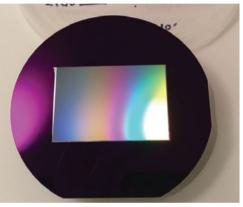


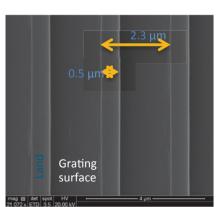
- High dynamic range, front surface monochromatic spectrum of same grating.
- Contact gratings can be as coarse as 5 l/mm



### JPL electron beam lithography tool



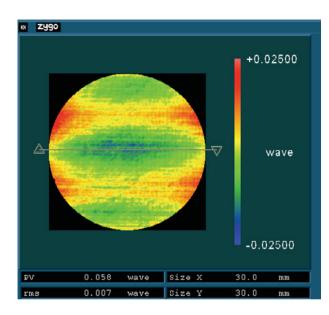




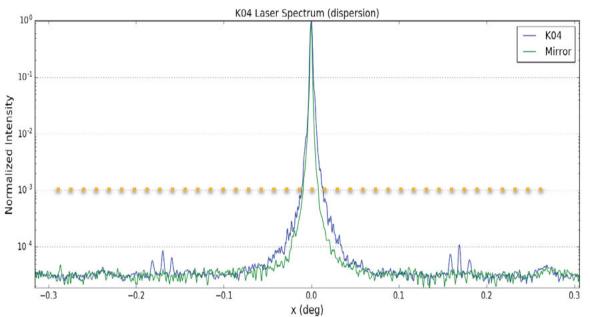
Electron beam lithography allows us to make very fine-pitch gratings of 1 µm or below.







- Front surface interferogram in Littrow at 18° for a 30 mm beam.
- Phase PV is 0.06 waves at 633 nm, corresponding to ~2.0 μm in immersion.



- High dynamic range, front surface monochromatic spectrum of same grating.
- Ghost specification is <10<sup>-3</sup> of the central peak; these are <10<sup>-4</sup>.



# **Degree of Difficulty**

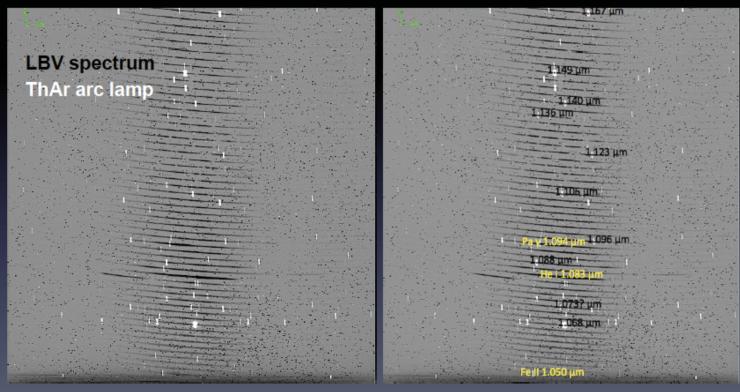
Degree of difficulty: scales with  $[(n-1) \sin \delta]/\lambda$  for grisms and with  $[2n\sin \delta]/\lambda$  for immersion gratings. Scales with length of grating divided by width of slit in units of diffraction limited width.

Forcast grism 1:  $\delta$ =11°  $\lambda$ = 30  $\mu$ m, DD=1 Other Forcast grisms: DD=1-10 JWST grisms: DD= 9 IG for IGRINS: DD=120 IG for GMTNIRS: DD= 250





### iSHELL: new J0 mode (for He I 1.083 μm)



At 80K Silicon transparent to full intensity at 1.065  $\mu$ m half intensity at 1.050  $\mu$ m, zero intensity at 1.035  $\mu$ m

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### **Deployed Technology**

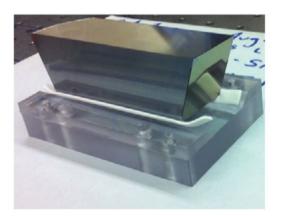
### Silicon grism flight parts:

- James Webb Space Telescope NIRCam
- FORCAST on SOFIA Telescope

Echelle immersion gratings:

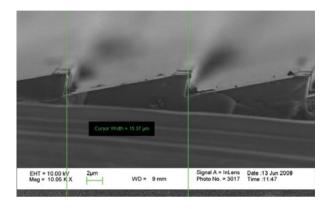
- Immersion GRating Infrared Spectrograph (IGRINS) installed at McDonald Observatory
- iShell for the NASA InfraRed Telescope Facility
- Planned GMT Near InfraRed Spectrograph (GMTNIRS) for the Giant Magellan Telescope.



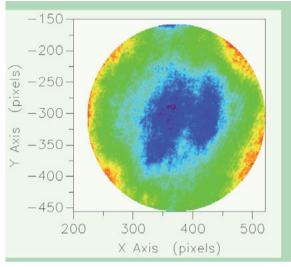


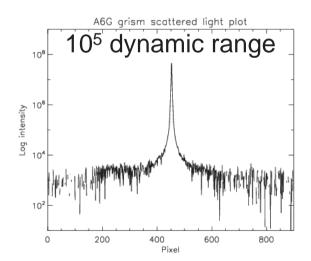


### JWST NIRCam Grisms



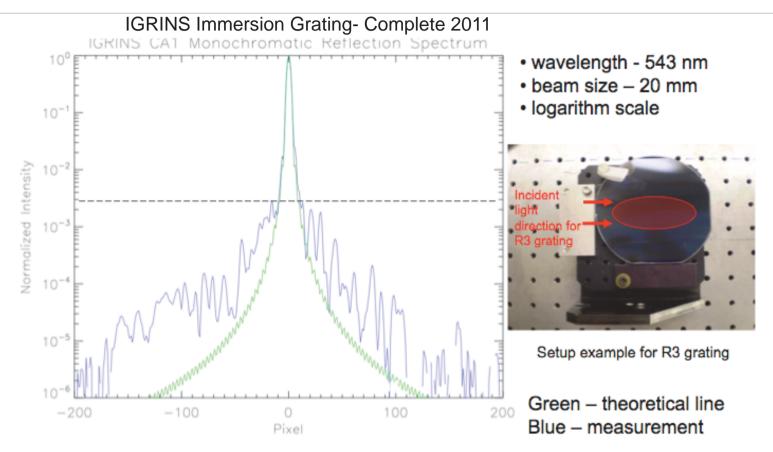
### $\lambda$ /100 rms grating surface





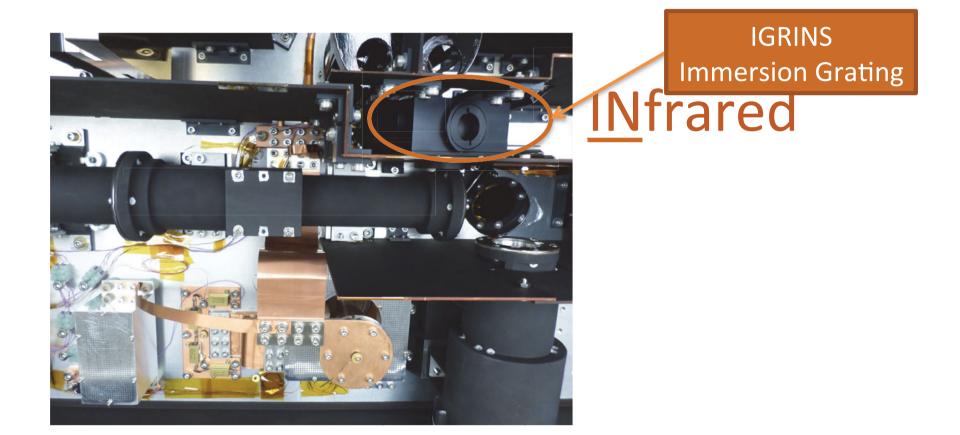




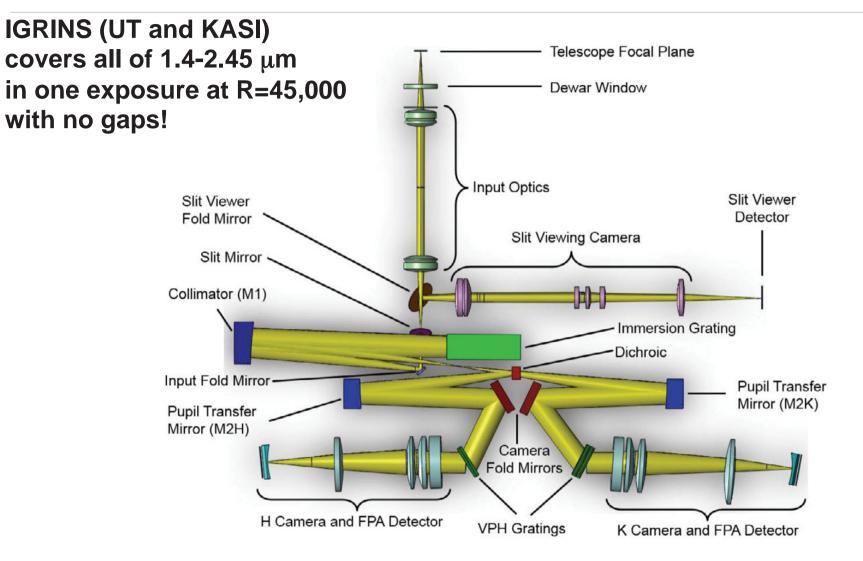


Surface is  $\lambda/6$  peak to valley. Efficiency ~80%



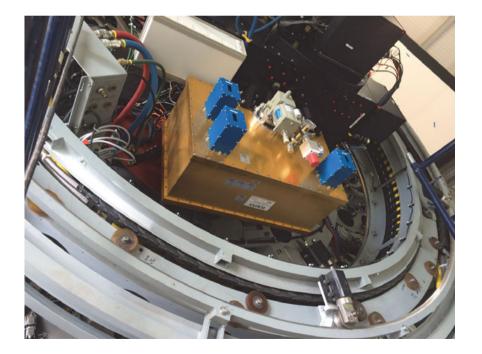


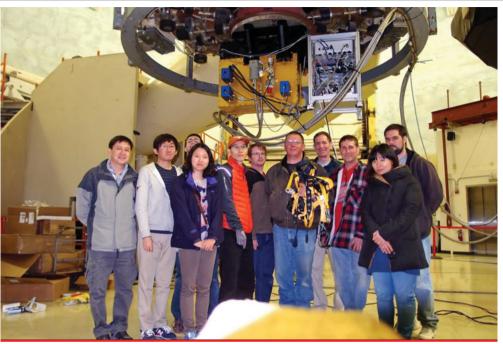






IGRINS on McDonald 2.7m 2014-2016, no cold moving parts, warm calibration system with moving mirrors





IGRINS on 4.3m DCT 2016- and McDonald 2.7m 2017-No moving parts, sky emission and absorption lines, white spot



IGRINS Spectrum of a star, not corrected for telluric absorption

