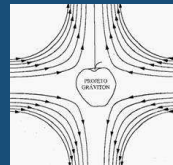




Gravitational Waves: Challenges and New Physics

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Instituto Nacional de Pesquisas Espaciais
Graviton Group



GRAVITON GROUP

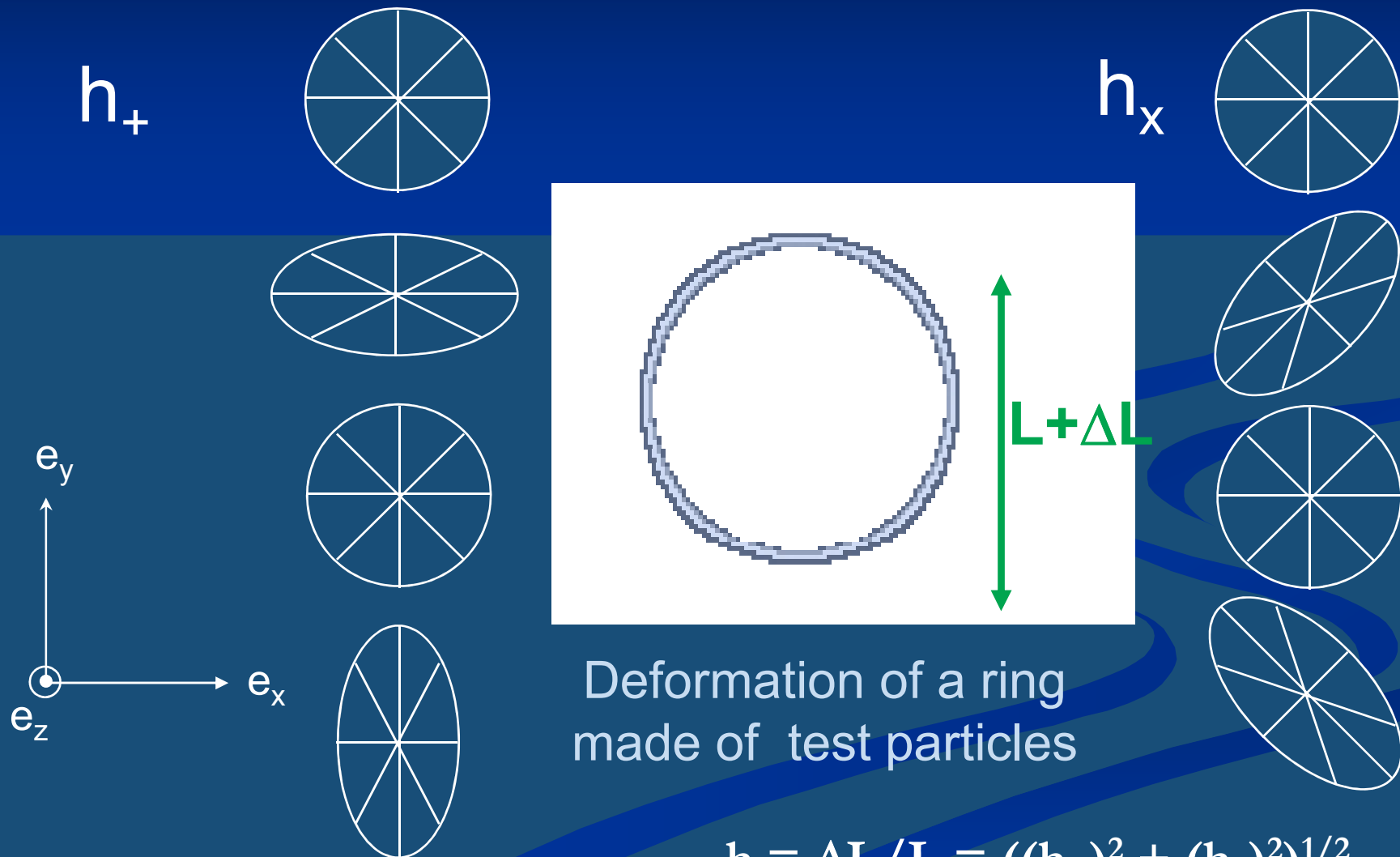
INPE

Campos do Jordão, April 26, 2009

Gravitational Waves

- 1916 - Einstein proved mathematically, from his theory of General Relativity, that gravitational waves should exist.
- These gravitational waves, caused by the accelerated movements of masses, would be distortions of the *spacetime* which would be travelling in the Universe with the speed of light.

Two polarizations “+” and “x”



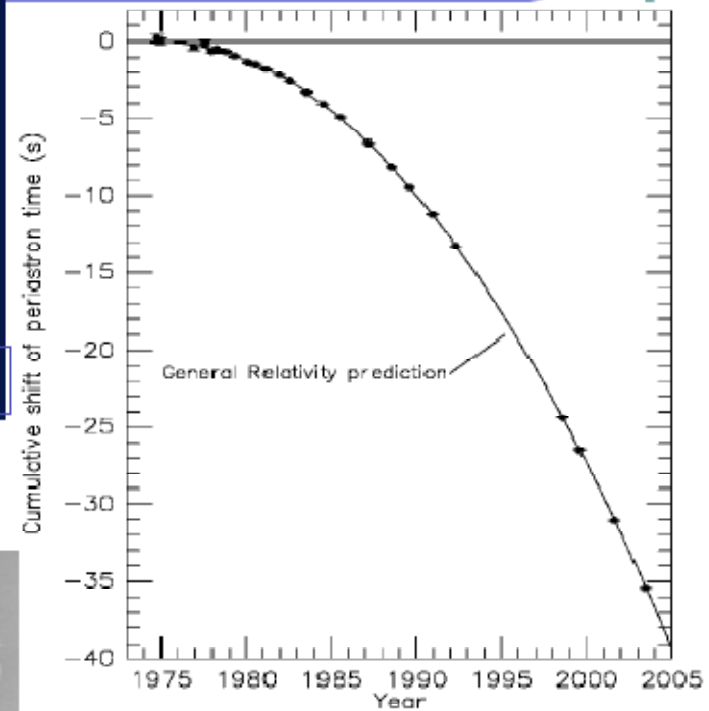
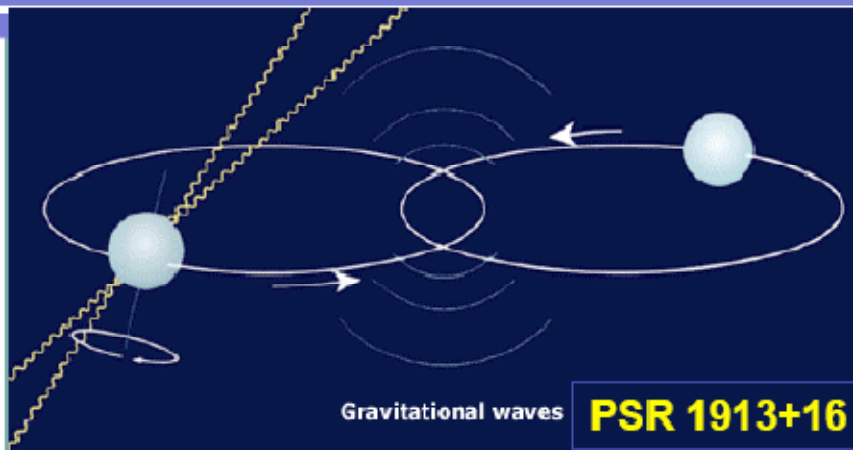
Deformation of a ring
made of test particles

$$h = \Delta L/L = ((h_+)^2 + (h_x)^2)^{1/2}$$

Gravitational Waves (GWs): do they really exist?

- Half century of theoretical debate (1916-1965);
- A direct detection hasn't been confirmed yet after four decades (1965-2009) of observational search.
- A good evidence (an indirect observation):
PSR 1913+16 (Taylor & Hulse 1974)
(System of two neutron stars with periods of 7h45min, in which one of them is a Pulsar)

First verification of GWs

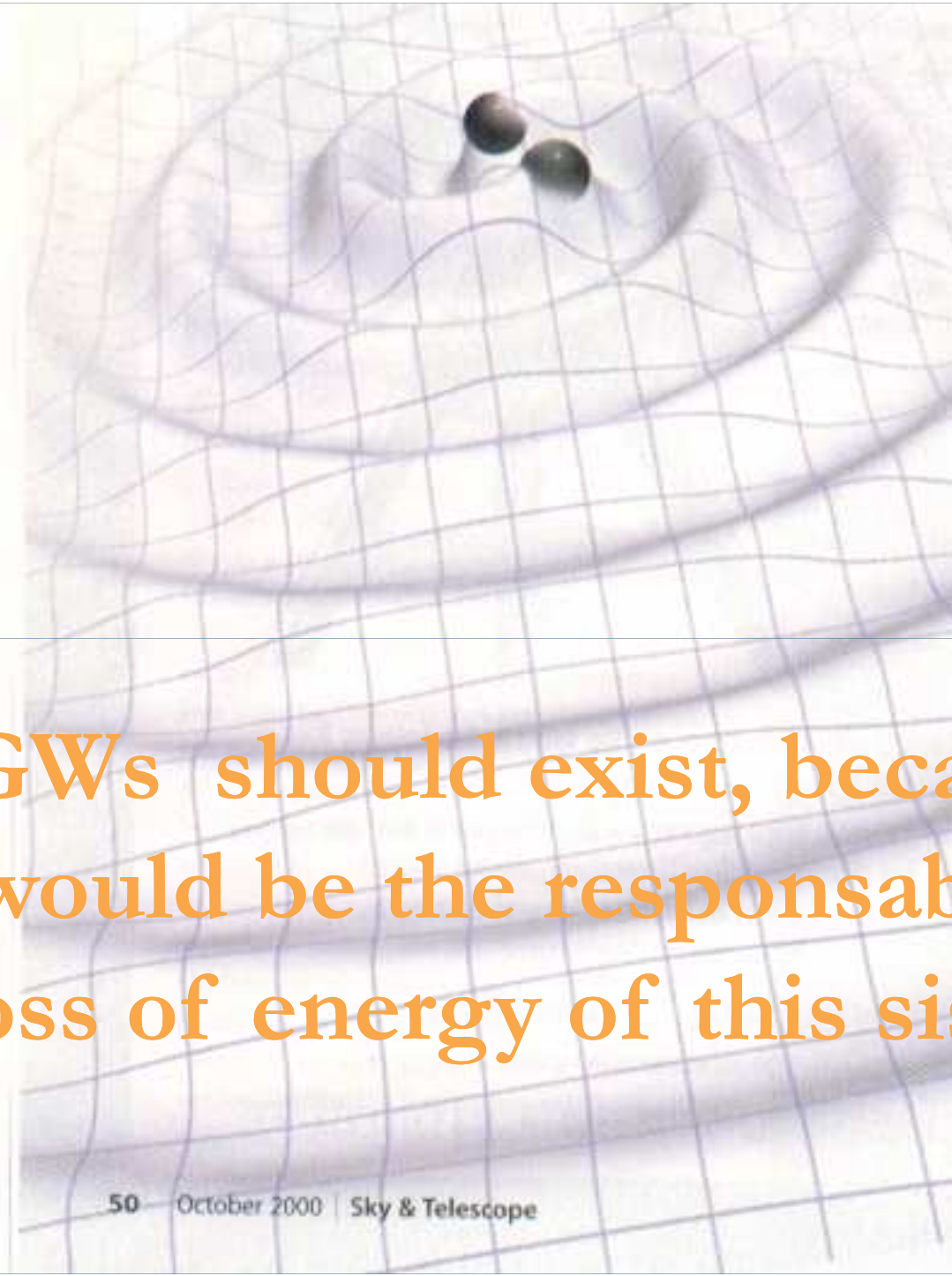


Nobel 1993

Hulse & Taylor



$$\frac{\dot{P}_{b,corrected}}{\dot{P}_{b,GR}} = 1.0013 \pm 0.0021$$



⇒ GWs should exist, because they would be the responsible for the loss of energy of this sistem.

$\nabla \cdot \mathbf{E}_g = -4\pi G\rho$ Newton's law (analogue to Gauss' law)

$\nabla \cdot \mathbf{B}_g = 0$ analogue to Gauss' law for magnetism

$\nabla \times \mathbf{E}_g = -\frac{1}{c} \frac{\partial \mathbf{B}_g}{\partial t}$ analogue to Faraday's law of induction

$\nabla \times \mathbf{B}_g = \frac{1}{c} \left(-4\pi G\mathbf{J} + \frac{\partial \mathbf{E}_g}{\partial t} \right)$ analogue to Ampère's circuital law

with Maxwell's correction

\mathbf{E}_g is the static gravitational field (or gravitoelectric field);

\mathbf{B}_g is the gravitomagnetic field;

ρ is mass density;

\mathbf{v} is the velocity of the mass flow generating the gravitomagnetic field;

$\mathbf{J} = \rho \mathbf{v}$ is mass current density;

G is the gravitational constant;

c is the speed of propagation of gravity (the speed of light in general relativity).

For a test particle of mass m and instantaneous velocity \mathbf{v}_m , the net (Lorentz) force acting on it due to a GEM field is described by the "Lorentz" force equation in gravitation:

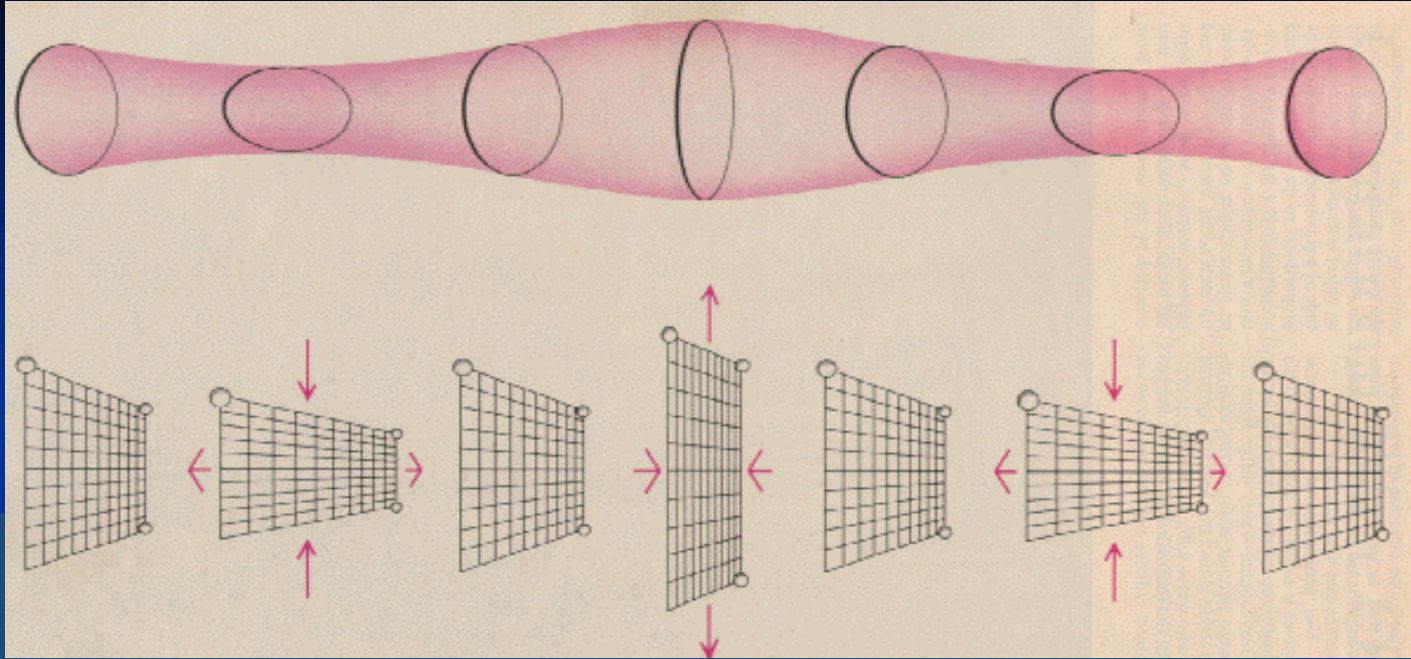
$$\mathbf{F}_m = m \left(\mathbf{E}_g + \frac{\mathbf{v}_m}{c} \times 2\mathbf{B}_g \right)$$

In vacuum:

$$\nabla \times \mathbf{E}_g = -\frac{1}{c} \frac{\partial \mathbf{B}_g}{\partial t}$$

$$\nabla \times \mathbf{B}_g = \frac{1}{c} \frac{\partial \mathbf{E}_g}{\partial t}$$

The changing gravitomagnetic field creates a changing gravitoelectric field through the analogue of Faraday's law for gravitation. That gravitoelectric field, in turn, creates a changing gravitomagnetic field through the analogue of Maxwell correction to Ampère's law for gravitation. This perpetual cycle allows these waves to propagate through space at velocity c .



**Why have they not
been detected yet ?**

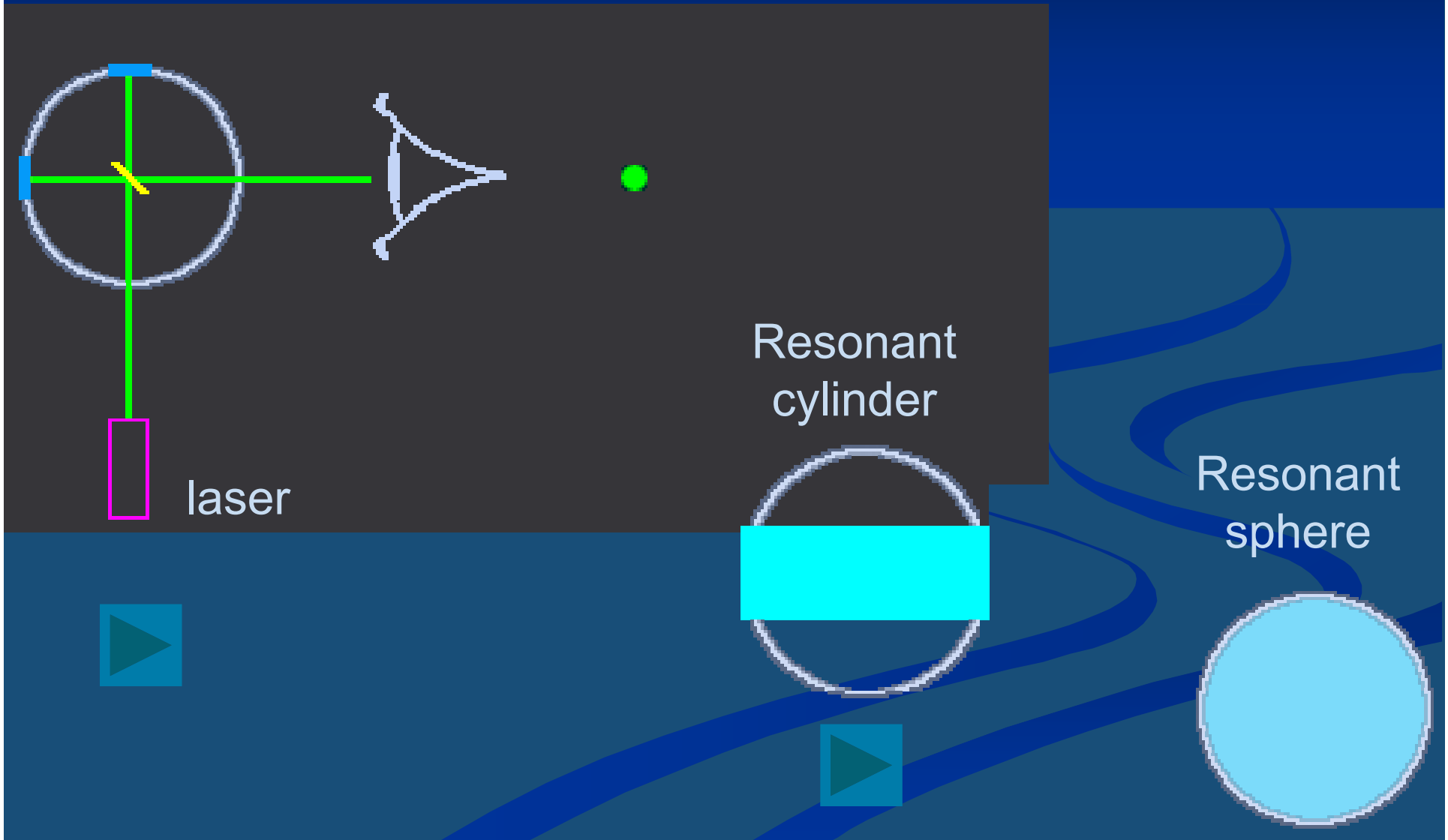
**Because they are hard
to be detected .**

**Why are they so hard
to be detected ?**

Gravitational Waves

- Two polarizations (“+” and “x”) according to General Relativity
 - Wave frequencies: 10^{-18} Hz to 10^{10} Hz
 - Very tiny amplitudes: $h = \Delta L/L < 10^{-19}$
 - Energies associated to $\Delta L \sim 10^{-19}$ m amplitudes:
for the Schenberg detector $\rightarrow \sim 6 \times 10^{-28}$ Joules
or 290 gravitons !
- This is 10^{-9} times the energy of a photon of light !

Gravitational Wave Detectors





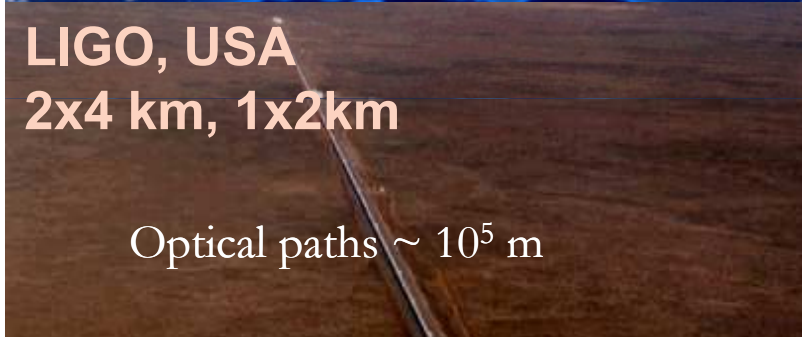
VIRGO, Italy
3 km



GEO600
Germany

tectors

by Arlette
de Waard



LIGO, USA
2x4 km, 1x2km

Optical paths $\sim 10^5$ m



TAMA300, Japan



AIGO, Australia



by Arlette de Waard

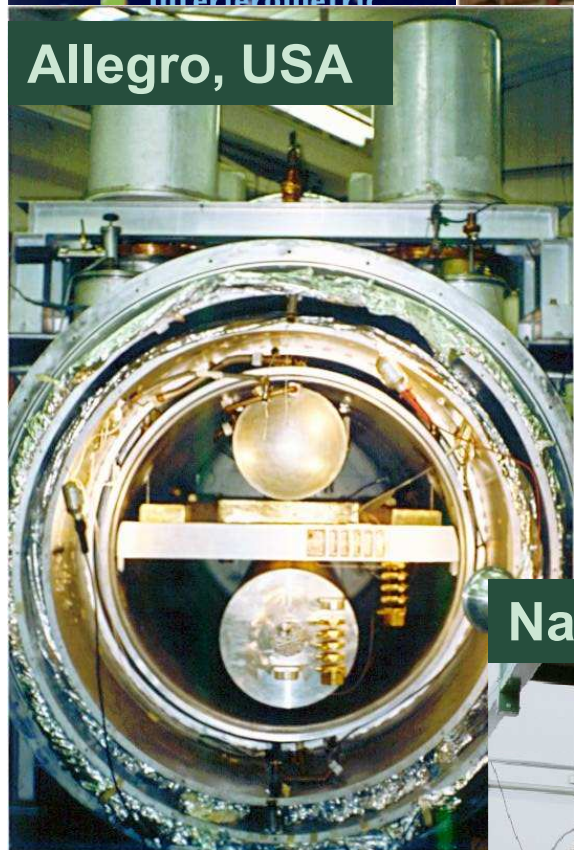
**Explorer
Switzerland**



Auriga, Italy



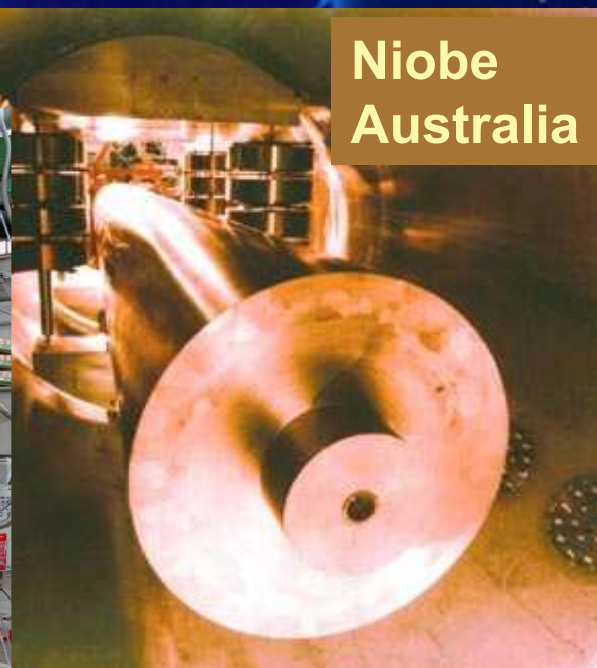
Allegro, USA



Nautilus, Italy



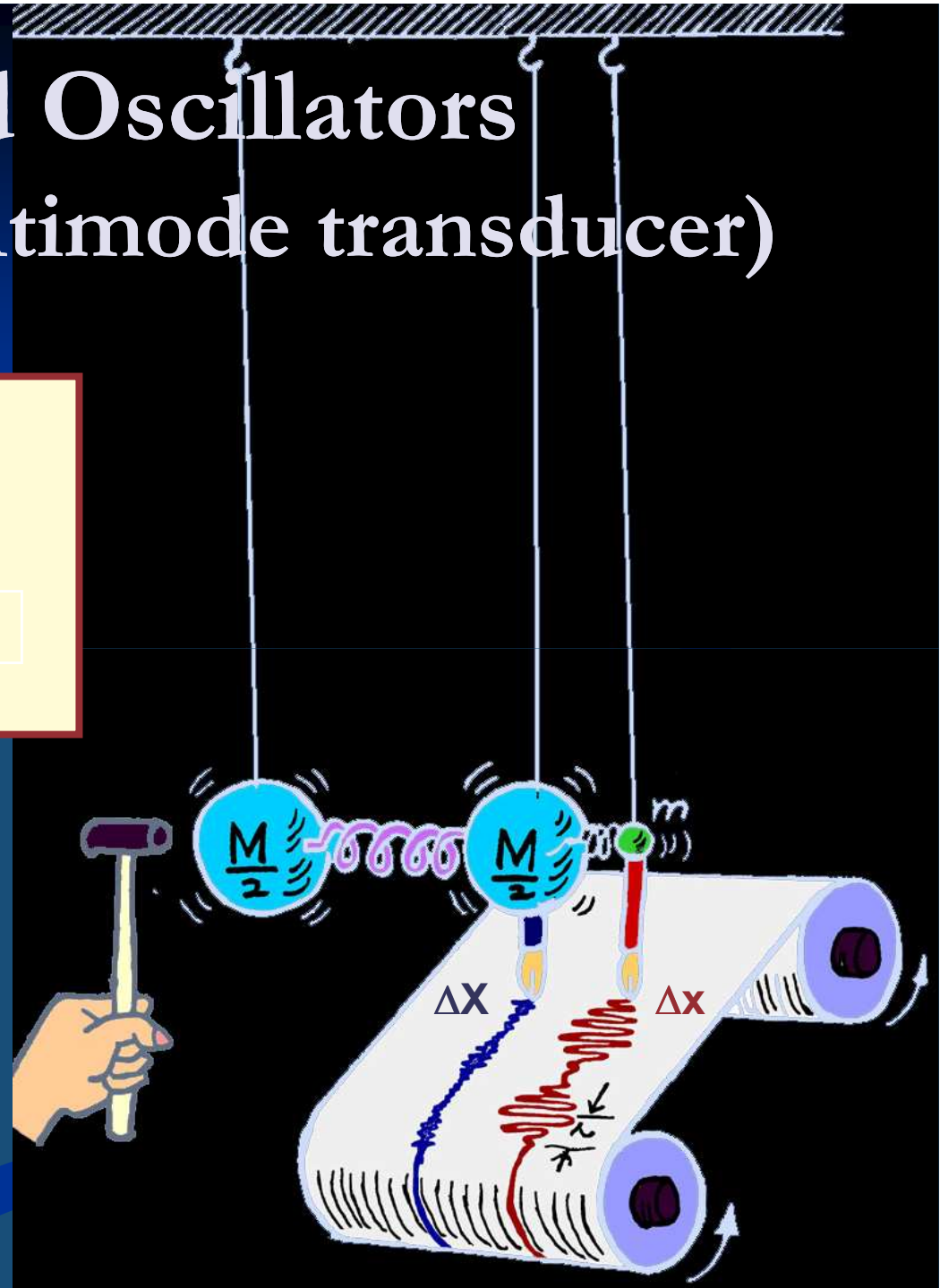
**Niobe
Australia**



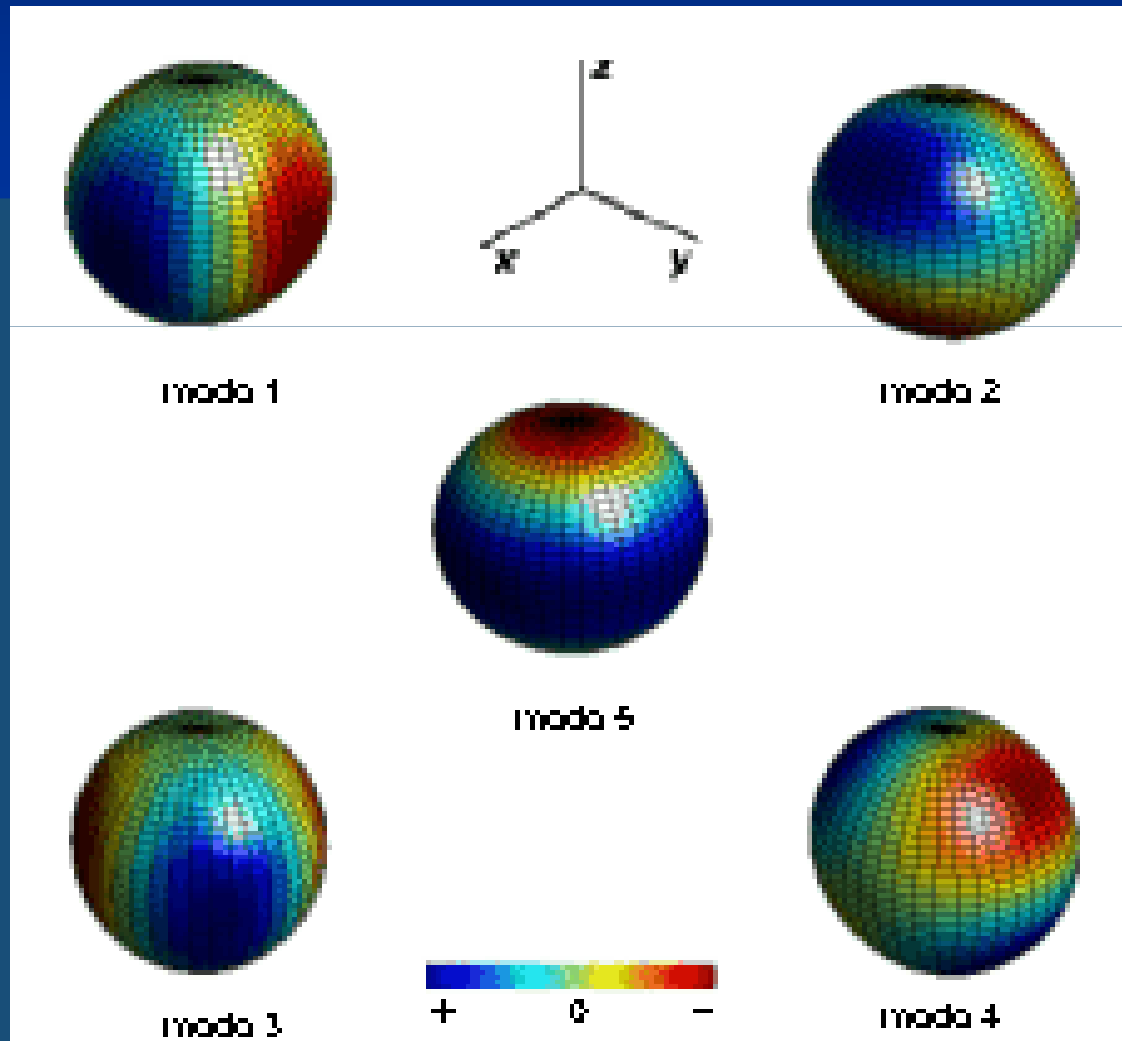
Coupled Oscillators (antenna + multimode transducer)

$$\Delta x = \sqrt{\frac{M}{m}} \Delta X$$

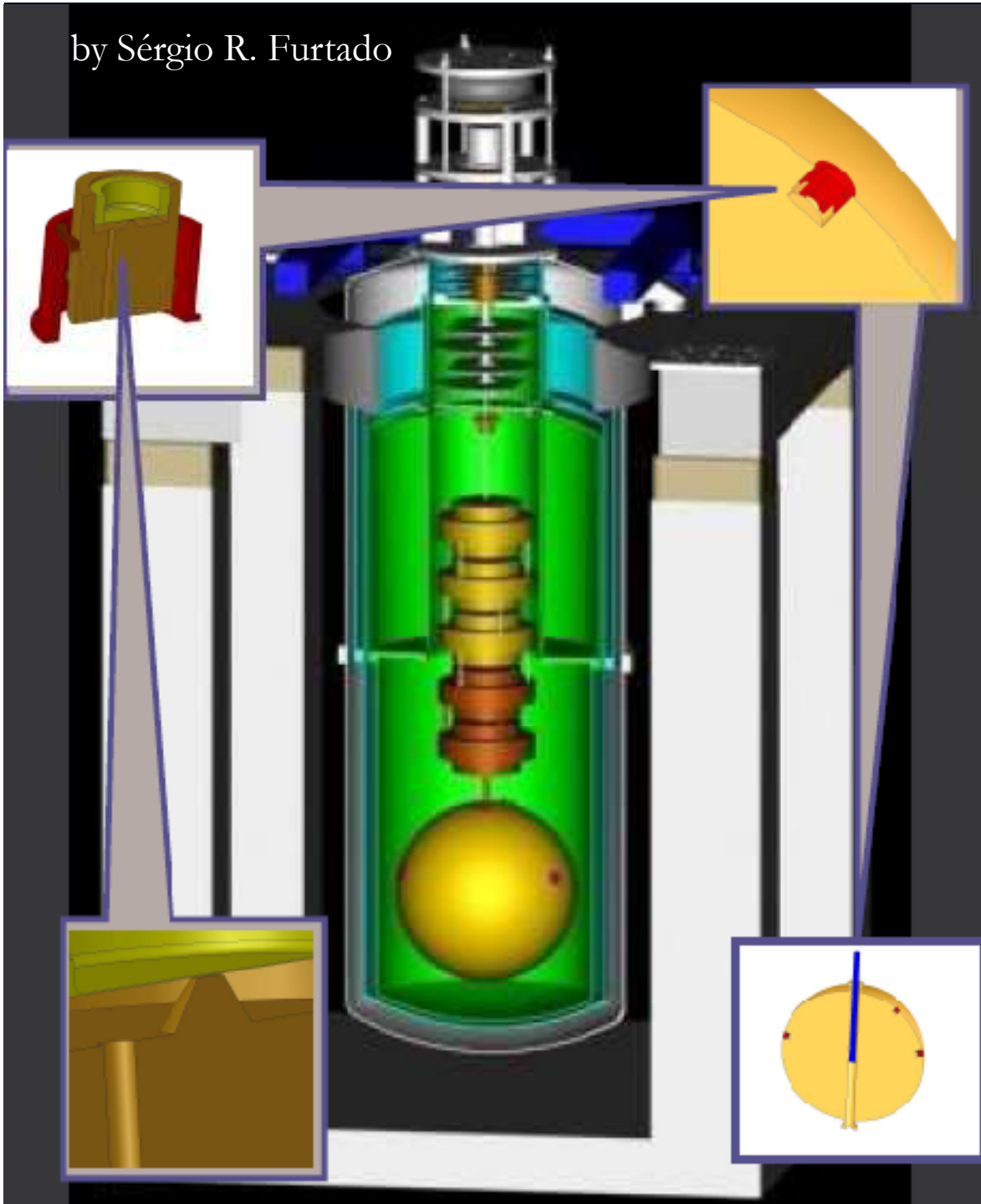
Ho Jung Paik



Quadrupole modes of the Mario Schenberg detector's sphere



by Sérgio R. Furtado



287 kg is the effective mass for each sphere's quadrupole mode;

$$\begin{aligned} & (5 \times 287 \text{ kg} = 1435 \text{ kg} \\ & \quad > 1150 \text{ kg} = M_{\text{sphere}}) \end{aligned}$$

**287 kg, 53 g and 10 mg
amplitude gain ~ 5300**

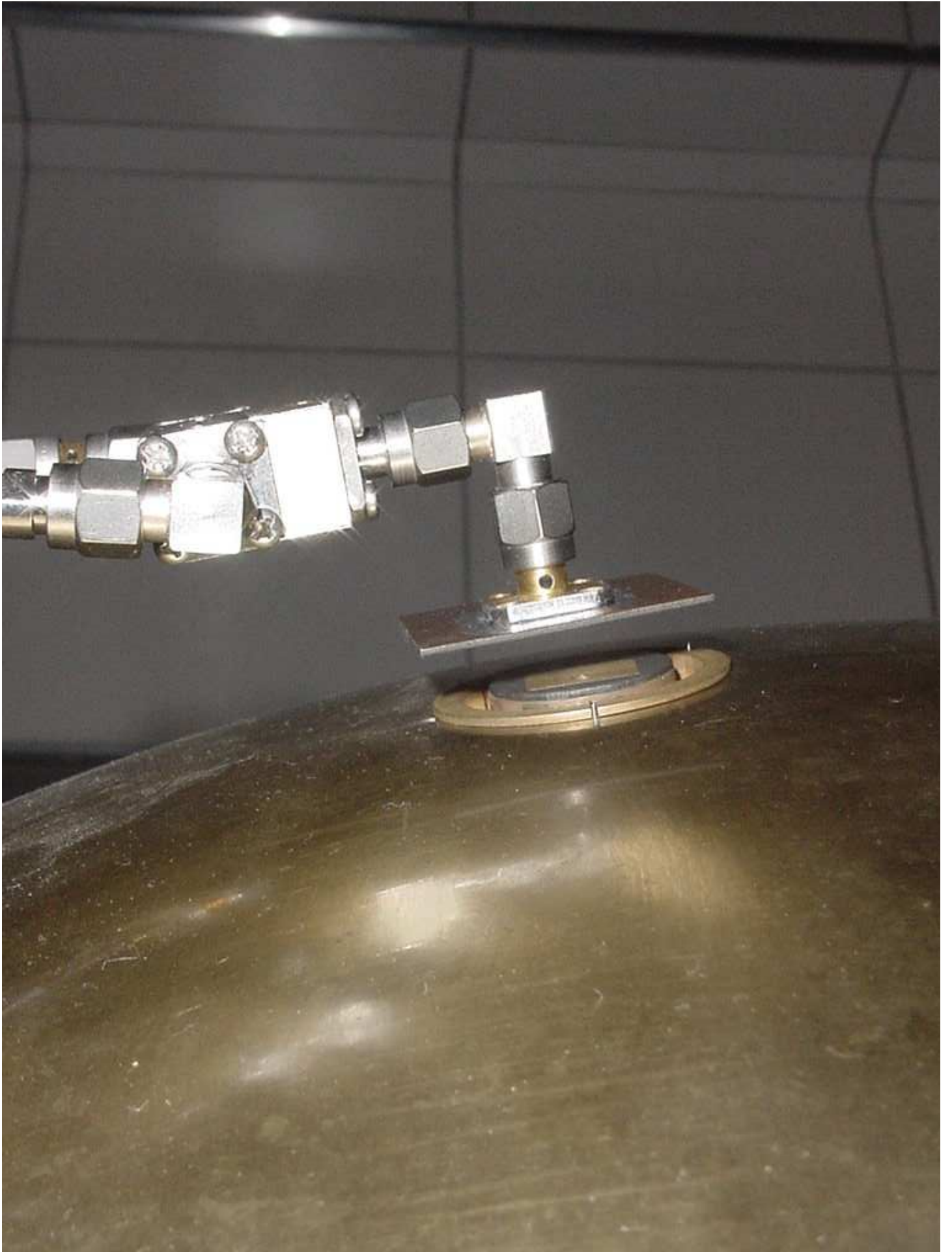
$$f_0 \sim 10 \text{ GHz}$$

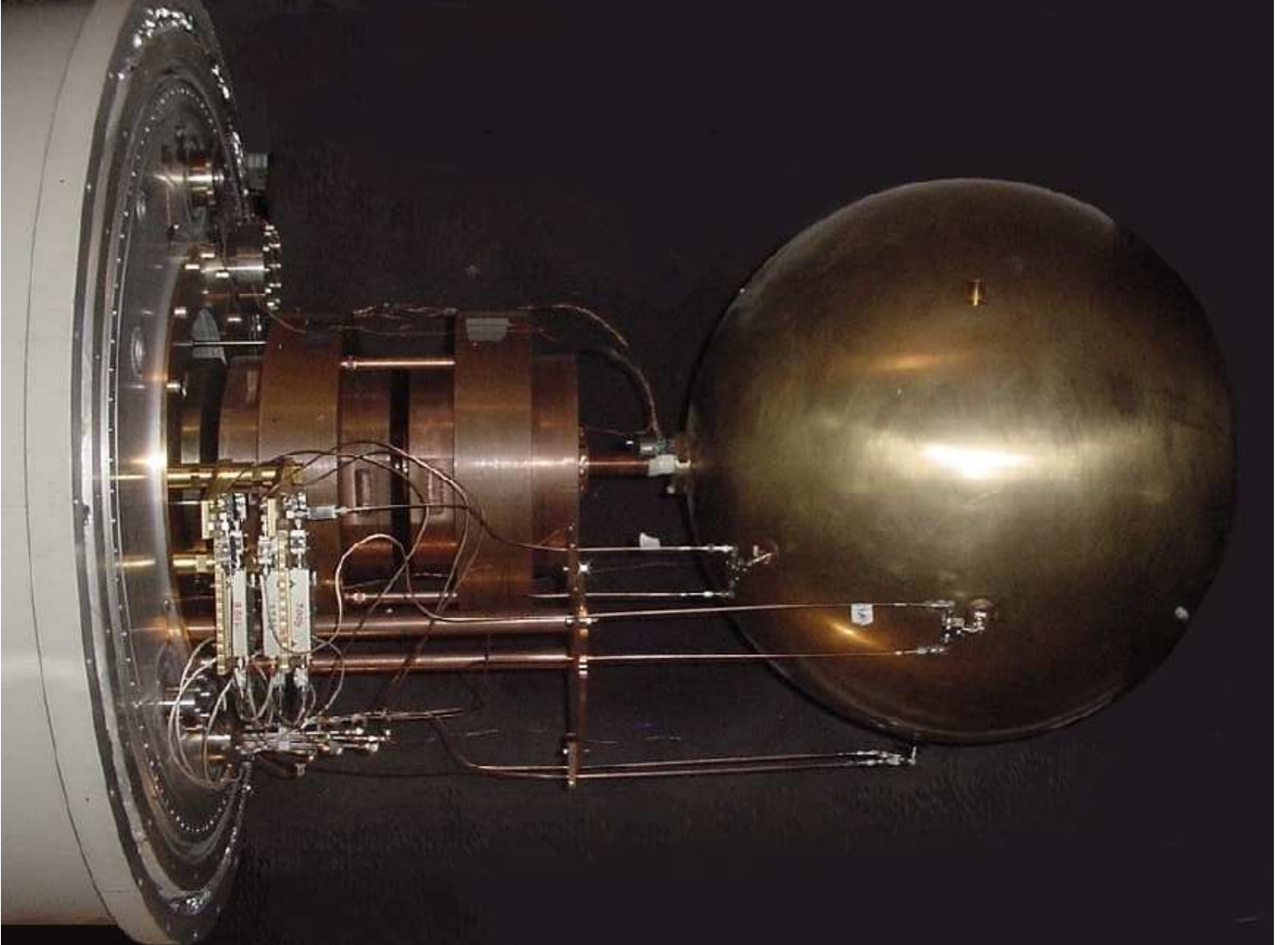
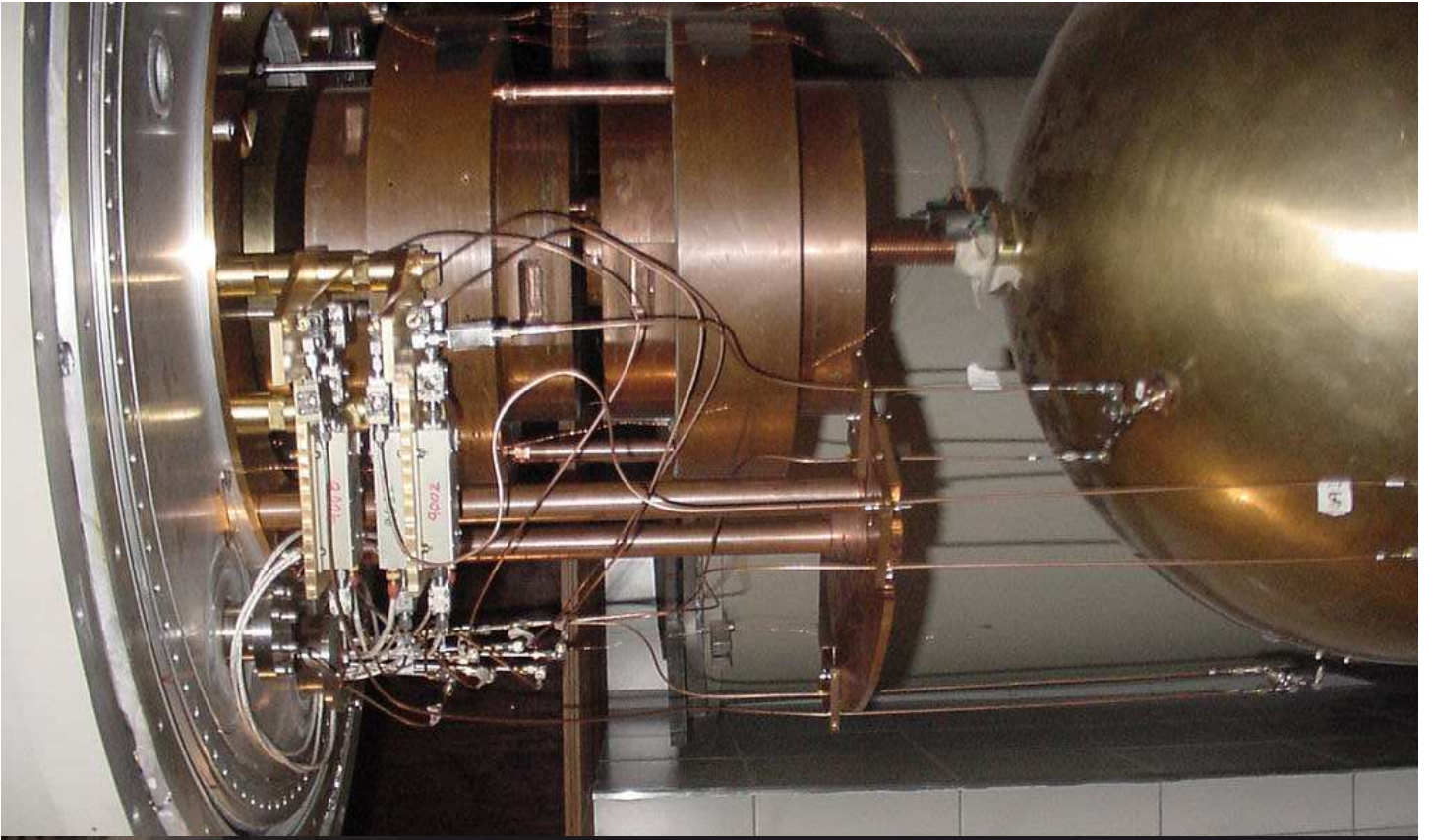
$$df/dx \sim 1 \text{ GHz / micron}$$

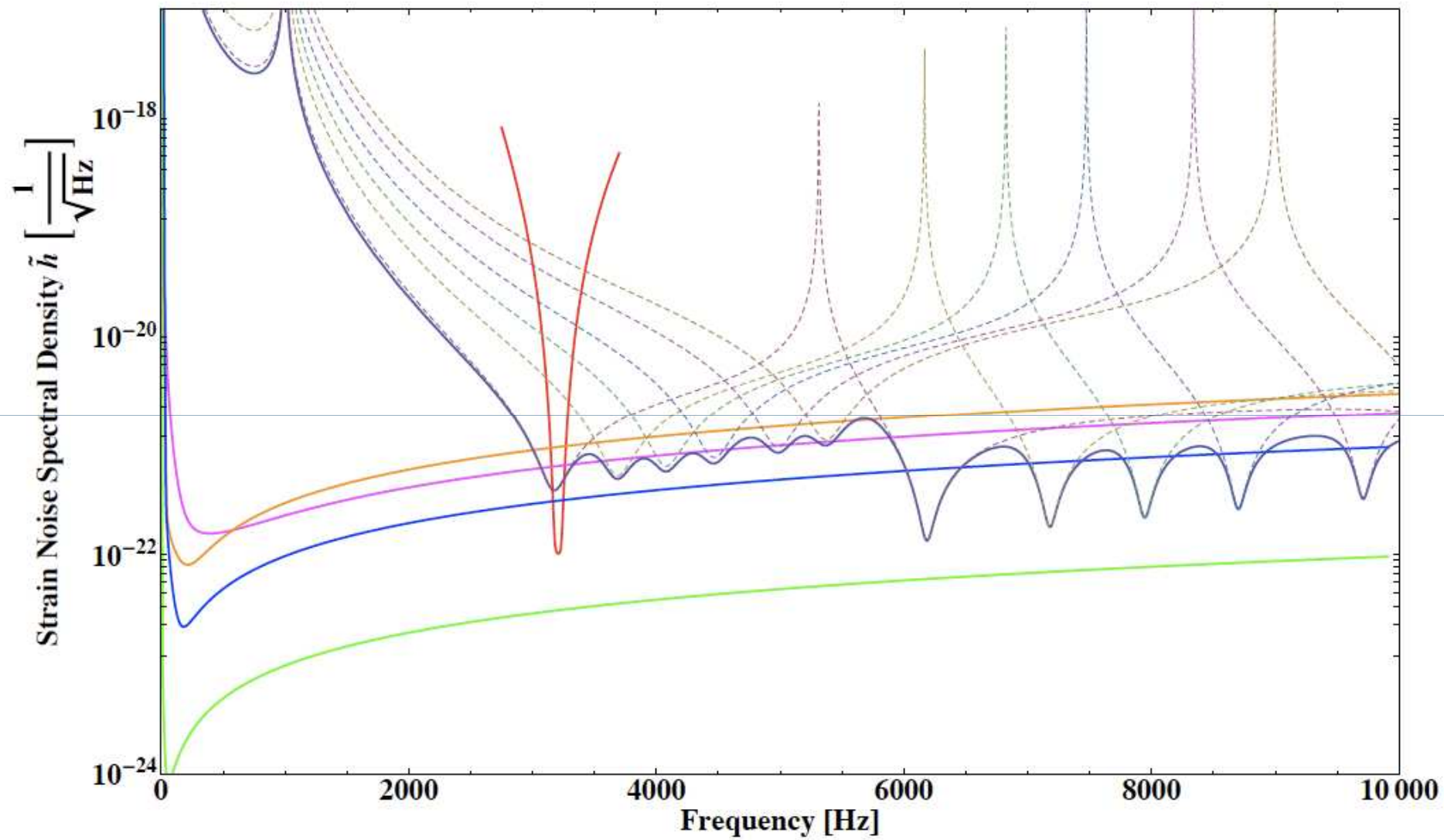
or

non-resonant

with $df/dx \sim 5 \text{ THz / micron}$







by Guilherme L. Pimentel

From the output of six 6 transducers
tuned to the quadrupole modes

$$\Psi(\theta, \phi, \omega) = \sum_i^5 a_i(\omega) \Psi_i(\theta, \phi)$$



spherical harmonics

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$

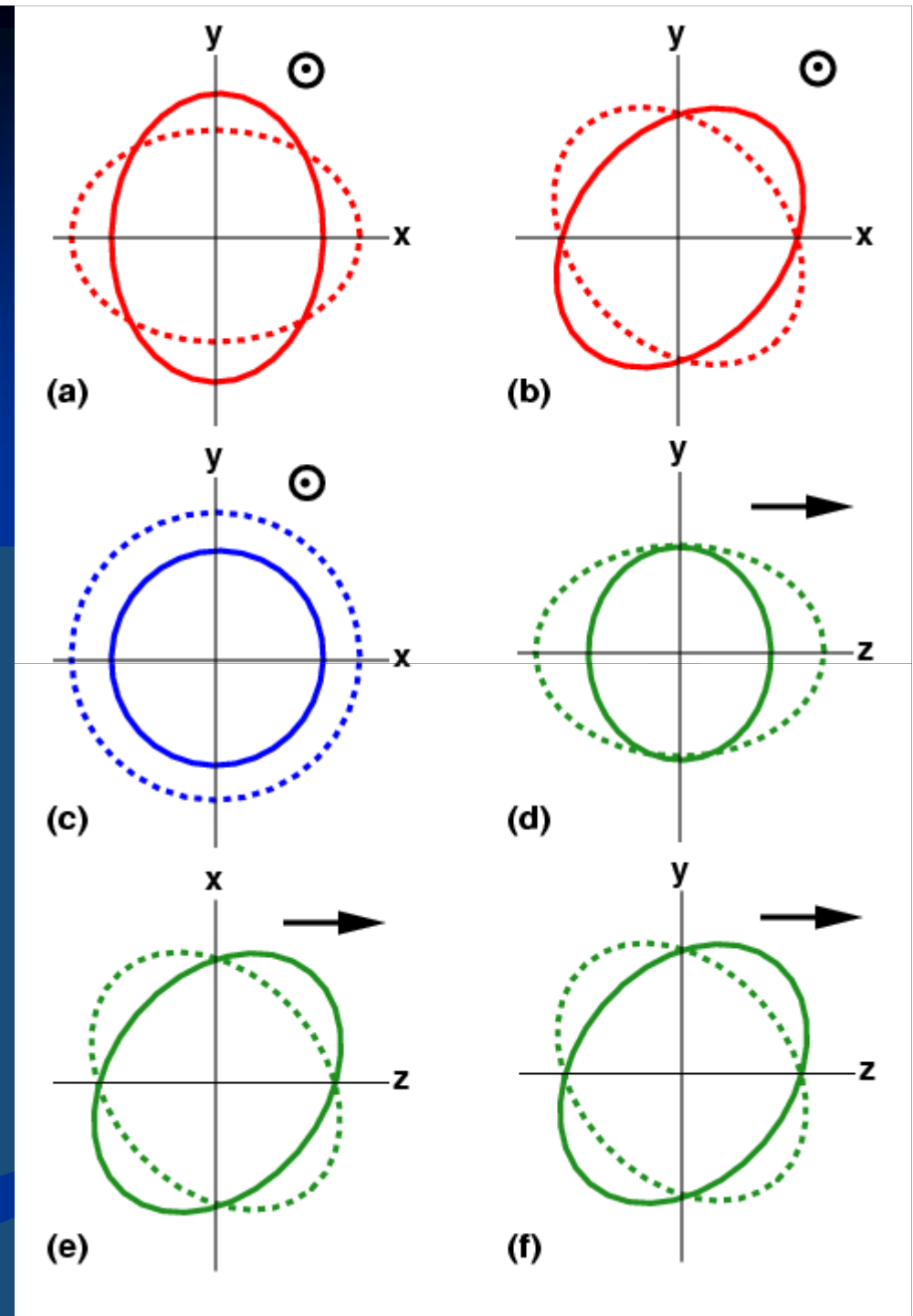
$$h_{xx} + h_{yy} + h_{zz} = 0$$

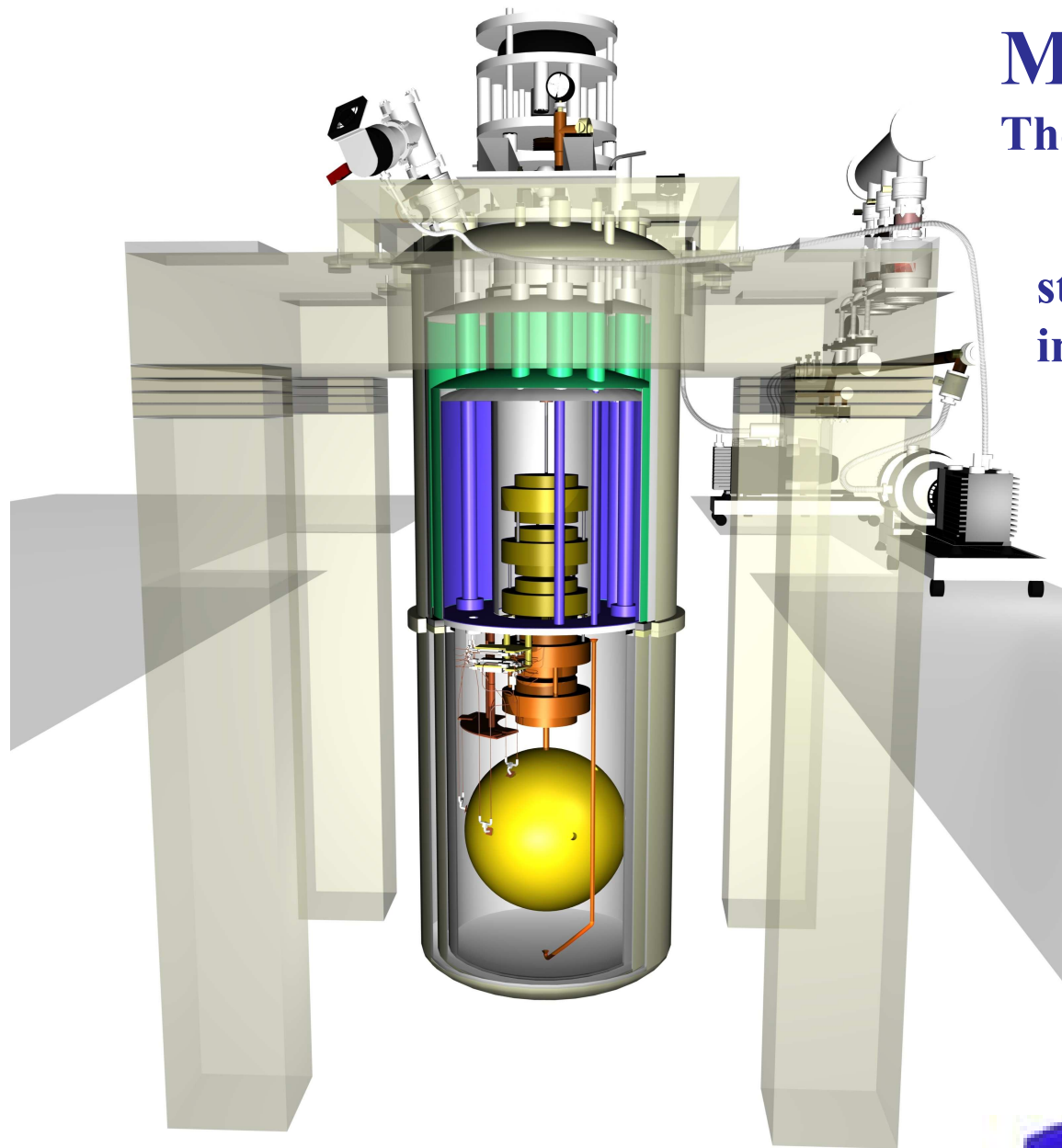
$$Ih = \begin{bmatrix} h_{xx} & h_{xy} & h_{xz} \\ h_{yx} & h_{yy} & h_{yz} \\ h_{zx} & h_{zy} & h_{zz} \end{bmatrix}$$

5 independent
components

Test of General Relativity and other Theories of Gravitation

- Polarizations
- Speed
- Quantization





Mario SCHENBERG

The Gravitational Wave Detector (Brazil)

**started commissioning operation
in the 8th of September, 2006.**

**It involves a
collaboration
between**

INPE

ITA

IFSP

UNICAMP

UNIFESP

**USP, among others,
and it has been
supported by**



Gravitational Wave Detectors

(V. Fafone)

● Interferometric

● Resonant-Mass



First gravity wave coincidence experiment between resonant cryogenic detectors: Louisiana-Rome-Stanford

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Summary. The results of a coincidence search for short bursts of gravitational radiation with cryogenic resonant-mass detectors are reported. No significant excess of coincidences at zero time delay were found. The data have been used to set an improved observational upper limit on the flux of impulsive gravitational waves that may be impinging on the Earth.

employs a resonant capacitive transducer (Rapagnani, 1982) matched to a d.c. SQUID amplifier (Carelli, 1985).

The performance of the three detectors during this coincidence experiment did not reach the design goals or previously achieved levels by the Stanford detector in either sensitivity or in non-Gaussian disturbance level (Boughn, 1982). Despite this situation,

the limit that we are able to set on the rate of gravity wave pulses impinging on the Earth is better than that set by any previous observations.

Key words: gravitational waves - detectors, gravitational waves - coincidence experiment

Class. Quantum Grav. 26 (2009) 085009 (30pp) doi:10.1088/0264-9381/26/8/085009

Gravitational wave burst search in the Virgo C7 data

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Abstract

A search for gravitational wave burst events has been performed with the Virgo C7 commissioning run data that have been acquired in September 2005 over 5 days. It focused on unmodeled short duration signals in the frequency range 150 Hz to 2 kHz. A search aimed at detecting the GW emission from the merger and ring-down phases of binary black hole coalescences was also carried out. An extensive understanding of the data was required to be able to handle a burst search using the output of only one detector. A 90% confidence level upper limit on the number of expected events given the Virgo C7 sensitivity curve has been derived as a function of the signal strength, for unmodeled gravitational wave searches. The sensitivity of the analysis presented is, in terms of the root sum square strain amplitude, $h_{\text{rss}} \simeq 10^{-20} \text{ Hz}^{-1/2}$. This can be interpreted in terms of a frequentist upper limit on the rate $\dot{R}_{90\%}$ of detectable gravitational wave bursts at the level of 1.1 events per day at a 90% confidence level. From the binary black hole search, we obtained the distance reach at 50% and 90% efficiency as a function of the total mass of the final black hole. The maximal detection distance for non-spinning high and equal mass black hole binary system obtained by this analysis in C7 data is $\sim 2.9 \pm 0.1$ Mpc for a detection efficiency of 50% for a binary of total mass $80 M_{\odot}$.

PACS numbers: 04.80.Nn, 04.30.Tv, 95.30.Sf, 95.85.Sz

(Some figures in this article are in colour only in the electronic version)

1. Introduction

Virgo [1] is a 3 km long arm power-recycled Michelson interferometer located near Pisa, Italy, whose goal is to detect gravitational waves (GW) emitted by astrophysical sources extending out past the Virgo cluster. The commissioning of the detector started in 2003 and regular data taking campaigns have been organized after each important milestone. The last commissioning run (C7) took place in September 2005 and lasted for 5 days. The best achieved sensitivity was $h \simeq 7 \times 10^{-22} \text{ Hz}^{-1/2}$ at 300 Hz. The Virgo design sensitivity at this frequency is expected to be better by an order of magnitude assuming that 10 W enters into the interferometer. However, during the C7 run, Virgo was running with a reduced light power, 0.7 W, because the backscattering in the mode-cleaner cavity of the light reflected by the recycling mirror

First joint gravitational wave search by the AURIGA–EXPLORER–NAUTILUS–Virgo Collaboration

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⁷ Laboratoire d'Année-le-Vieux de Physique des Particules (LAPP), IN2P3/CNRS, Université de Savoie, Annecy-le-Vieux, France
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⁹ LAL, Université Paris-Sud, IN2P3/CNRS, Orsay, France
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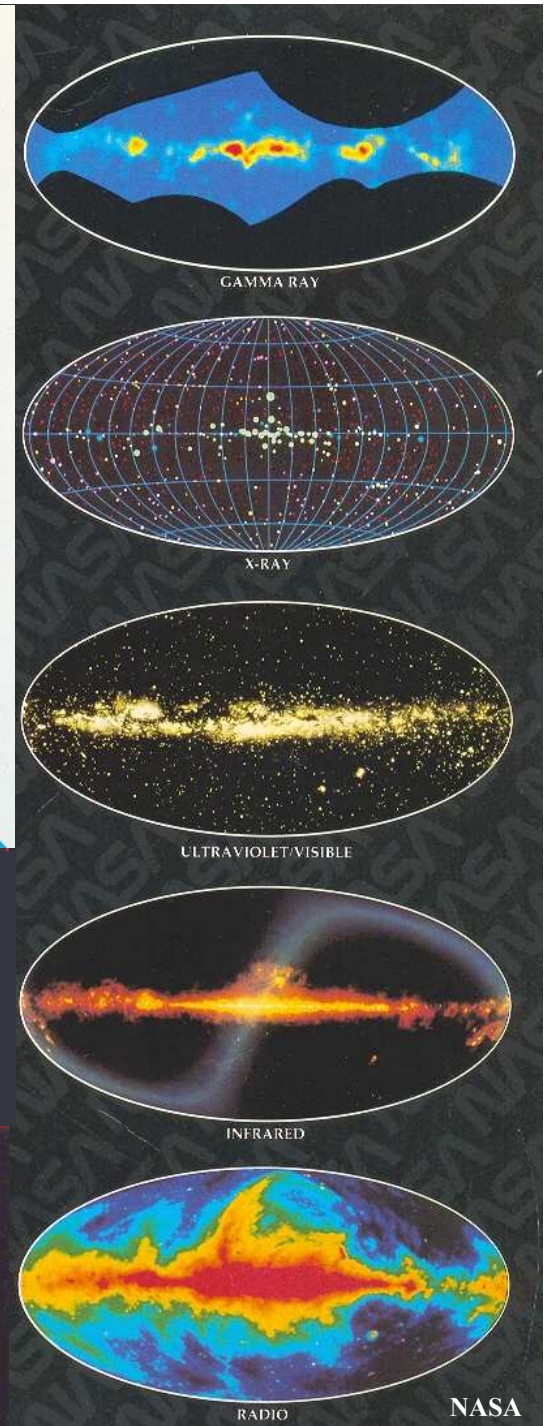
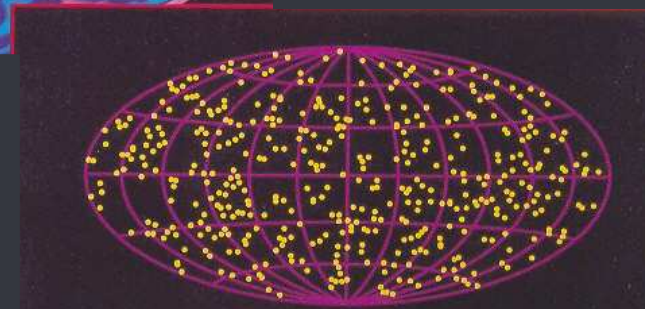
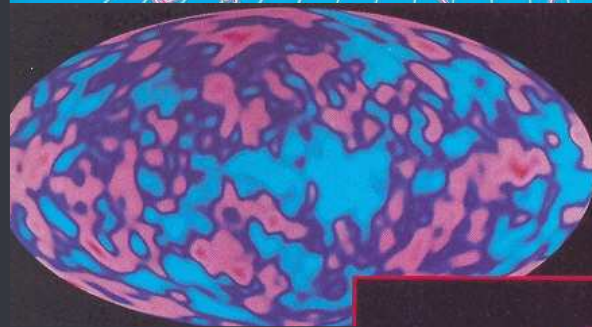
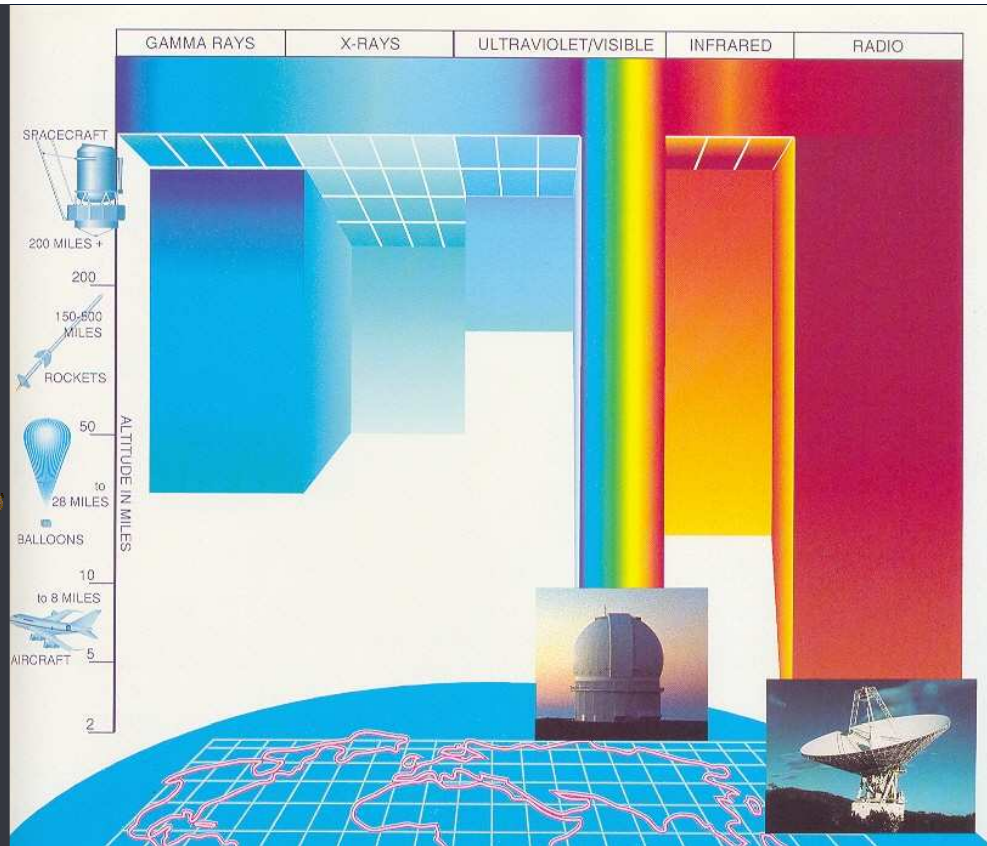
Published 30 September 2008

Online at stacks.iop.org/CQG/25/205007

Abstract

We present a methodology of network data analysis applied to the search for coincident burst excitations over a 24 h long data set collected by AURIGA, EXPLORER, NAUTILUS and Virgo detectors during September 2005. The search of candidate triggers was performed independently on each of the data sets from single detectors. We looked for two-fold time coincidences between these candidates using an algorithm optimized for a given population of sources and we calculated the efficiency of detection through injections of templated signal waveforms into the streams of data. To this end we have considered the

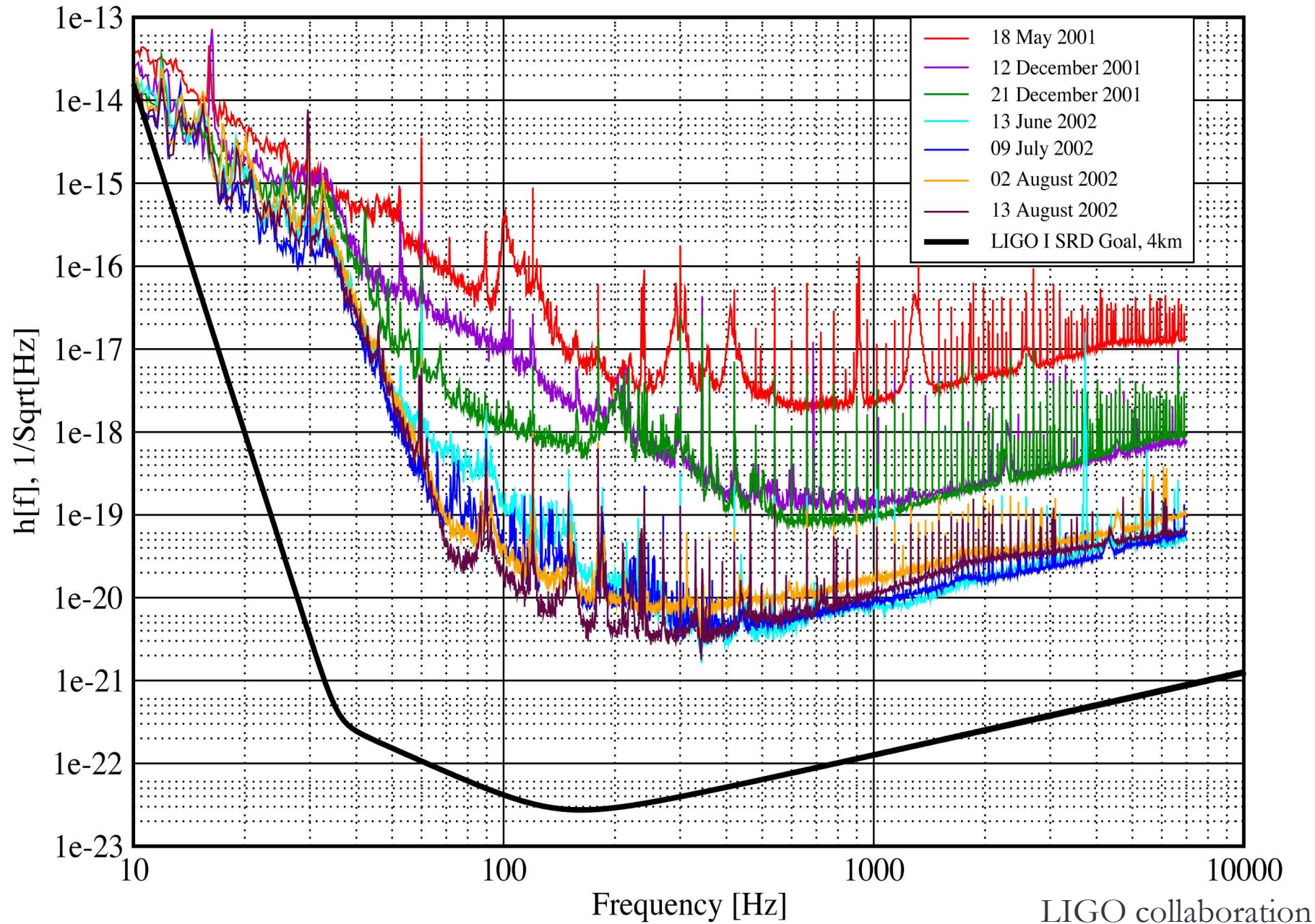
Gravitational spectrum (a new window to observe the Universe)



Strain Sensivities for the LIGO Livingston 4km Interferometer, E7 to S1

18 May 2001 - 13 August 2002

LIGO-G020451-00-E

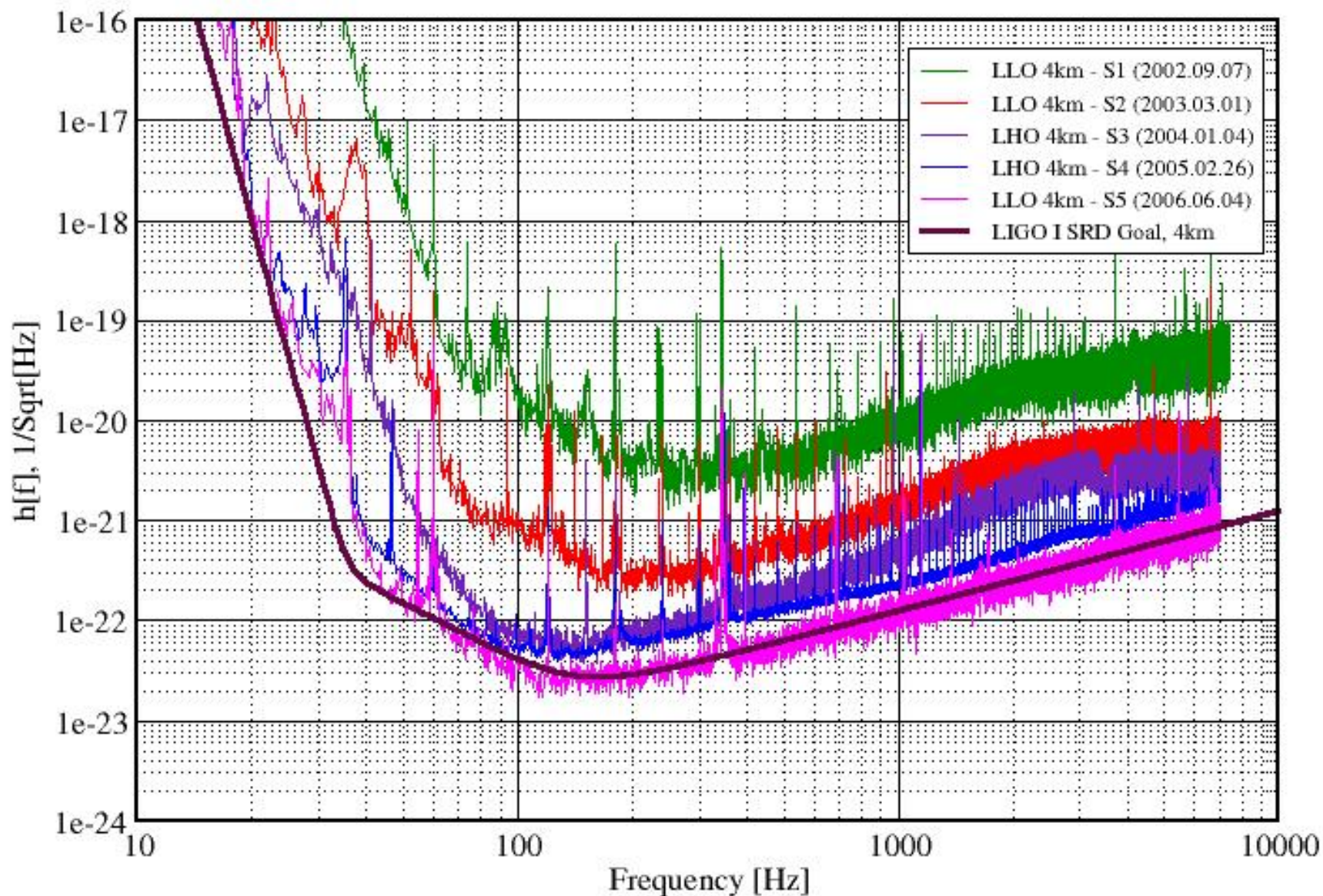


LIGO collaboration

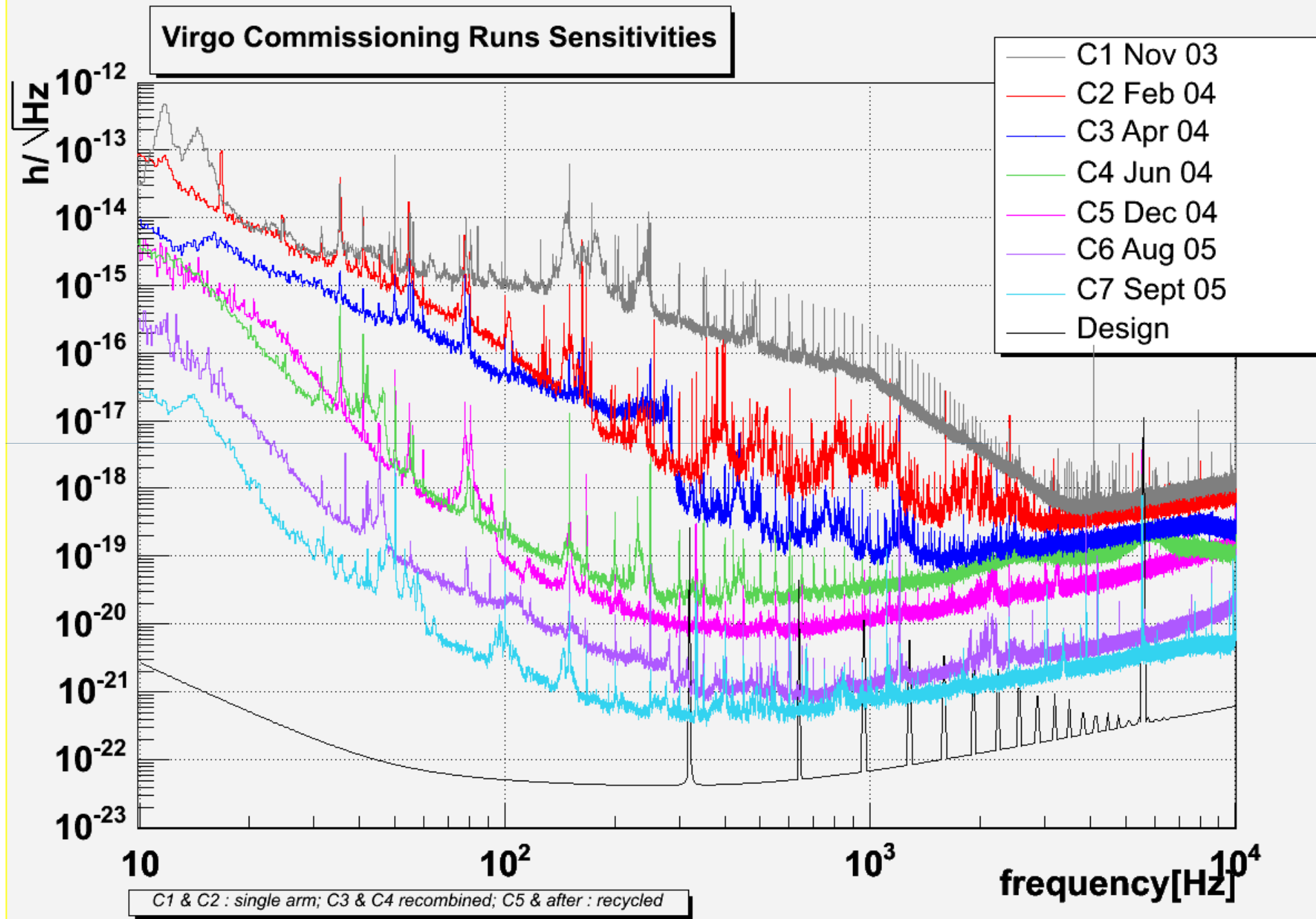
Best Strain Sensivities for the LIGO Interferometers

Comparisons among S1 - S5 Runs

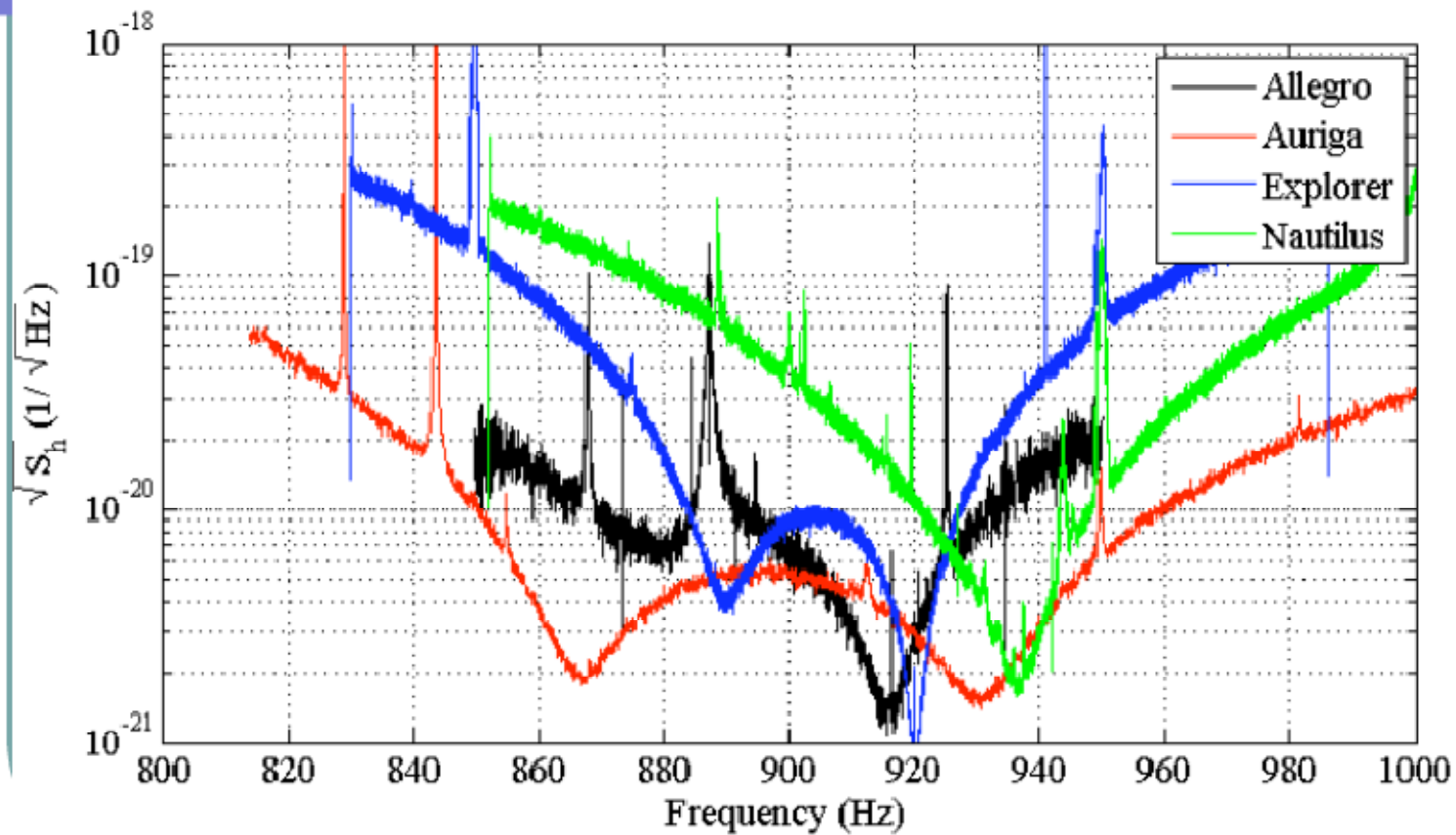
LIGO-G060009-02-Z



LIGO collaboration



Sensitivity of Resonant Detectors



credit: Kostas Kokkotas

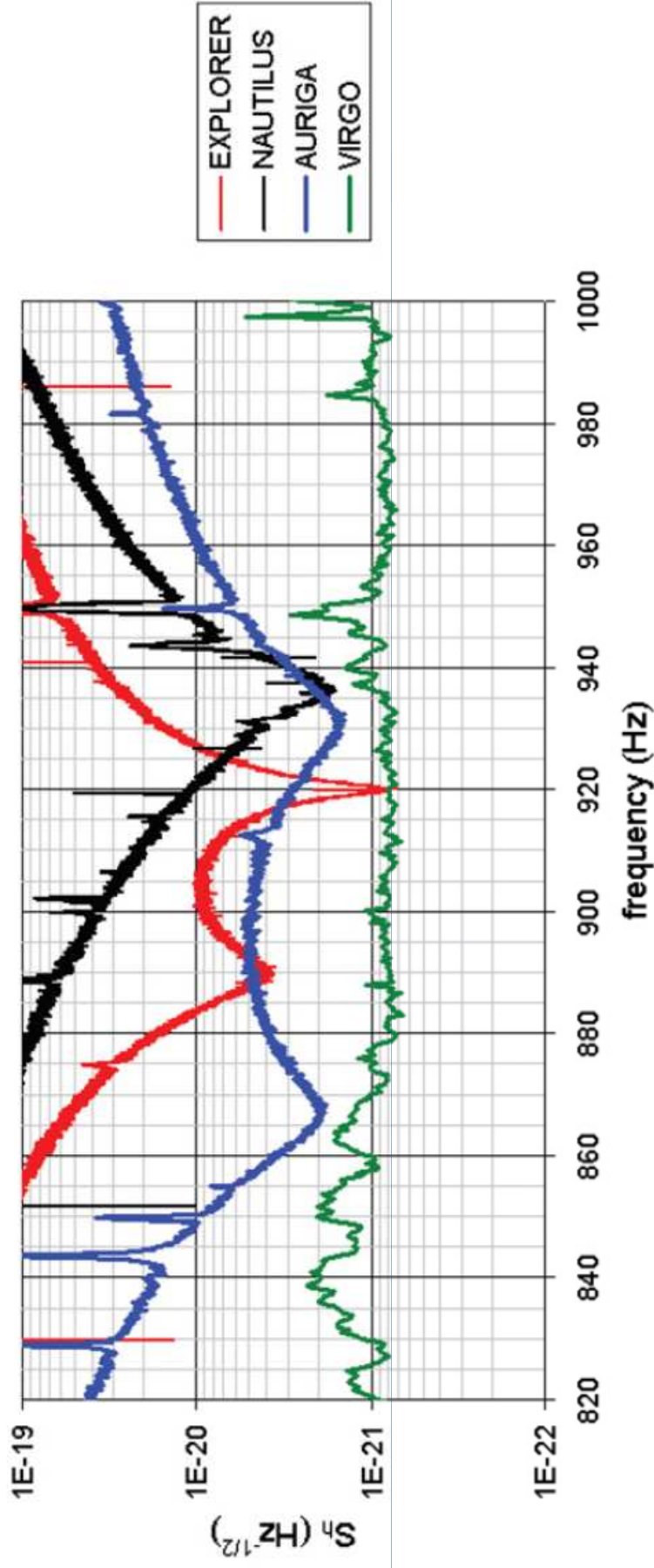
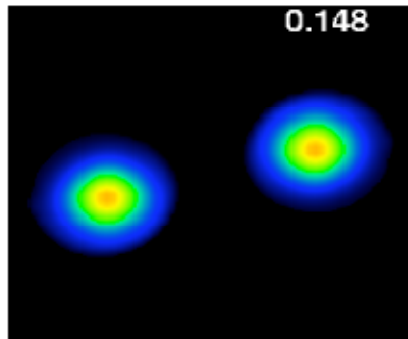
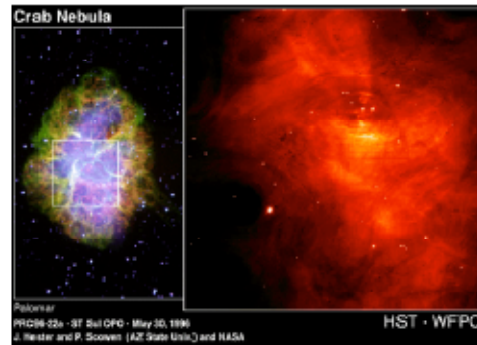


Figure 1. Typical spectral density of calibrated noise for the three resonant bar detectors during 2005 and for the Virgo interferometer in September 2005.

GW sources for Interferometric and Acoustic Detectors



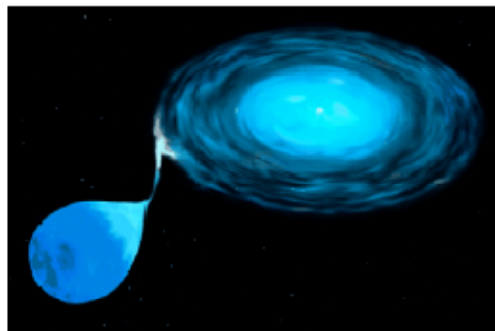
BH and NS Binaries



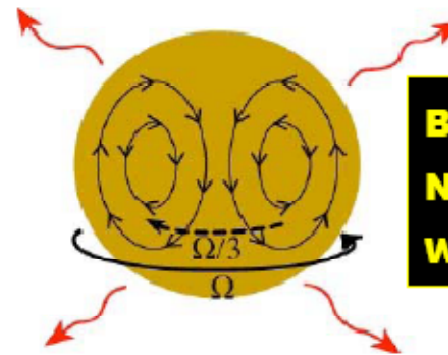
Supernovae, BH/NS formation

$$L_{GW} \sim \left(\frac{M}{R}\right)^5$$

$$h \sim \varepsilon \cdot \left(\frac{M}{r}\right) \cdot \left(\frac{M}{R}\right)$$



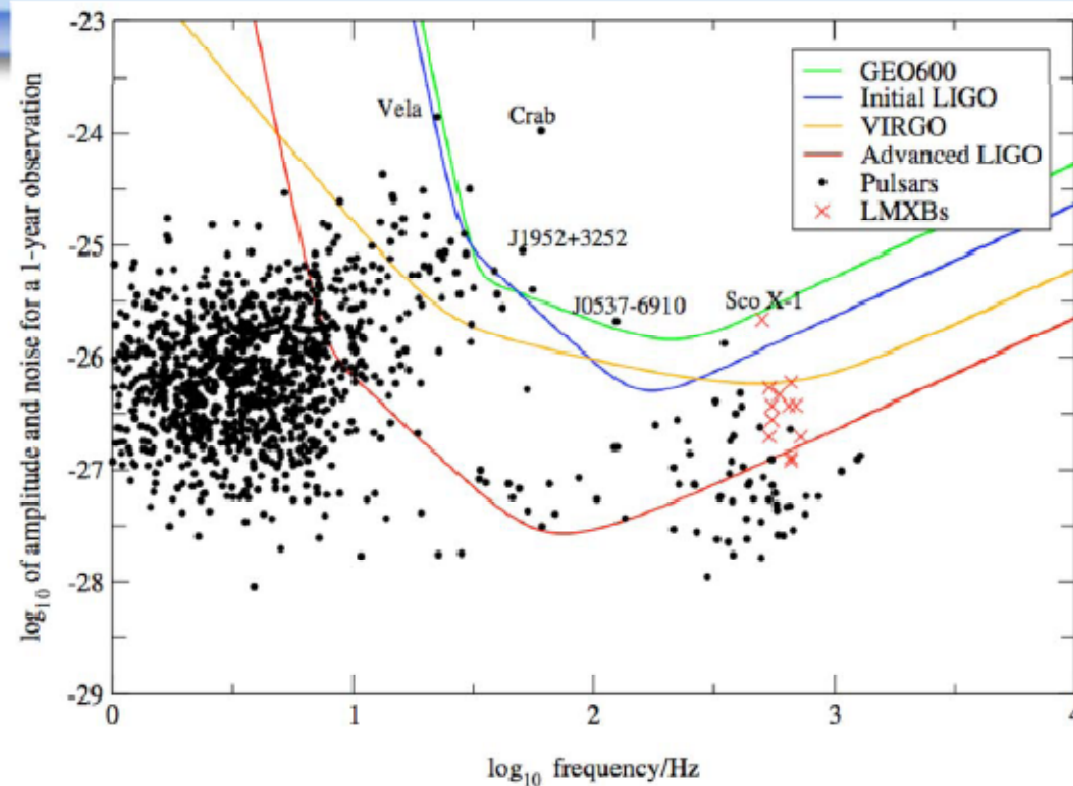
Spinning neutron stars in X-ray binaries



Young Neutron Stars

- Black Holes : $M/R=0.5$**
- Neutron Stars : $M/R\sim 0.2$**
- White Dwarfs : $M/R\sim 10^{-4}$**

Slowdown of radio pulsars

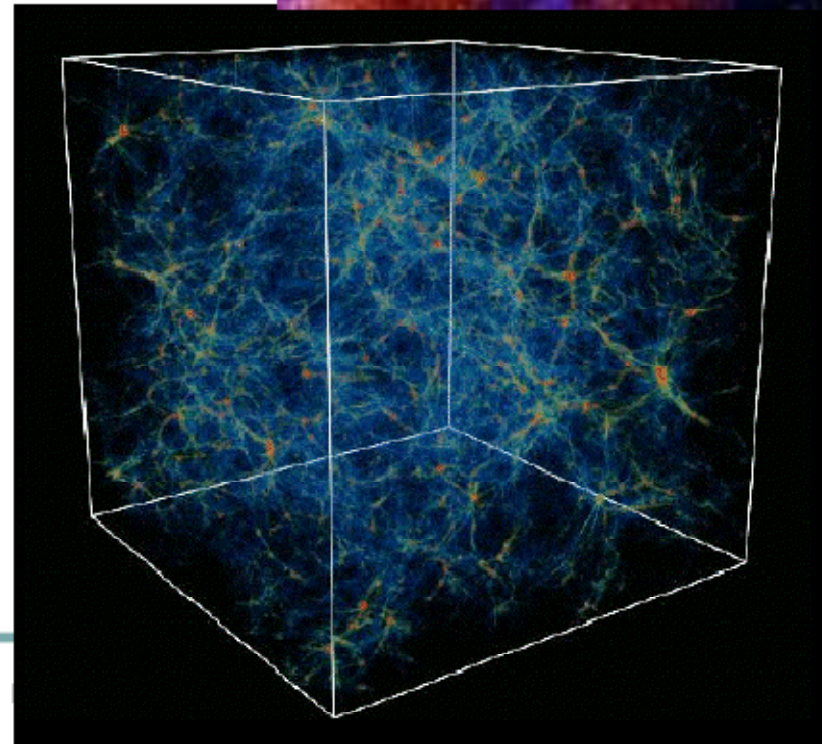


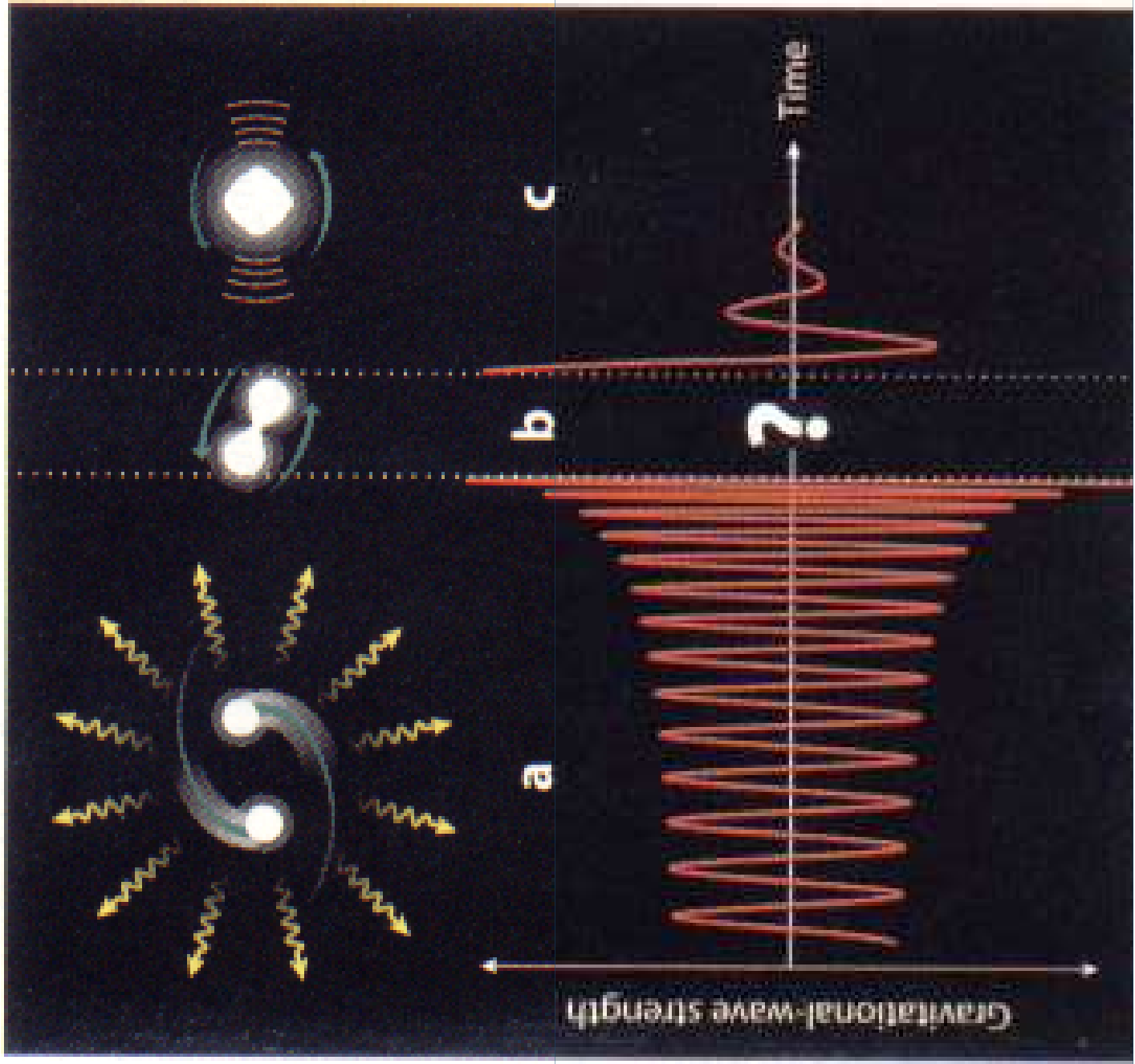
Upper limits on GW emission from triaxial NS can be used to decide which pulsars to target

+ stochastic sources

(contribute to a noisy background)

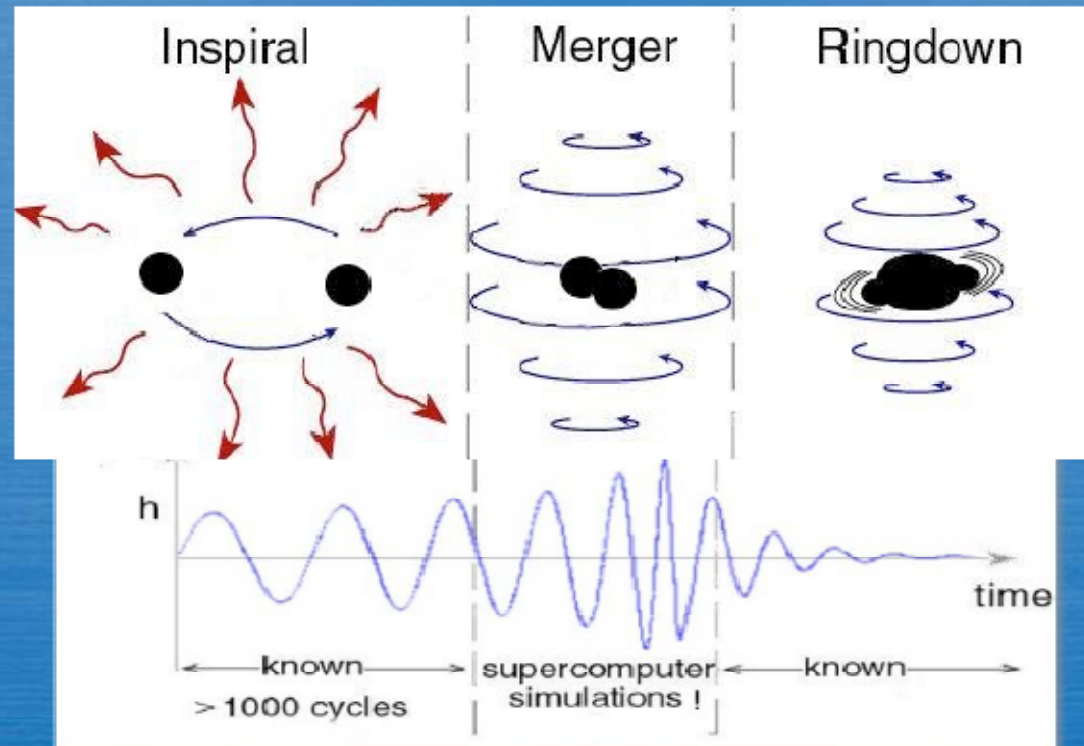
- ✓ Big Bang,
- ✓ early expansion of the Universe,
- ✓ cosmic strings,
- ✓ unresolved sources...







The emerging picture...



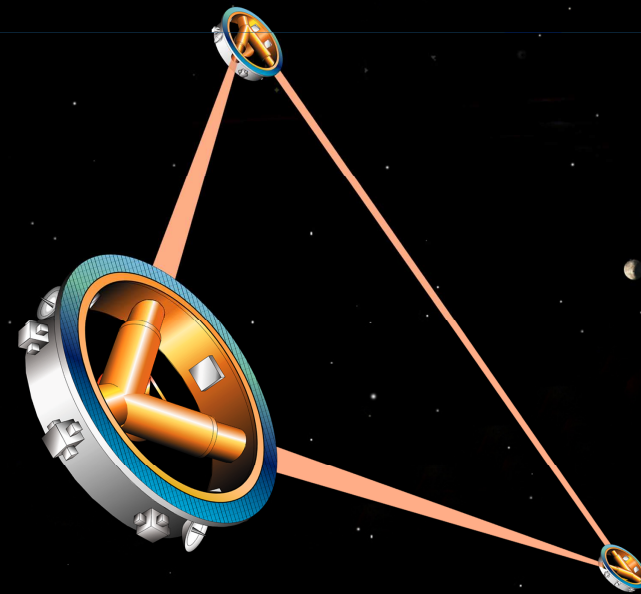
credit: Kostas Kokkotas

Space Interferometers:

LISA, ALIA and BBO

<http://lisa.jpl.nasa.gov/>

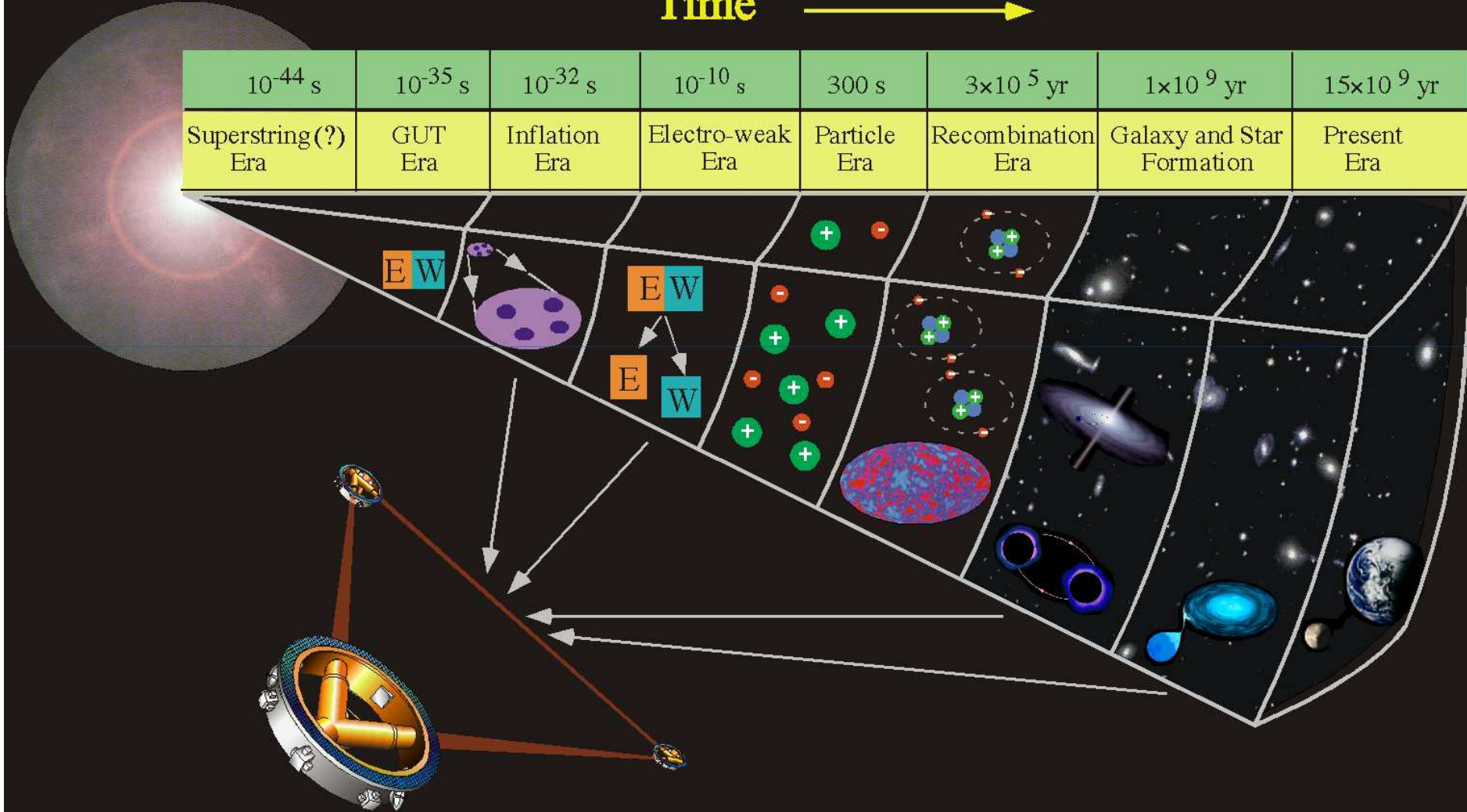
http://www.esa.int/esaSC/120376_index_0_m.html



Big Bang

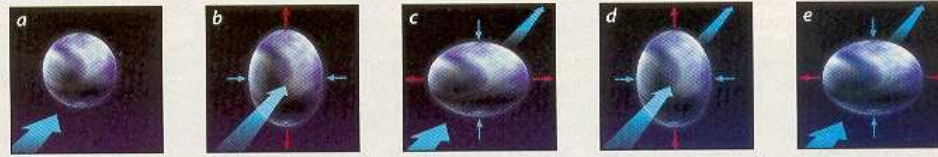
Time \longrightarrow

10^{-44} s	10^{-35} s	10^{-32} s	10^{-10} s	300 s	3×10^5 yr	1×10^9 yr	15×10^9 yr
Superstring (?) Era	GUT Era	Inflation Era	Electro-weak Era	Particle Era	Recombination Era	Galaxy and Star Formation	Present Era



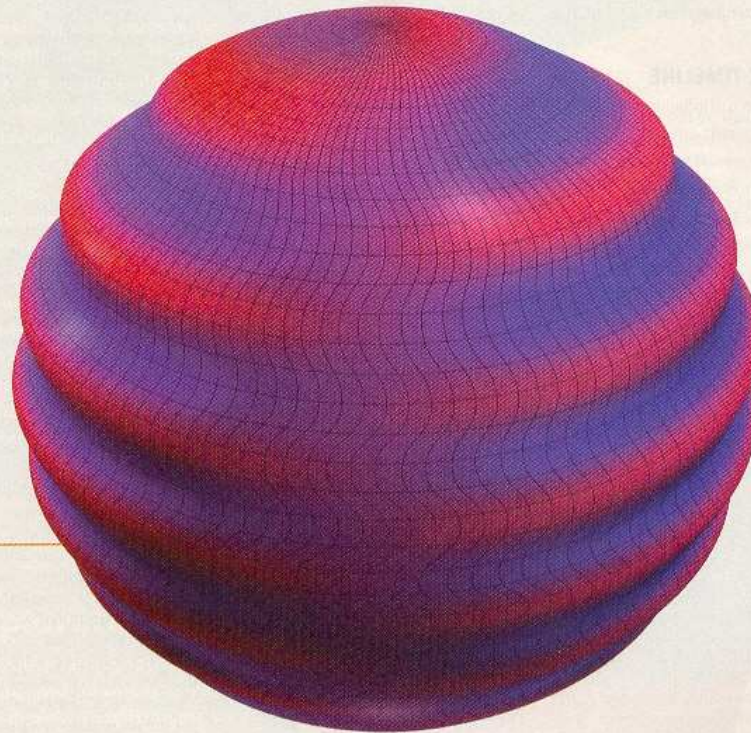
by Robert R. Caldwell
and Marc Kamionkowski

Scientific American



GRAVITATIONAL WAVES

Although gravitational waves have never been directly observed, theory predicts that they can be detected because they stretch and squeeze the space they travel through. On striking a spherical mass (a), a wave first stretches the mass in one direction and squeezes it in a perpendicular direction (b). Then the effects are reversed (c), and the distortions oscillate at the wave's frequency (d and e). The distortions shown here have been greatly exaggerated; gravitational waves are usually too weak to produce measurable effects.



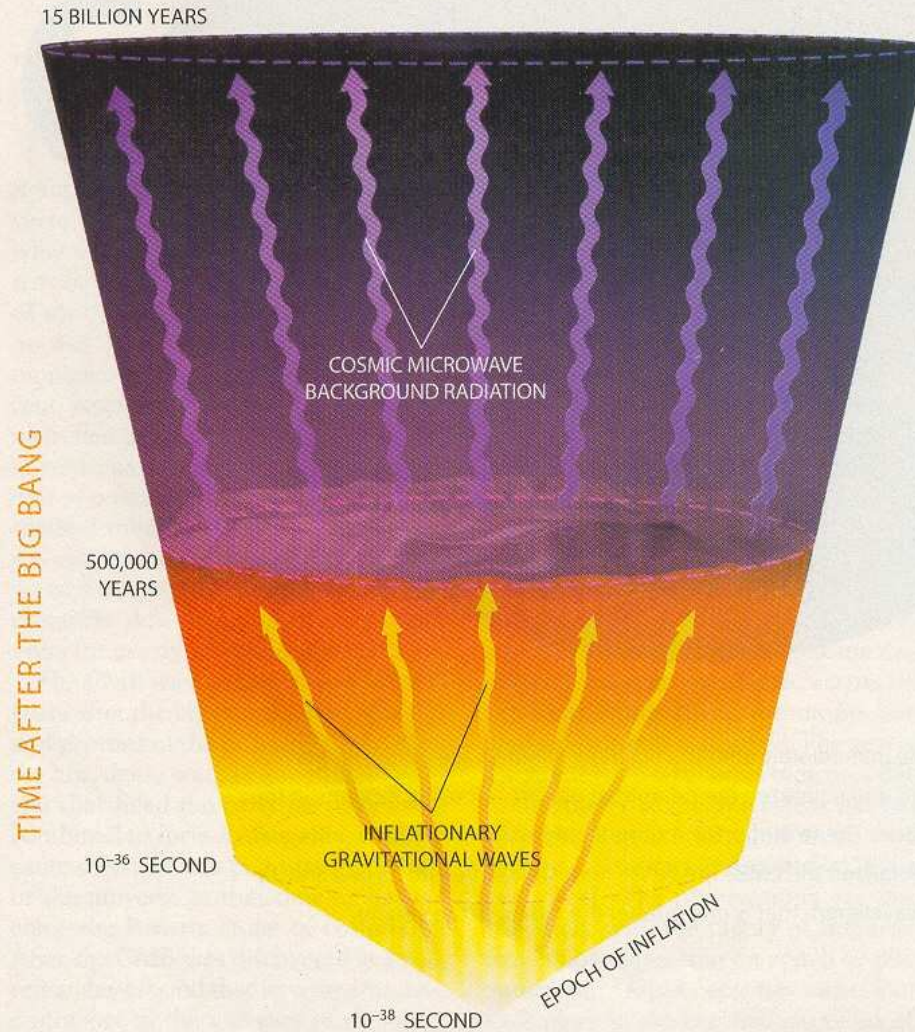
DISTORTED UNIVERSE

The fantastically rapid expansion of the universe immediately after the big bang should have produced gravitational waves. These waves would have stretched and squeezed the primordial plasma, inducing motions in the spherical surface that emitted the CMB radiation. These motions, in turn, would have caused redshifts and blueshifts in the radiation's temperature and polarized the CMB. The figure here shows the effects of a gravitational wave traveling from pole to pole, with a wavelength that is one quarter the radius of the sphere.

ILLUSTRATIONS BY SLM FILMS

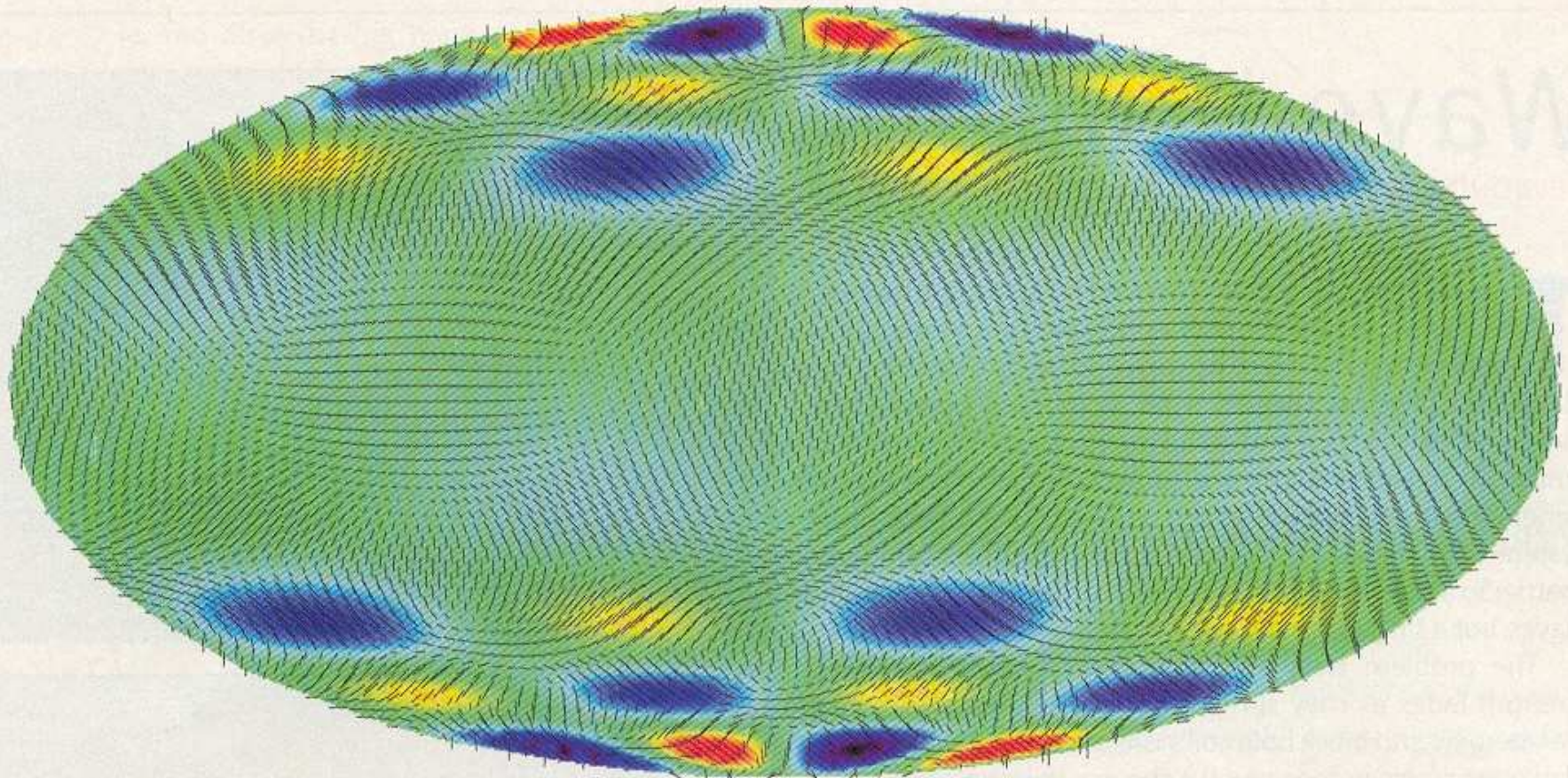
■ COSMIC TIMELINE

During the epoch of inflation—the tremendous expansion of the universe that took place in the first moments after the big bang—quantum processes generated a spectrum of gravitational waves. The waves echoed through the primordial plasma, distorting the CMB radiation that was emitted about 500,000 years later. By carefully observing the CMB today, cosmologists may detect the plasma motions induced by the inflationary waves.



by Robert R. Caldwell
and Marc Kamionkowski

Scientific American



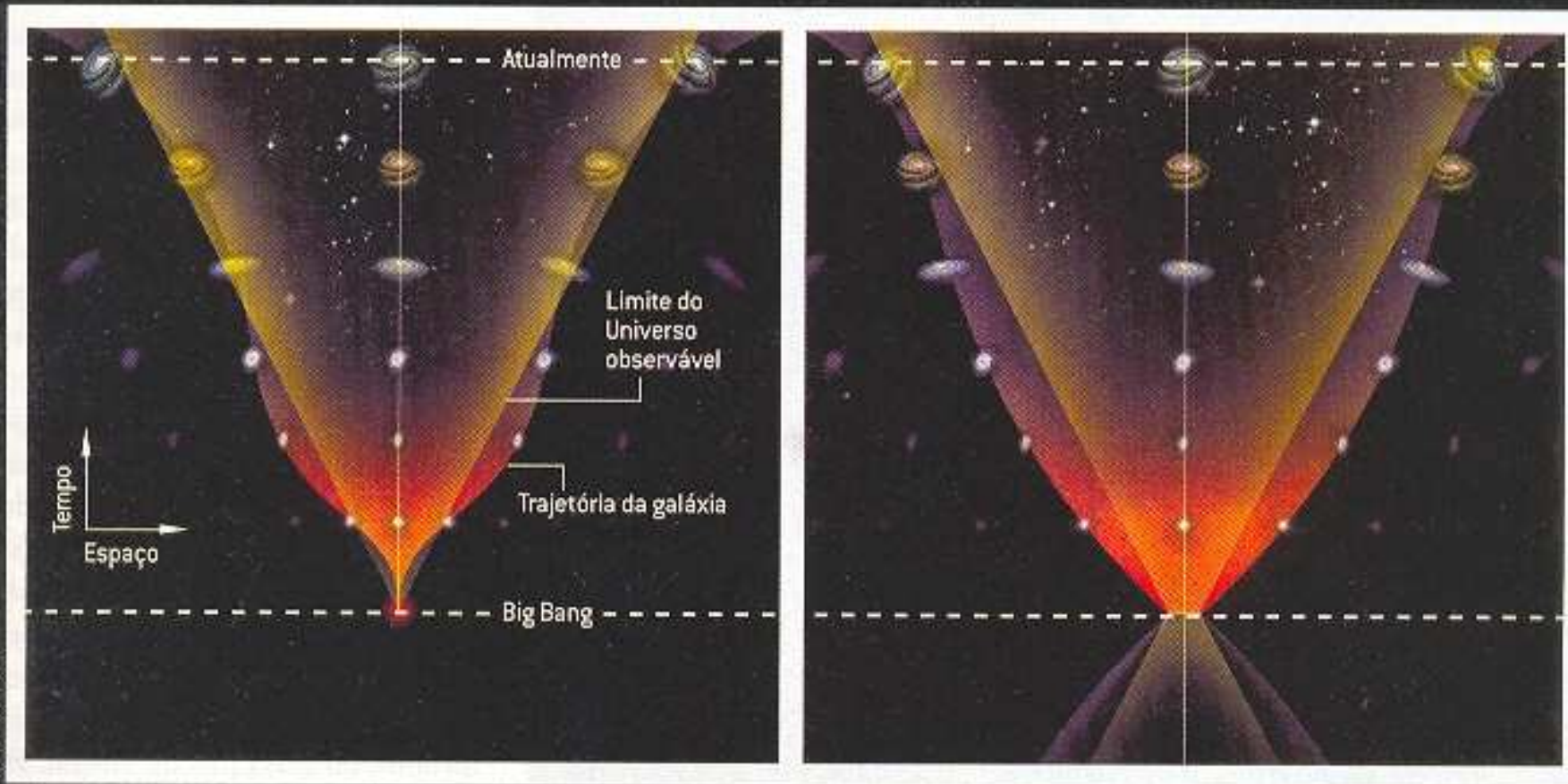
■ RELIC IN THE RADIATION

Inflationary gravitational waves would have left a distinctive imprint on the CMB. The diagram here depicts the simulated temperature variations and polarization patterns that would result from the distortions shown in the bottom illustration on page 39. The red and blue spots represent colder and hotter regions of the CMB, and the small line segments indicate the orientation angle of the polarization in each region of the sky.

DUAS VISÕES DO PRINCÍPIO

Gabriele Veneziano, 2004
Scientific American Brasil

Em nosso Universo em expansão, as galáxias se afastam rapidamente umas das outras, a uma velocidade proporcional à distância entre elas: duas galáxias separadas por uma distância de 500 milhões de anos-luz se afastam a uma velocidade duas vezes maior do que duas galáxias separadas por uma distância de 250 milhões de anos-luz. Por isso, todas as galáxias que vemos devem ter se originado em um mesmo ponto e ao mesmo tempo – o Big Bang. A conclusão é válida mesmo quando a expansão cósmica passou por períodos de aceleração e desaceleração; em diagramas de espaço-tempo (abaixo), as galáxias seguem trajetórias sinuosas que as levam para dentro e para fora da região observável do espaço (triângulo amarelo). No entanto, a situação torna-se incerta no instante exato em que as trajetórias das galáxias (ou suas predecessoras) começaram a divergir.



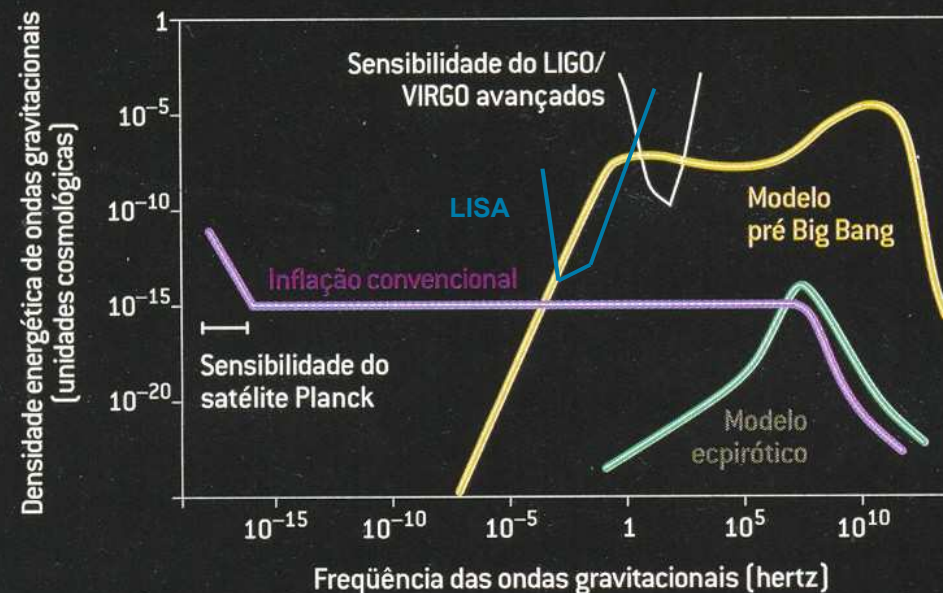
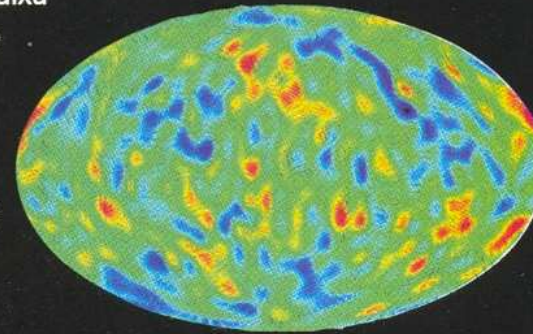
SAMUEL VELASCO

Na cosmologia do Big Bang convencional, que se baseia na teoria da relatividade geral de Einstein, a distância entre quaisquer duas galáxias era zero há um tempo finito. O tempo perde significado antes desse instante.

Nos modelos mais sofisticados, que incluem efeitos quânticos, quaisquer duas galáxias devem ter surgido a uma certa distância mínima uma da outra. Esses modelos abrem a possibilidade de um Universo pré-Big Bang.

EFEITO DE FUNDO

Observar o Universo pré-Big Bang pode parecer uma tarefa sem futuro, mas uma forma de radiação poderia sobreviver desde aquela época: a radiação gravitacional. Essas variações periódicas na campo gravitacional poderiam ser detectadas indiretamente, por seu efeito na polarização do fundo cósmico de microondas (*vista simulada, abaixo*), ou diretamente, em observatórios terrestres. Os cenários pré-Big Bang e ecpirótico prevêem mais ondas gravitacionais de alta freqüência e menos de baixa freqüência do que os modelos convencionais de inflação (*abaixo*). Medidas existentes de vários fenômenos astronômicos não conseguem fazer a distinção entre esses modelos, mas futuras observações do satélite Planck, assim como dos observatórios LIGO e VIRGO, devem ser capazes disso.



Exotic events, anticipated from theory, which may be confirmed when gravitational wave astronomy becomes a reality:

- Binaries formed by primordial black holes (MACHOS?),
- Cosmic strings,
- Cosmic bubbles,
- Boson stars,
- Quark stars (strange stars) (test of QCD),
- Phenomena not yet known or anticipated, and surprises.

If nature has **more dimensions** than the four of **spacetime**, gravity will have “transit” in these extra dimensions, and **special effects on the intensity, propagation and polarization** of gravitational waves may occur different from the ones anticipated by the General Relativity. The study of these effects on the gravitational waves would be, therefore, a tool for the investigation of extra dimensions in the Universe and to formulate an auto consistent theory of quantum gravity.

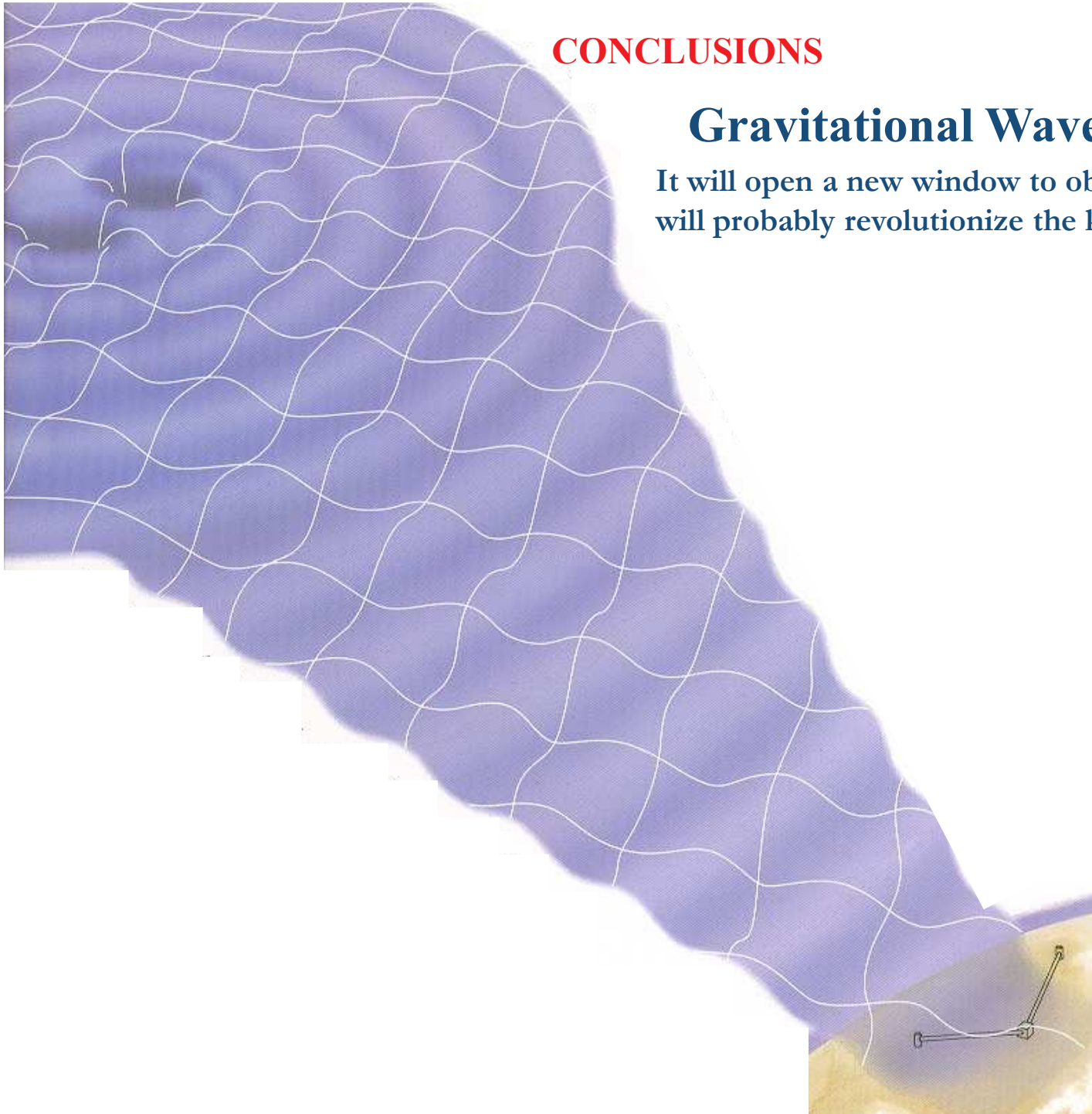
Universal GW Radio Broadcast (100 MHz FM)

Carrier produced by a rotating 10^{-4} solar mass black hole and modulated by a controlled accretion flow.

CONCLUSIONS

Gravitational Wave Astronomy

It will open a new window to observe the Universe, which will probably revolutionize the knowledge we have of it.



CONCLUSIONS

The sources of gravitational waves can be:

- galactic,
- extragalactic,
- cosmological,
- from previous universes.

The spectrum goes from 10^{-18} Hz to 10^{10} Hz.

If extra-dimensions exist, gravitational wave have transit in them.

The probability of new revolutionary discoveries is very high.

Because spherical antennas are able to determine the origin of the signal in the sky and its polarizations, they will probably play an important rule on these new discoveries.

Thank you !

