

XII IAG-USP Advanced School on Astrophysics

THE IONIZED INTERSTELLAR MEDIUM

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SUMMARY

- brief introduction
 - definition, phases
 - same physical processes → same tools for analyzing the spectra
- methods for the determination of physical parameters
 - empirical methods: T_e , n_e , Z
 - photoionization models: physical conditions, Z , ionizing source(s), n_H

definition: "*ionized medium*"

• H is the most abundant element → *ionized medium* \equiv presence of H^+

• H ionization potential = 13.6 eV

∴ if H^+ ions are present,

- incident radiation with $E > 13.6$ eV
and/or
- energetic ionizing particles
and/or
- winds, shock waves, ...

in most ionized regions
the gas is mainly
photoionized

(cosmic rays can ionize H at the borders of neutral clouds)

• **Ionized Interstellar Medium can be found in different phases**

≠ gas densities: $\sim 0.01 \rightarrow 10^9 \dots \text{cm}^{-3}$

≠ ionizing/heating sources: star(s), power-law spectrum, shocks, ...

≠ gas temperatures: $\sim 500 \rightarrow 10^{6-7} \dots \text{K}$

≠ ionization degrees: $H^+/H = >0 \rightarrow 1$

≠ sizes: $< 1 \text{pc} \rightarrow \text{galactic}$



Diffuse Ionized Gas

Extended Low-Density HII Regions

intercloud gas

Hot Ionized Medium

superbubbles

HII regions

QSO absorption-line systems

Planetary Nebulae

LINERS

Warm Ionized Medium

BLRs

NLRs

Active Galactic Nuclei

Diffuse Ionized Medium

SuperNova Remnants

Reynolds layer

InterStellar Matter

starburst galaxies

Diffuse Ionized Gas

Extended Low-Density HII Regions

intercloud gas

Hot Ionized

⇒ ionized plasmas with same physics

different physical conditions (and emitted spectra)
mainly due to
different heating/ionizing sources

⇒ can be analyzed with similar techniques

Plane

SuperNova Remnants

Diffuse Ionized Medium

Reynolds layer

InterStellar Matter

starburst galaxies

→ the ionized interstellar gas produces

- continuum emission
- emission-line spectrum
- can present absorption lines when observed against continuum sources

→ the details of the spectrum are determined by the physical conditions of the emitting gas



→ observed line (and continuum) intensities can be used as probes of the excitation conditions of the gas, providing robust information on:

- properties of the ionizing/heating source(s)
- gas density (distribution)
- chemical composition
- gas temperature

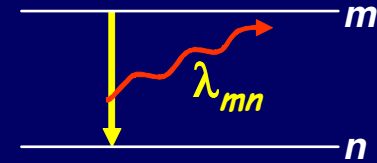
→ many of these regions are also useful laboratories for atomic physics

→ LTE relations (Boltzmann, Saha) do not apply at these densities

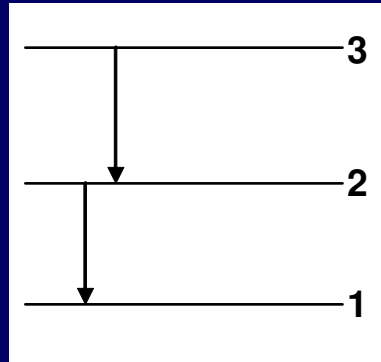
empirical methods - I. The gas temperature

⇒ forbidden line intensities

$$L_{\lambda_{mn}} = n_m(X^i) A_{mn} E_{mn} V = n(X^i) n_e \epsilon_{\lambda_{mn}}(T_e, n_e) V$$



- lines from 2 ≠ upper levels with **very** ≠ E_{exc} → line intensity ratio = $f(T_e)$

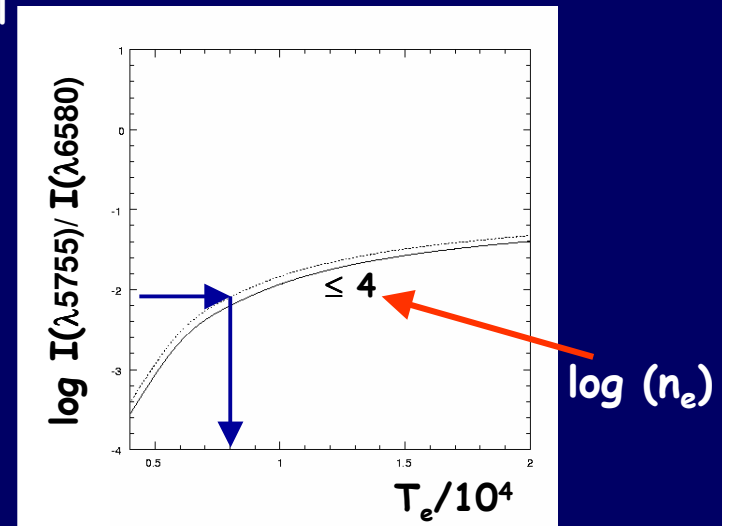
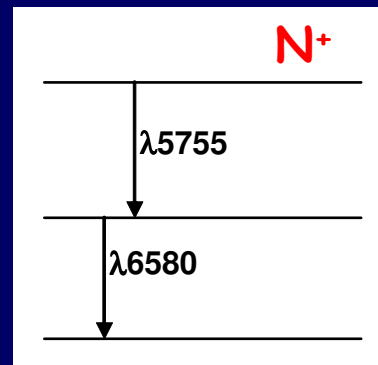


$$\frac{I_{\lambda_{32}}}{I_{\lambda_{21}}} = \frac{\epsilon_{\lambda_{32}}(T_e, n_e)}{\epsilon_{\lambda_{21}}(T_e, n_e)} = f(T_e, n_e) \sim f(T_e)$$

T_e

observed obtained from the statistical equilibrium of the energy levels

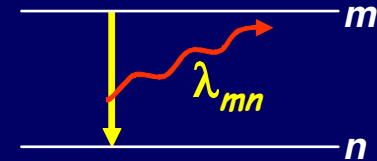
exs.: [OIII]λ4363/5000
[NII] λ5755/6580
[SIII]λ6312/9532



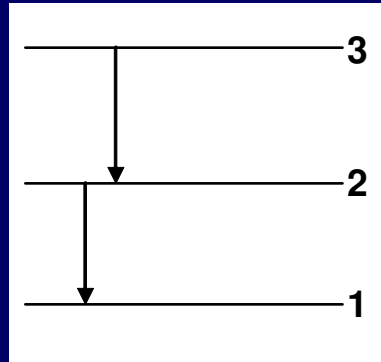
empirical methods - I. The gas temperature

⇒ forbidden line intensities

$$L_{\lambda_{mn}} = n_m(X^i) A_{mn} E_{mn} V = n(X^i) n_e \epsilon_{\lambda_{mn}}(T_e, n_e) V$$



- lines from 2 \neq upper levels with **very** $\neq E_{exc}$ → line intensity ratio = $f(T_e)$

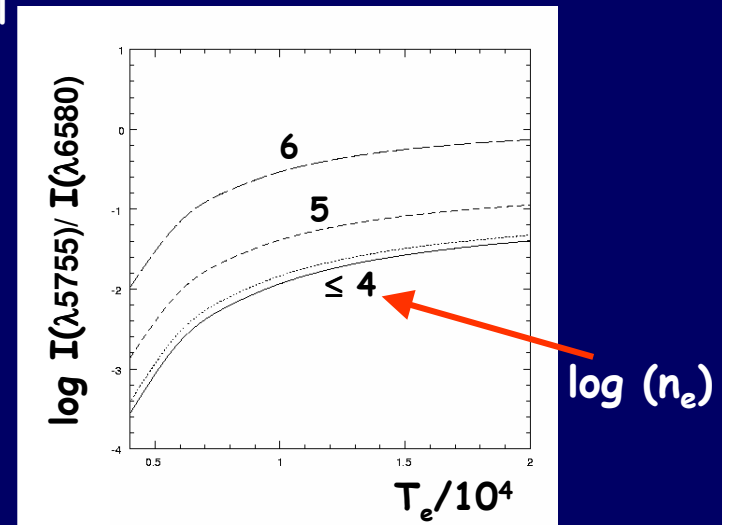
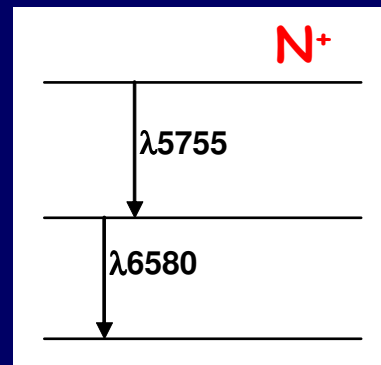


$$\frac{I_{\lambda_{32}}}{I_{\lambda_{21}}} = \frac{\epsilon_{\lambda_{32}}(T_e, n_e)}{\epsilon_{\lambda_{21}}(T_e, n_e)} = f(T_e, n_e) \sim f(T_e)$$

T_e

observed obtained from the statistical equilibrium of the energy levels

exs.: [OIII] λ 4363/5000
 [NII] λ 5755/6580
 [SIII] λ 6312/9532
 .
 .



empirical methods - I. The gas temperature

⇒ optical continuum measurements

- recombination continuum/H recombination line = $f(T_e)$

two choices for the continuum:

$$\frac{I(\lambda 4861)}{I(H\beta)}$$

$$\frac{I_\nu(\lambda 3646^-) - I_\nu(\lambda 3646^+)}{I(H\beta)}$$

Balmer discontinuity

⇒ radio continuum measurements

- at sufficiently low ν , nebulae are optically thick $\Rightarrow I_\nu \sim B_\nu(T_b) \Rightarrow T_b \rightarrow T$

transfer equation
with $\tau_\nu \gg 1$

⇒ radio recombination lines

(as shown in L. Rodriguez and
M. Elitzur lectures)

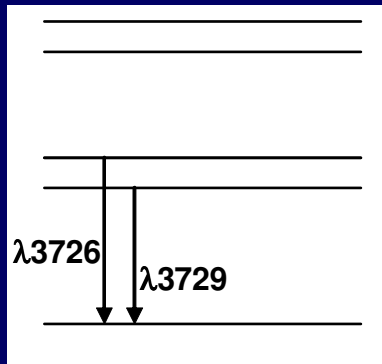
(W. Villas-Boas' lecture)

empirical methods - II. The gas density

⇒ forbidden line intensities

$$L_{\lambda_{mn}} = n(X^i) n_e \epsilon_{\lambda_{mn}}(T_e, n_e) V$$

- lines from 2 \neq upper levels with $\sim E_{exc}$ & $\neq A_{rad}$ or C_{dexc} → line intensity ratio = $f(n_e)$



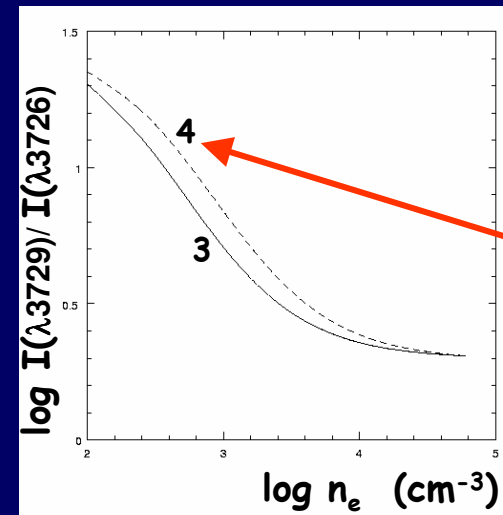
ex.: O^+

$$\frac{I_{\lambda_{31}}}{I_{\lambda_{21}}} = \frac{\epsilon_{\lambda_{31}}(T_e, n_e)}{\epsilon_{\lambda_{21}}(T_e, n_e)} = f(T_e, n_e) \sim f(n_e)$$

observed

exs.: [OII] λ 3729/3926
[SII] λ 6716/6731
many IR lines

n_e

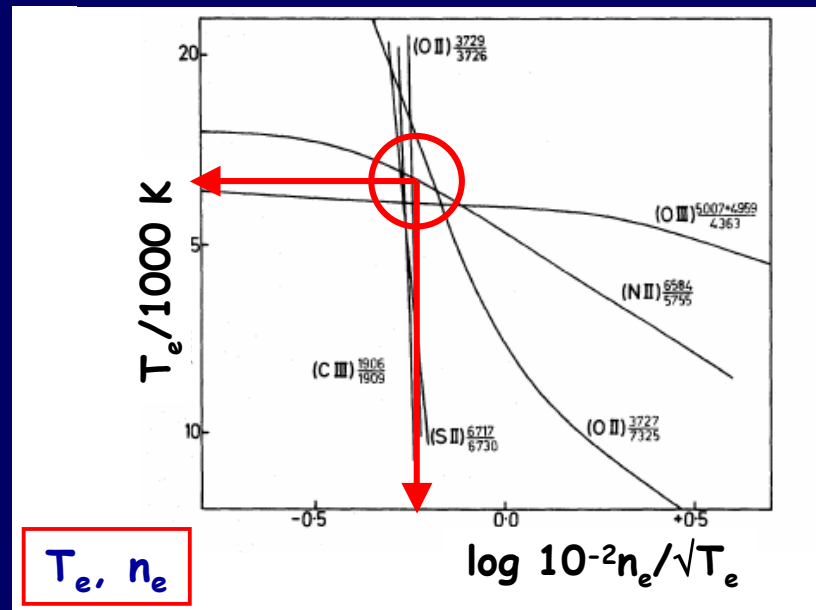


⇒ radio recombination lines

(W. Villas-Boas' lecture)

empirical methods - The gas temperature and density

⇒ combining all the information from temperature- and density-sensitive lines ratios for a given nebula



T_e, n_e

NGC6302 (Aller et al. 81)

lines of a given ion

$$\frac{I_{\lambda_{mn}}}{I_{\lambda_{kl}}} = \frac{\epsilon_{\lambda_{mn}}(T_e, n_e)}{\epsilon_{\lambda_{kl}}(T_e, n_e)} = f(T_e, n_e)$$

observed

empirical methods - III. The gas chemical composition

→ abundance of ions with observed lines

• collisionally excited line: $L_{\lambda_{mn}} = n_m(X^i) A_{mn} E_{mn} V = n(X^i) n_e \epsilon_{\lambda_{mn}}(T_e, n_e) V$

• recombination line:
(for example, H β) $L_{H\beta} = n(H^+) n_e \alpha_{H\beta}^{eff} V = n(H^+) n_e \epsilon_{H\beta}(T_e) V$

$$\frac{n(X^i)}{n(H^+)} = \frac{I_{\lambda_{mn}}}{I_{H\beta}} \frac{\epsilon_{H\beta}(T_e)}{\epsilon_{\lambda_{mn}}(T_e, n_e)} = \underbrace{\frac{I_{\lambda_{mn}}}{I_{H\beta}}}_{\text{observed}} \times f(T_e, n_e) \rightarrow \frac{n(X^i)}{n(H^+)} \left(\equiv \frac{X^i}{H^+} \right)$$

from the methods
previously discussed

empirical methods - III. The gas chemical composition

$$\frac{X}{H} \sim \frac{X^o}{H^+} + \frac{X^+}{H^+} + \frac{X^{++}}{H^+} + \dots$$

- ions with observed lines

→

$$\frac{X^i}{H^+} = \frac{I_{\lambda_{mn}}}{I_{H\beta}} \times f(T_e, n_e)$$

- ions with no observed lines

→ ionization correction factors (icf)

(based on the similarity of the ionization potentials)

⇒ elemental abundances:

$$\frac{He}{H} = \frac{He^o + He^+ + He^{++}}{H^o + H^+} \sim \frac{He^+ + He^{++}}{H^+}$$

$$\frac{O}{H} \sim \frac{O^o + O^+ + O^{++} + \dots}{H^+} = icf(O) \times \frac{O^+ + O^{++}}{H^+}$$

$$icf(O) \sim \frac{He^+ + He^{++}}{He^+}$$

$$\frac{N}{H} \sim \frac{N^o + N^+ + N^{++} + \dots}{H^+} = icf(N) \times \frac{N^+}{H^+}$$

$$icf(N) = icf(O) \left(1 + \frac{O^{++}}{O^+} \right)$$

etc.

statistical methods for the gas chemical composition (strong line method)

↳ to derive metallicities in giant extragalactic HII regions
when the electron temperature cannot be determined
([OIII] λ 4363 too weak)

→ based on calibrations from { HII regions determinations
and
theoretical models

for example,
$$R_{23} = \frac{[OII]\lambda 3727 + [OIII]\lambda 5007}{H\beta} \leftrightarrow \frac{O}{H}$$

Pagel et al. (1979) + improvements

(also O_{32} , S_{23} , N_2 , ...)

However ...

- \neq methods for estimating the gas temperature provide \neq values for T_e
→ nebulae are not isothermal
- temperature and ionization distribution depend strongly on
 - the spectral distribution of the incident radiation
 - chemical abundance (cooling)
- detailed calculation shows:
 - simple relations for the icfs do not necessarily hold
 - charge transfer reactions can be very important for some ions
 - ions with different ionization potentials can coexist
 - degeneration of the relation of R_{23} and other indexes with O/H
(double value methods)

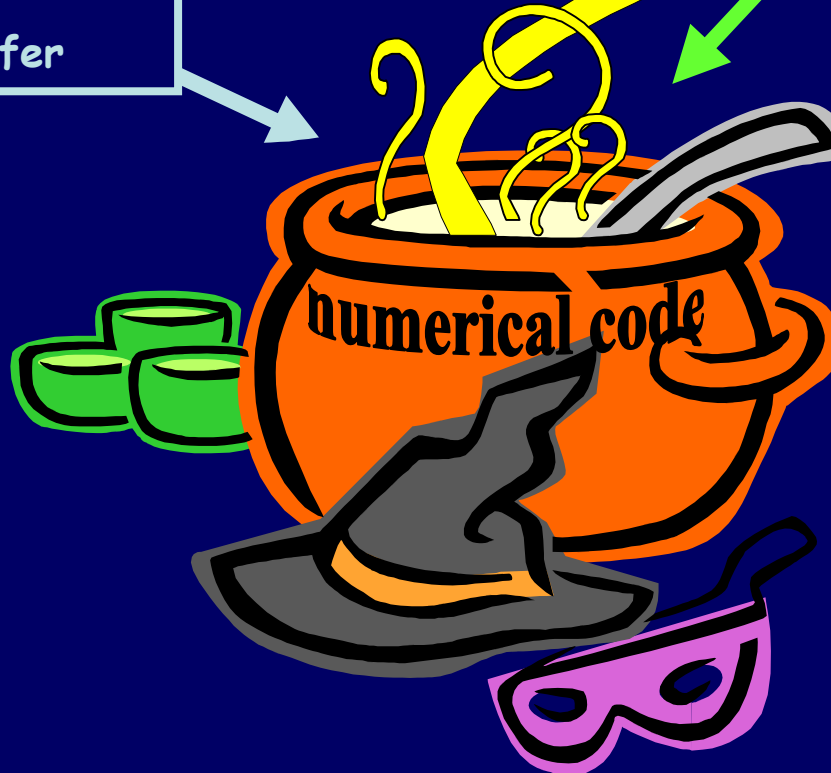
⇒ more detailed calculations are needed

one-dimensional photoionization codes

- elements
- geometry
- reasonable hypotheses
- physical processes
- atomic data
- radiation transfer

physical
conditions
= $f(r)$

- input parameters:
- ionizing radiation spectrum
 - gas density
 - chemical abundance



- most abundant elements ($Z=1 \rightarrow 30$)

- elements
- geometry
- reasonable hypotheses
- physical processes
- radiation transfer
- atomic data

- spherical symmetry
- plane-parallel geometry

- ionization equilibrium $dn(X^m)/dt = 0$
- thermal balance $G = L$

- all processes known to be important in a low-density gas

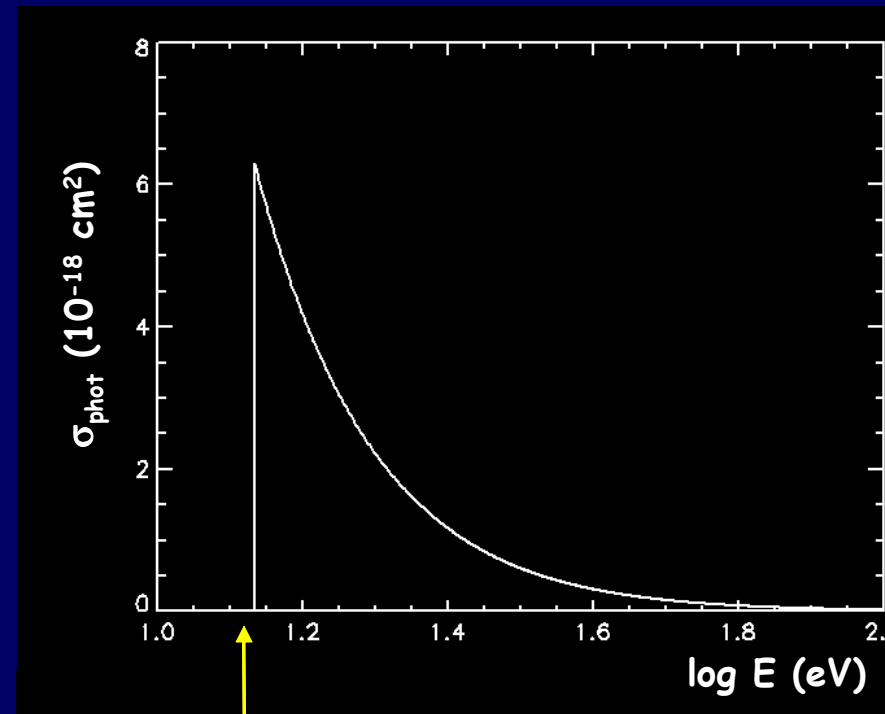
⇒ main assumptions of one-dimensional photoionization codes:

ionization equilibrium:

$$\frac{dn(X^m)}{dt} = 0$$

ionization rate = recombination rate

- photoionization (+ Auger effect)
(primary and diffuse radiation)



13.6 eV

⇒ main assumptions of one-dimensional photoionization codes:

ionization equilibrium:

$$\frac{d n(X^m)}{d t} = 0$$

ionization rate = recombination rate

- photoionization (+ Auger effect)
(primary and diffuse radiation) $X^i + h\nu \rightarrow X^{i+1} + e^-; X^i + h\nu \rightarrow X^{i+n} + ne^-$
- collisional ionization $X^i + e^- \rightarrow X^{i+1} + e^- + e^-$
- charge transfer reactions (ion.) $X^i + Y_j \rightarrow X^{i+1} + Y_{j-1}$
- dielectronic recombination $X^i + e^- \rightarrow X^{i*} \rightarrow X^{i-1} + h\nu_1 + h\nu_2 + \dots$
- radiative recombination $X^i + e^- \rightarrow X^{i-1} + h\nu_1 + h\nu_2 + \dots$
- charge transfer reactions (rec.) $X^i + Y_j \rightarrow X^{i-1} + Y_{j+1}$

(3-body recombination; ionization by cosmic rays; ...)

⇒ main assumptions of one-dimensional photoionization codes:

thermal equilibrium:

energy gains = energy losses

energy input:

- photoionization



cooling:

- recombination (+ cascade)



- thermal bremsstrahlung



- coll. exc. + rad. deexcitation



- two-photon decay



equations for the thermal balance
f [n(Xⁱ), T_e, n_e; ionizing spectrum]

equations for the ionization balance
f' [n(Xⁱ), T_e, n_e; ionizing spectrum]

**iterative
calculation**

```
graph LR; A["equations for the thermal balance  
f [n(Xi), Te, ne; ionizing spectrum]"] --- B["equations for the ionization balance  
f' [n(Xi), Te, ne; ionizing spectrum]"]; B --- C["iterative  
calculation"]
```

- most abundant elements ($Z=1 \rightarrow 30$)

- elements
- geometry
- reasonable hypotheses
- physical processes
- radiation transfer
- atomic data

- spherical symmetric
- plane-parallel geometry

- ionization equilibrium $dn(X^m)/dt = 0$
- thermal balance $G = L$

- all processes known to be important

- updated data

- on-the-spot approximation
- outward-only approximation
- Monte Carlo techniques
- escape probability
- ...

- OB stars: individual, group, starburst = $f(\text{age})$
- PN central stars
- AGN sources (power-law spectrum, starburst, both, ...)
- cosmic background radiation (Σ AGNs, ...)
- ...

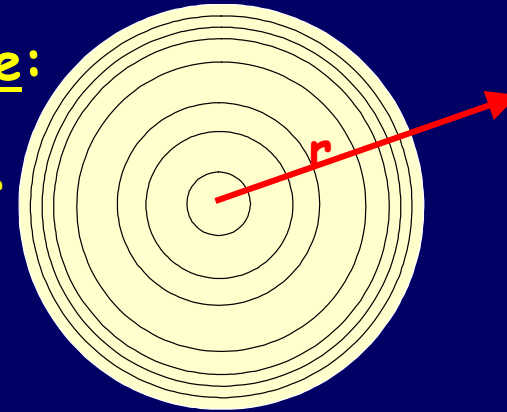
input parameters:

- ionizing radiation spectrum
- gas density
- chemical abundance

from line fitting or
typical values

Example of the output of a photoionization code:

- spherical symmetry
- source at the center



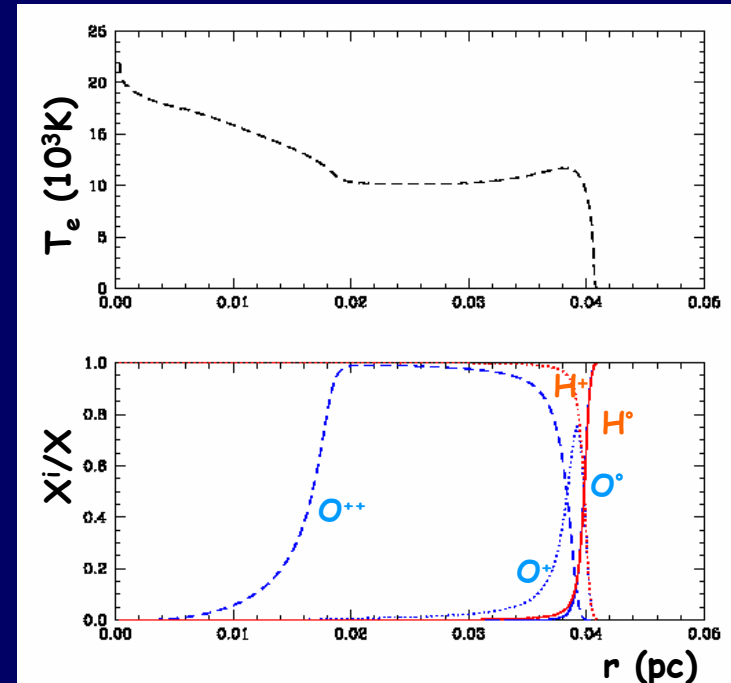
physical conditions = $f(r)$

$$n(X^i, r), n_e(r), T_e(r)$$

$$\varepsilon(\lambda, r)$$

$$L_\lambda = \oint \varepsilon(\lambda, r) dV$$

$$I_\lambda / I_{H\beta}$$



$$\begin{aligned} T_* &= 100\text{kK}; L_* = 3000 L_\odot \\ n_H &= 10^4 \text{ cm}^{-3} \\ Z &= Z_{\text{typical}} \end{aligned}$$

⇒ photoionization models can be used to:

a) statistical studies

to derive the general properties of a given class of emission-line objects

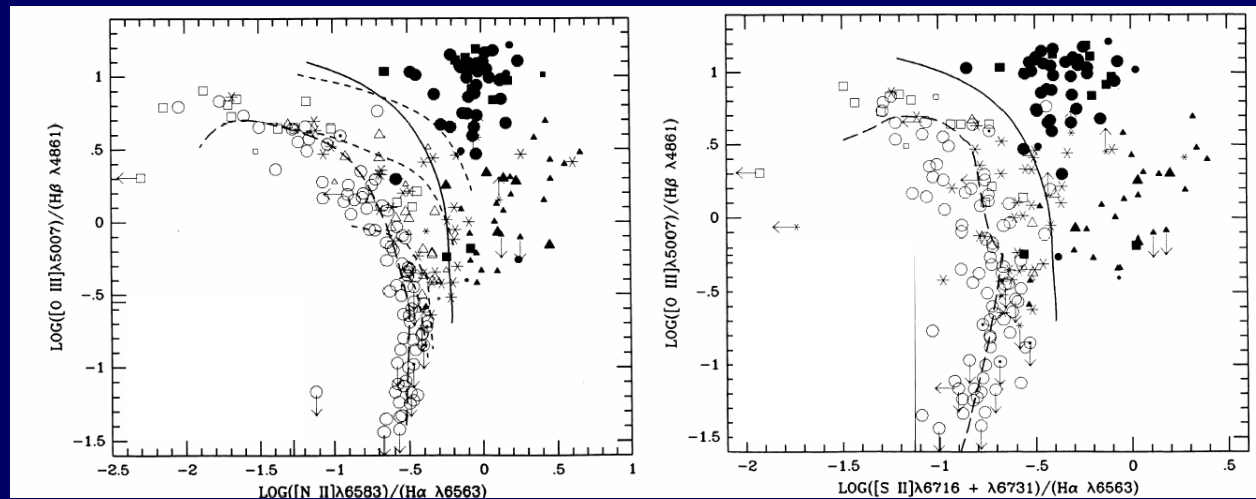
grid of models

+

observations

diagnostic diagrams

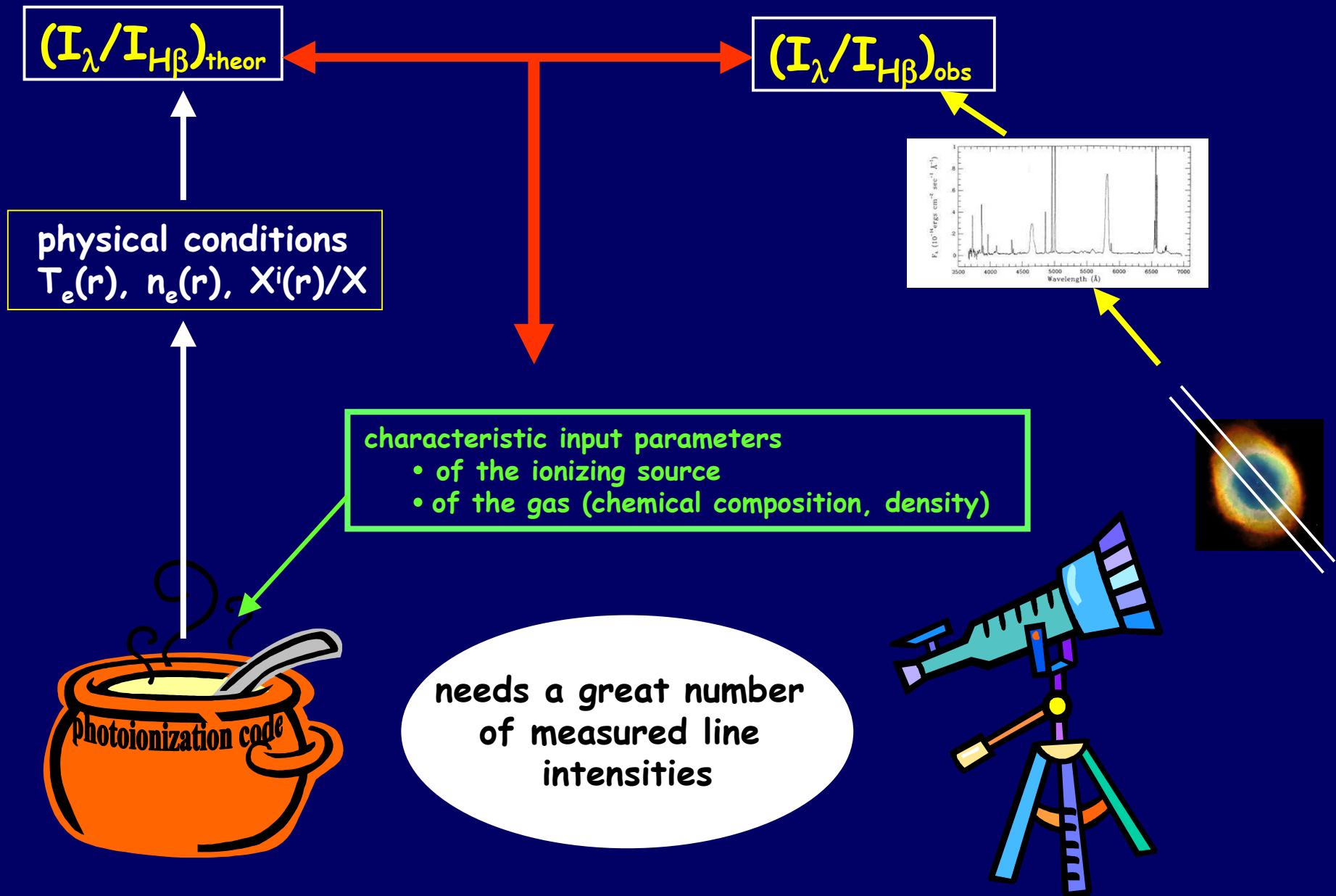
ex.:



- | | | | |
|---|------------------------|---|-----------|
| ○ | H II region | ● | Seyfert 2 |
| ⊙ | H II region in nucleus | ■ | NLRG |
| △ | Starburst galaxy | ▲ | Liner |
| □ | H II galaxy | * | NELG |

Baldwin et al. (81)

b) detailed analysis of individual objects



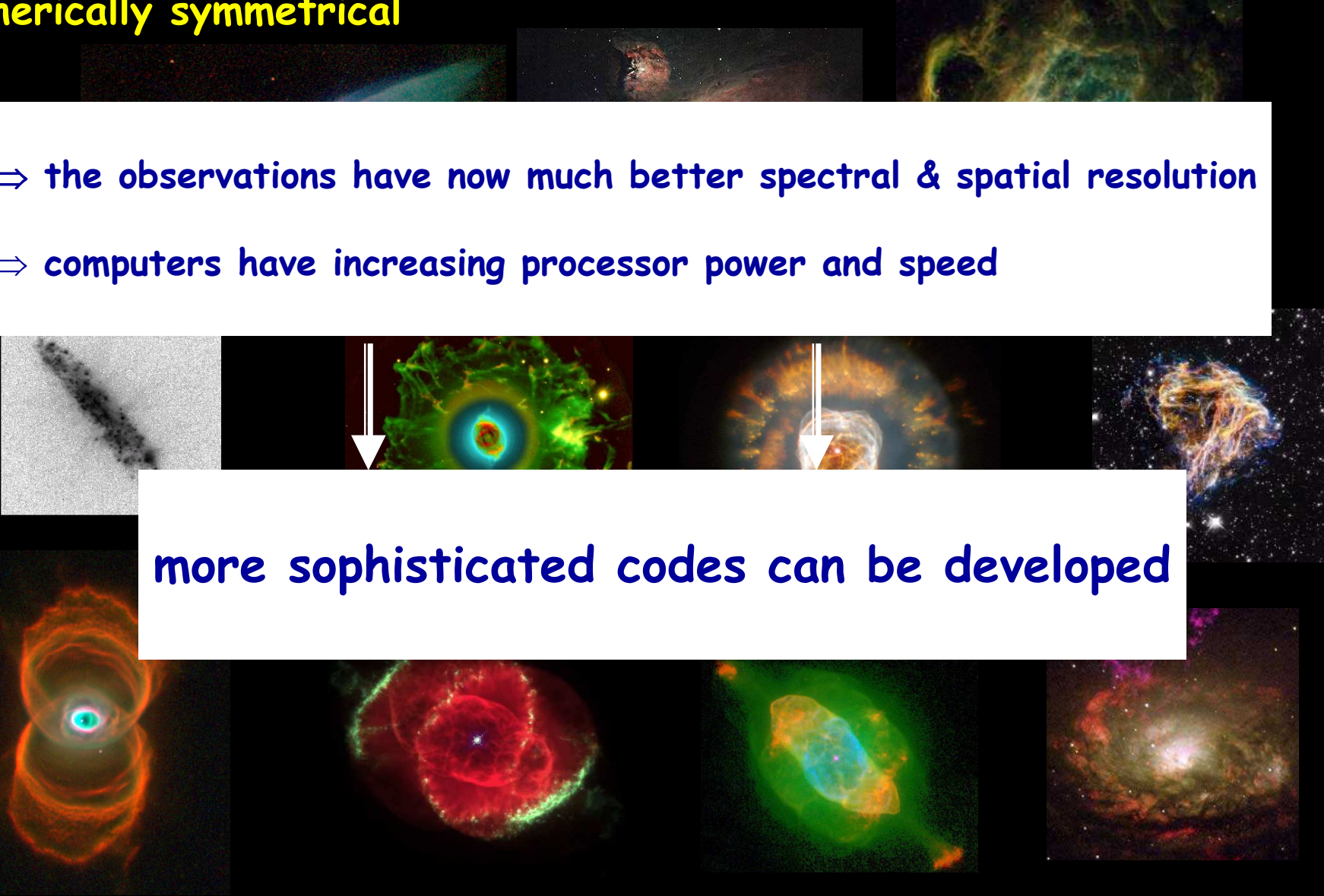
some more ingredients ...

- dust
 - most codes
- shocks
 - Viegas & Contini, Dopita (M.Elitzur lecture)
- line radiative transfer
 - Elitzur & Ferland, Netzer, Collin-Souffrin, ..
- relativistic electrons
 - Viegas & Gruenwald
- molecules in the ionized region
 - Aleman & Gruenwald
- hydrodynamical equations
 - + microphysical processes (simplified)
 - Frank, Mellena, ...
- time dependent models
 - Tylenda, Szczerba
- chemodynamical models
 - Friaça & Terlevich

- the photoionization codes are, in general, one-dimensional
- but most photoionized regions are neither **homogeneous nor spherically symmetrical**

- ⇒ the observations have now much better spectral & spatial resolution
- ⇒ computers have increasing processor power and speed

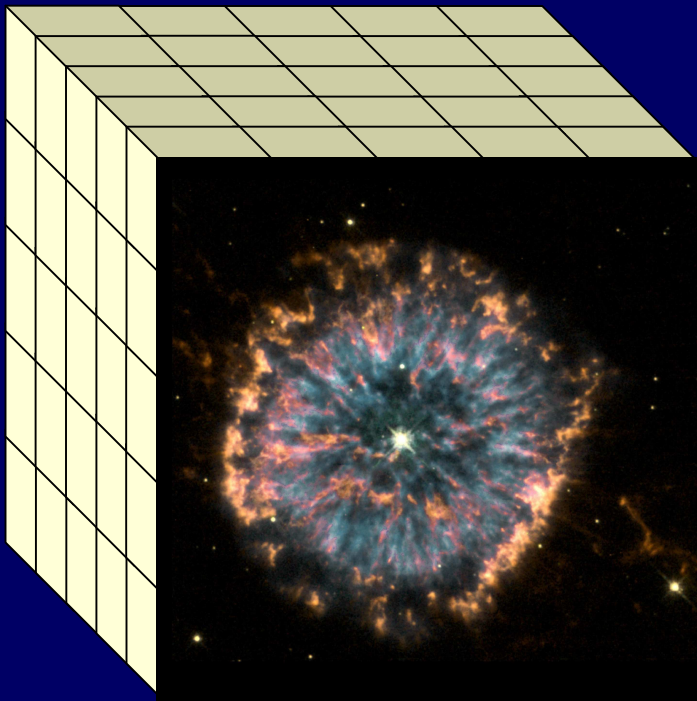
more sophisticated codes can be developed



3D photoionization code



Gruenwald et al. (1997)

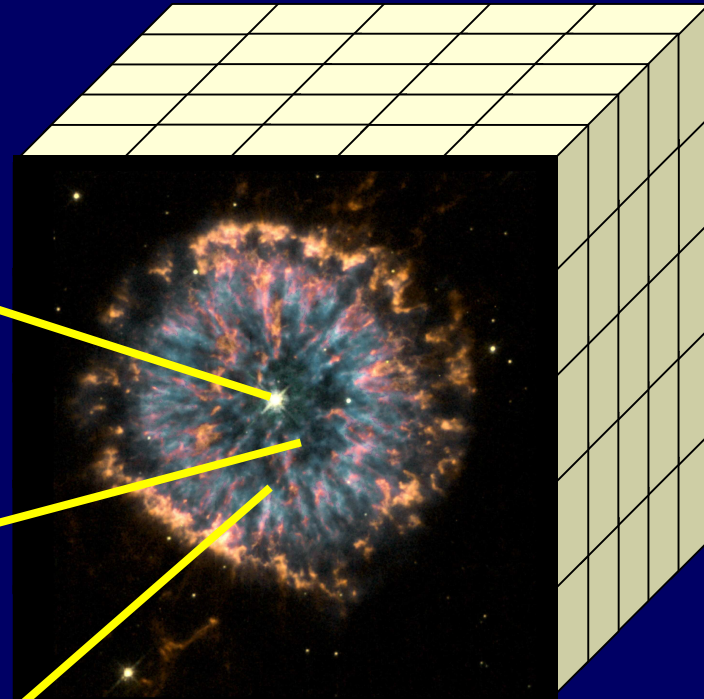


input parameters:

spectral distribution of
the ionizing radiation

chemical composition (X/H)

density distribution
 $n_H(x, y, z)$



for each cell:

hypotheses

homogeneous
physical conditions

ionization equilibrium
 $d(X^i/X)/dt = 0$

thermal equilibrium
 $G = L$

input parameters

- chemical composition
- density (x, y, z)

calculation

incident radiation (diluted)
• from the primary source
• from other cells

iterative calculation

$X^i/X(x, y, z)$
 $n_e(x, y, z)$
 $T_e(x, y, z)$

ex.: results for a nebula with a dense waist of matter:

$$\begin{matrix} X^i/X(x,y,z) \\ T_e(x,y,z) \\ n_e(x,y,x) \end{matrix}$$

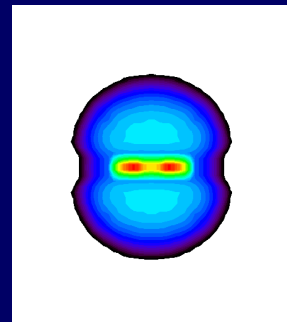
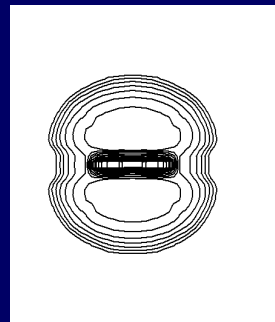
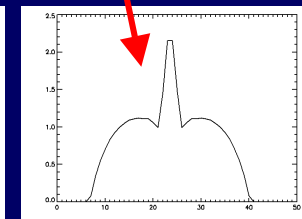
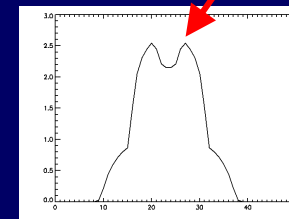
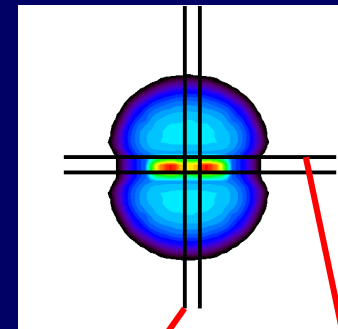
$$\epsilon_\lambda(x,y,z)$$

$$L_\lambda = \int \epsilon_\lambda(x,y,z) dV$$

$$I_\lambda / I_{H\beta}$$

monochromatic images

F_λ distribution in a given aperture



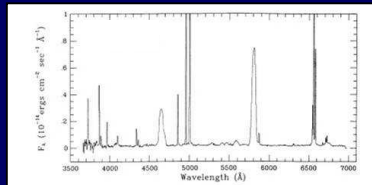


typical input parameters:

- for the ionizing radiation (T_* , L_*)
- for the gas (chemical composition, spatial structure)

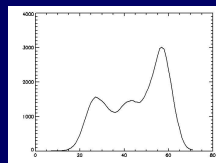
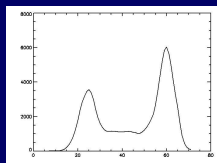
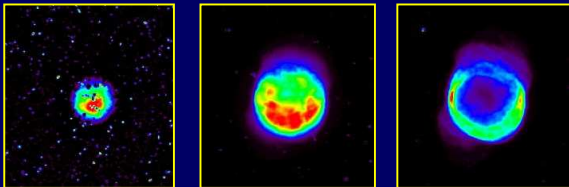
observations

numerical code



$I_\lambda / I_{H\beta}$

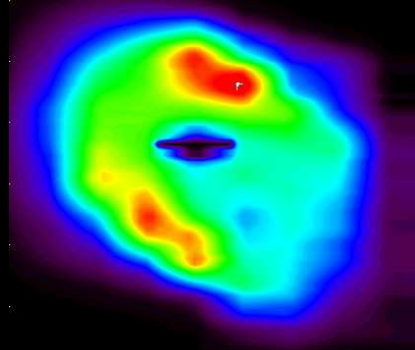
monochromatic images



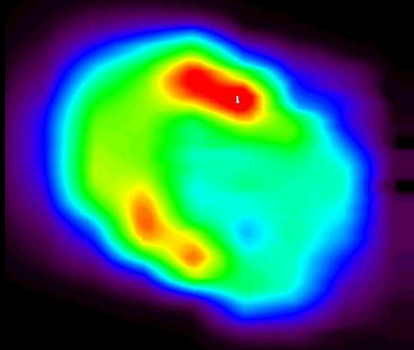
F_λ distribution in a given aperture

NGC 3132

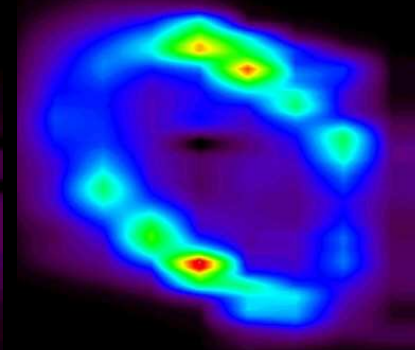
LNA-1.6m



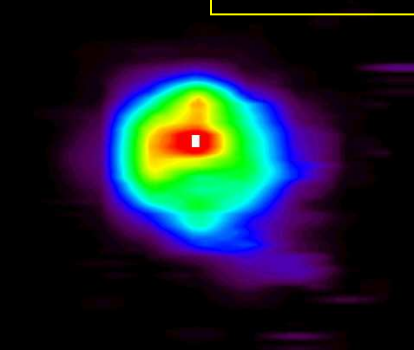
H β



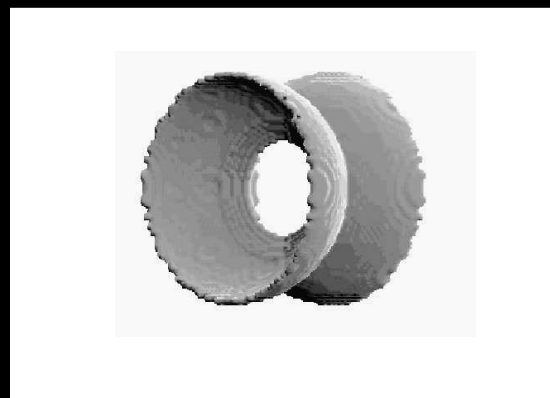
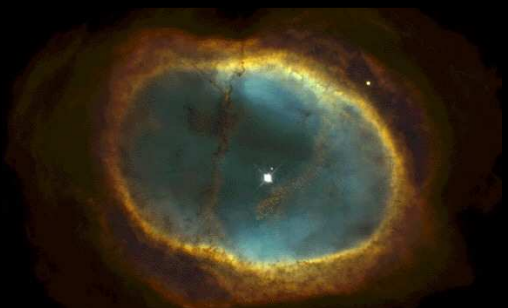
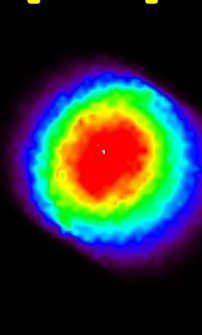
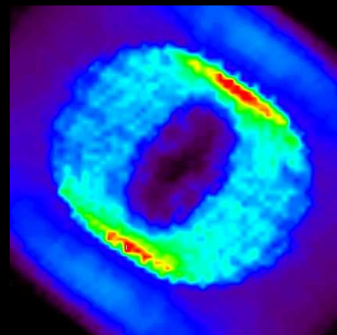
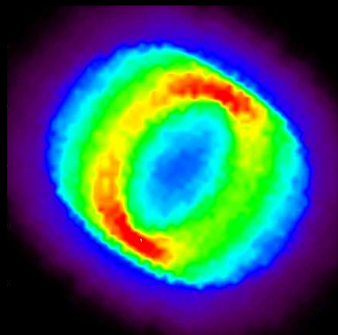
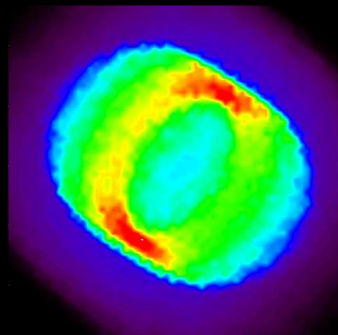
[OIII]



[NII]



[HeII]



conclusion: the geometry is not ellipsoidal, as suggested by the observed images

what can be improved?

- more accurate atomic data → Opacity Project
Iron Project
- better resolution of the observations (spectral and spatial)
- better resolution of the models (higher processor power and speed)
- information about grain composition and structure
- more reliable model atmosphere (enormous progress in last years)
- increasing sensibility of the observations
- coupling of photoionization codes with hydrodynamical calculations
- detailed analysis of additional heating - winds, shocks, ...
- for giant H II regions - correction of stellar absorption
- ...



Some references for the study of (photo-)ionized regions:

- basic physics of photoionized nebulae
 - Osterbrock (1989) - "Astrophysics of Gaseous nebulae and AGN"
 - Aller (1984) - "Physics of Thermal Gaseous Nebulae"
- empirical methods for abundance determination; icfs
 - many authors, mainly Peimbert & collaborators
- models for very dense photoionized regions, including with line transfer
 - Elitzur & Ferland, Collin-Souffrin, Netzer, Kwan & Krolik
- models including shocks (SNR, AGN)
 - Viegas & Contini, Dopita
- DIM, WIM, DIG, ...
 - Dettmar, Reynolds, Rand, Walterbos, Fergunson, ...
- dust in ionized regions
 - Stasinska & Szczerba, van Hoof, ...
- free photoionization code "Cloudy" - Ferland & co-workers
<http://www.nublado.org>