Astronomy & Asrophysics



Observing Accretion Disks I:

Basics of Observing Disks and Continuum Emission

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2 **Topics I would like to cover**

How we observe disks

- Spectral energy distributions → thermodynamics
- Emission lines → velocity field and structure
- Variability → time scales and instabilities
- Indirect imaging → spatially-resolved structure

Things we learn from observations

- What structures do we see in disks?
- How well does the theory work?
- How much faith should we have in the results?

3 General Philosophy



Intro and Continuum Emission

- The basics of observing disks
- Spectral energy distributions
- **Emission Lines and Indirect Imaging**
 - Eclipse mapping
 - Emission lines from disks and diagnostic value
 - Doppler tomography

AGN Disks and Their Emission Lines

- The disk as the source of the emission lines
- Outflows as seen through absorption lines

Organization of Lecture I

Basics First

6

- Sizes, masses, accretion rates, temperatures
- Distinction between CVs, XRBs, AGNs, YSOs
- Direct Imaging of Disks
 - Rare but extremely valuable
- Spectral Energy Distributions
 - Assumptions and basic calculations
 - Applications to specific types of systems
 - Lessons from comparison of models with data

Where do we find accretion disks?

- Interacting binary stars
 - CVs, LMXBs, (short period) Algol binaries
 - perhaps some HMXBs and some symbiotic stars
- Active galactic nuclei (AGNs)
- Protostars (YSOs)

Other types of disks

- Debris disks in main seq. stars (e.g., β Pictoris)
- Outflow disks (Be and other hot stars)

8 **Conventions**

Measure lengths in units of the gravitational radius

$$r_{\rm g} \equiv \frac{GM_{\bullet}}{c^2} = M_{\bullet} \qquad (G = c = 1)$$

For the Sun (1 M_{\odot}) $r_g = 1.5 \text{ km}$ For a 10⁸ M_{\odot} black hole $r_g = 1 \text{ AU}$

For compact notation, adopt a dimensionless radial coordinate

Kepler's laws take a very simple form (in Greek), for example

$$\xi \equiv \frac{r}{r_{\rm g}} = \frac{r}{M_{\bullet}}$$
$$\beta = \frac{1}{\sqrt{\xi}}$$
(where $\beta = v/c$

Geometry and Energetics

Inner disk radius:

9

radius of star or close to event horizon

• Outer disk radius:

set by tidal truncation in binaries (or gravitational instability in case of AGNs?)

• Mass of central object:

can be very hard to measure, especially in AGNs

• Accretion rate:

$$L_{\rm bol} = \eta_{\rm rad} \frac{GM\dot{m}}{r_{\rm in}} = 4\pi d^2 \left(\frac{F_{\rm obs}}{\zeta_{\rm bol}}\right)$$

10 Inner Disk Radius

- radius of a WD ~ 10^8 - 10^9 cm (M < 1.4 M_{\odot}) $R_{\rm WD} \approx 5 \times 10^8 \left(M/{\rm M}_{\odot} \right)^{-0.8}$
- radius of a NS ~ 10-15 km (M=1.4-2.7 M_{\odot})

M/R depends on NS equation of state

radius of event horizon for non-rotating BHs

 $R_{\rm S}=2GM/c^2$

- ~ 6 km for Stellar BH (M=4 M $_{\odot}$)
- ~ 2 AU for Massive BH ($M = 10^8 M_{\odot}$)

11 Outer Disk Radius

- CVs and LMXBs
 - $R_D \sim 0.6 R_{L1} \sim 0.3 a$





• AGNs: ???

radius of marginal self gravity

e.g., Goodman & Tan (2004, ApJ, 608, 108)

$$R_{\rm sg} \approx 3100 \left(\frac{\alpha_{0.3}}{\mu}\right)^{1/3} \left(\frac{\ell_E}{\epsilon_{0.1}}\right)^{1/6} \left(\frac{\kappa}{\kappa_{\rm es}}\right)^{1/2} M_8^{1/2} \,\mathrm{AU}$$

 $R > R_{sg}$: terra incognita

	YSO	CV	NS XRB	10 ⁸ M⊚ AGN
$\boldsymbol{\xi}$ in	~I0 ⁷	> 1800	4–7	2
ξout	~10 ¹⁰	10 ⁵ —10 ⁶	10 ⁵ —10 ⁶	>1500

 $\xi = r/r_g = rc^2/GM$

12

13 Temperature and Disk Mass

• Temperature, non-rotating black hole disks,

$$T_{\rm in} \propto M^{-1/4} \left(\frac{L}{L_{\rm Edd}}\right)^{1/4}$$

 $^{\rm L} \approx 20 \left(\frac{M}{10^8 M_{\odot}}\right)^{-1/4} \left(\frac{L}{L_{\rm Edd}}\right)^{1/4} \, {\rm eV}$

- Disk Mass:
 - Estimate: mass should be accreted in a viscous time $M_{\rm D} \sim \dot{m} \ \tau_{\rm visc}(\xi_{\rm out})$
 - Mass also derived from fits to the SED

	YSO	CV	NS XRB	10 ⁸ M _☉ AGN
M _D (M _☉)	~ 0.1–1	~ 10-11	~ 10-11	~ 10 ⁵
T _{max}	~ 10 ² K	~10 ⁴ K	~10 ⁷ K	~10 ⁵ K
		few eV	I–2 keV	~20 eV
Ideal band	FIR	UV	X-ray	EUV

15 Can We Image Disks Directly?

Nearest CVs

• 1 R_{\odot} @ 50 pc \rightarrow 0.6 milli-arcsec

Nearest AGNs

- 1000 AU @ 1 Mpc \rightarrow 1 milli-arcsec
- but a few large disks imaged with HST in optical

Nearest YSOs

- 100 AU @ 150 pc → 0.6 arcsec
- imaged with the HST and with ground-based interferometers (near-IR and mm-wave bands)

16 Gaseous Disk in M87



HST/WFPC2 Hα + [N II] image Ford et al. 1994, ApJ, 435, L27

17 Gaseous Disk in M81





HST/WFPC2 images: 57.5 × 40.3 arcsec Devereux et al. 1997, ApJ, 481, L71

18 Direct Images of Protostellar Disks



Gomez's Hamburger IRAS 18059-3211: disk in silhouette (Hubble Heritage)

19 Direct Images of Protostellar Disks



HST/NICMOS images of YSOs in Taurus

20 The luxury of imaging...

- Combine morphology and SED constraints to infer detailed sructure.
 - e.g., dust properties from synthetic scattered light images and model S.E.D.s



21 Sizes of Protostellar Disks from Interferometry

Size-luminosity relation in Herbig Ae/Be stars



Monier et al. 2005, ApJ, 624, 832

22 Disk structure from size measurements

"Optically-thin Cavity" Disk Model



2 **M** 80 Monier et al. 2005, ApJ, 624,

23 Can we isolate the disk without imaging it?

- Observed light includes:
 - Boundary Layer and central star
 - Jet (beamed)
 - Companion star
 - Host galaxy of AGN



- Disk light passes through
 - Atmosphere/wind
 - Curved spacetime



24 The Disk-Star Boundary Layer



Half of the accretion luminosity released in boundary layer

In CVs: $T \sim 10^8$ K if optically thin (e.g., Tylenda 1979) $T \sim 10^5$ K if otically thick (e.g., Pringle 1977)

25 Spectral Energy Distributions (S.E.D.s)

• Start with temperaure distribution of disk

$$T(r) = T_* \left(\frac{r}{R_{\rm in}}\right)^{-3/4} \left[1 - \left(\frac{r}{R_{\rm in}}\right)^{-1/2}\right]^{1/4}$$

where $T_* = \left(\frac{3 G M \dot{m}}{8\pi \sigma_{\rm SB} R_{\rm in}^3}\right)^{1/4}$

Assume local black body emission and integrate over the entire disk

$$f_{\nu} = \frac{\cos i}{d^2} \int_{R_{\star}}^{r_{\text{out}}} B_{\nu}[T(r)] \ 2\pi r dr$$

26 Resulting shape

600 Å 200 eV For XRBs 2 keV 5000 Å 20 eV For AGNs **6** μm 1 Ry For CVs 3 eV $1 \mu m$ 2 $\propto v^{1/3}$ В log $F_{
u}$ (arbitrary units) Α $\propto v^2$ $\propto e^{-hv/hT}$ 0 2 0 $\log (h\nu/kT_{out})$

27 Wavelength dependent disk size

 $f_{\nu}(r) \propto 2\pi r \ B_{\nu}[T(r)]$

Disk should appear smaller at shorter wavelengths



In principle, we can learn a lot about the disks if we can measure their SEDs and fit them with models.

test disk models

infer accretion rate

turn them into tools

Doing this in practice is not easy.



need distances and masses

need to observe in "difficult" bands

need to do very messy physics, with many uncertainties...

29 How messy and how uncertain...

The use of black bodies is not adequate...

- Disks have atmospheres with vertical structure
- Scattering is important in very hot disks
- Dust is present in very cold disks
- Energy generation need not be concentrated in the mid-plane of the disk

Need sophisticated atmosphere models

- with proper radiative transfer
- and relativity (special and general)



SED models: Radiative transfer in dusty disk, illuminated by central star

Monier et al. 2005, ApJ, 624, 832

31 Accretion rates from SED fitting

Evolution of protostellar disks in the Orion OB1 association.

Model fits to SEDs yield the accretion rate (among other parameters).

Accretion rate appears to drop as the stars age.

Disks are consumed ("viscous evolution").



Calvet et al., 2005, AJ, 129, 935

32 Disk size & structure from SED fitting

Fits of SED models indicate a gap in the inner disk. Gap manifests itself as a dip in the SED at 1 micron. Spitzer IRS data constrain the SED at the crucial wavelength.



figures from Calvet et al. 2005, 630, L185

33 Imaging combined with SED fitting

- Measurements of disk inner radii with Keck interferometer.
- Fix values in SED models.
- Infer vertical strucure of inner disk.

Eisner et al. 2004, ApJ, 613, 1049

"Optically-thin Cavity" Disk Model



sketch from Monier et al. 2005, ApJ, 624, 832

34 Applications to CVs



- Energy produced in disk midplane.
- Propagates through planeparallel atmosphere.
- Limb darkening,
 foreshortening, rotational
 broadening of lines.
- Ideal range to observe: near-UV, 912–2000 Å
- Relativity not needed

35 Examples of resulting spectra

- Sequence of increasing dM/dt (bottom to top)
- Continua get bluer with dM/dt
- Absorption lines get shallower

from Wade & Hubeny 1998, ApJ, 509, 350



36 Effect of rotation and limb darkening



Weak lines blend together into a pseudo-continuum NOT A BLACKBODY! from Wade & Hubeny 1998, ApJ, 509, 350

38 Comparison with observations

 Cannot match flux and slope of UV spectra of CVs at the same time

Wade 1988, ApJ, 335, 394 Puebla et al. 2008, AJ, 134, 1923

- Possible avenues to a solution:
 - Temperature profile of inner disk.
 - Optically thin gas in disk

Wind



39 Applications to Luminous AGNs

Average S.E.D.s of quasars



40 Emission models for thin disks in AGNs

- Historically, stellar atmospheres used to approximate disk atmosphere (matched in T and g). e.g., Sun & Malkan 1989, ApJ, 346, 68
- Shortcomings:
 - Disks have a different density than stars for the same combination of T and g
 different opacity
 - Temperature gradient of disk atmospheres different from stellar atmospheres
 different absorption lines

41 Models for AGN accretion disk SEDs

by Laor & Netzer 1989, MNRAS, 238, 897

- Calculation of H+He disk atmospheres based on vertical structure from α-disk models
 - free-free, bound-free opacity (latter is important)
 - comptonizaton, when necessary
 - various prescriptions for vertical temperature gradient
 - special and general relativity

42 Model examples



from Laor & Netzer 1989, MNRAS, 238, 897

43 More examples: effect of inclination



from Laor & Netzer 1989, MNRAS, 238, 897

44 **Comparison with data: mixed success**



• Fits not always good.

Model parameters not well constrained

Laor 1990, MNRAS, 246, 369

45 **Comparison with data: issues**



- FIR power-law needs to be invoked ad hoc
- No smooth connection to X-ray data

46 **Comparison with data: more issues**

- Ly edge: predicted but not observed. Breaks observed instead.
 - Need even better atmosphere models.
- Edges can be smoothed by scattering, but should lead to polarization.
 - Early measurements of polarization did not agree with predictions (too high or too low).
 - Latest polarization observations (Kishimoto et al. 2003, 2004, 2005) may have uncovered Balmer edge and disk thermal emisson in polarized light.

47 High-Resolution UV Spectra of BBB



Shang et al. 2005, ApJ, 619, 41

- Find breaks in UV spectra of quasars and determine slope on either side
- Compare break wavelengths and slopes with hot atmosphere models (from Hubeny et al. 2001)
- Reasonable slopes but none of the expected correlations between model parameters found.

48 Application to Low-Luminosity AGNs

Collection of LL-AGN SEDs from Ho 1999, ApJ, 516, 672



49 Main radiation processes involved

- From hot torus:
 - Synchrotron followed by inverse compton
 - Bremstrahlung
- From the thin disk
 - thermal emission
 (could be irradiated)
- From the jet
 - Synchrotron followed inverse compton



50 Fit to the SED of NGC 3998

from Ptak et al. 2004, ApJ, 606, 173



51 What we know about LL-AGN SEDs

- VLA/VLBA observations
 - compact, high-T_B radio cores and jets
 - jets required for radio portion of SED.
- Chandra X-ray observations
 - unresolved X-ray sources with power-law spectra
 - may be emission from ADAF/RIAF
- HST UV Observations
 - variable, unresolved sources at 2500 Å
 - does this agree with ADAF models?

52 New Average S.E.D. of LL-AGNs



53 Fits to individual objects: M81

Jet explains radio and X-rays Thin disk explains UV+optical small contribution from ADAF



54 Fits to individual objects: M84

Jet explains radio and X-rays sma Thin disk explains UV+optical

small contribution from ADAF



55 Fits to individual objects: NGC 4274



- Ambiguous case: more than one physically plausible solution possible
- Value of δ controls contribution of ADAF to the X-ray band and leads to different solutions

56 Applications to XRBs

- New generation of models based on latest theoretical developments.
 - Most sophisticated atmosphere models yet
 - Electron scattering and comptonization
 - Fully relativistic ray propagation
 - Boundary conditions at inner disk edge
 - Vertical energy dissipation profile
- Compare to BH-XRBs with known masses, distances, and inclinations

57 Application to XRBs and the B. H. spin

• Multi-color disk with modified blackbodies



- Adjust spin parameter to fit X-ray spectrum.
- Verify by checking if $L(T) \propto T^4$

as the system varies



figure from Laor & Netzer 1989, MNRAS, 238, 897

58 Applications to XRBs: Recap

- New generation of models based on latest theoretical developments.
 Davies et al. 2006, ApJ 621, 372
 - Most sophisticated atmosphere models yet
 - Electron scattering and comptonization
 - Fully relativistic ray propagation
 - Boundary conditions at inner disk edge
 - Vertical energy dissipation profile
- Compare to BH-XRBs with known masses, distances, and inclinations

59 **Results**

from Davies et al. 2006, ApJ 621, 372



60 Can this method be applied to AGNs?

- Easier to measure distances to AGNs
- But BH masses and disk inclinations are very difficult to get
- AGN disks have kT ~ 20 eV; anathema!
 - radiate in FUV; thermal spectrum unobservable
 - atmosphere models more difficult (cf, CVs)
- If X-ray emission models are right (Lecture 3), disks must be intensely radiated.

61 Where do we stand?

- SED models have only partial success.
 - Qualititative features of the data are reproduced.
 - Most successful application is to hot XRB disks.
 - Spectacular failure in AGN SEDs.
- Where could the problem be?
 - + Fundamental theory has open questions.
 - Atmosphere models are also incomplete (we know there are winds).
 - Irradiated atmospheres (e.g., current big picture in AGNs not self-consistent) !

End of Lecture I