

Density dependent magnetic field and the equation of state of baryonic matter.

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Outlook

Density
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- Used model.
- Equations.
- Chosen parameterization.

Lagrangian:

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$$\mathcal{L}_b = \sum_{b=n,p,hyp} \bar{\psi}_b [i\gamma_\mu \partial^\mu + q_b \gamma_\mu A^\mu - g_{\rho b} \tau_{3b} \gamma_\mu \vec{\rho}^\mu - m_b^* - g_{\omega b} \gamma_\mu \omega^\mu - \frac{1}{2} \mu_N k_b \sigma_{\mu\nu} F^{\mu\nu}] \psi_b \quad (1)$$

$$\begin{aligned} \mathcal{L}_m = & \frac{1}{2} (\partial_\mu \sigma \partial^\mu \sigma - m_\sigma^2 \sigma^2) - \frac{1}{3} b m_n (g_{\sigma N} \sigma)^3 - \frac{1}{4} c (g_{\sigma N} \sigma)^4 \\ & + \frac{1}{2} m_\omega^2 \omega_\mu \omega^\mu - \frac{1}{4} \Omega^{\mu\nu} \Omega_{\mu\nu} + \frac{1}{2} m_\rho^2 \vec{\rho}_\mu \cdot \vec{\rho}^\mu - \frac{1}{4} F^{\mu\nu} F_{\mu\nu} \\ & - \frac{1}{4} P^{\mu\nu} P_{\mu\nu} + \frac{1}{4!} \xi g_\omega^4 (\omega_\mu \omega^\mu)^2 + \Lambda_\omega (g_\rho^2 \vec{\rho}_\mu \cdot \vec{\rho}^\mu) (g_\omega^2 \omega_\mu \omega^\mu). \end{aligned} \quad (2)$$

$$\mathcal{L}_l = \sum_{l=e,\mu} \bar{\psi}_l [i\gamma_\mu \partial^\mu + q_l \gamma_\mu A^\mu - m_l] \psi_l \quad (3)$$

$$\mathcal{L} = \mathcal{L}_b + \mathcal{L}_m + \mathcal{L}_l \quad (4)$$

Lagrangian.

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- (ψ_b) - Baryon field
- (ψ_l) - Lepton field
- (σ) - Scalar meson
- (ω^μ) - Vectorial meson
- (ρ^μ) - Isovectorial meson
- $g_\sigma, g_\omega, g_\rho$ - Meson-baryon coupling constants
- γ_μ - Dirac Matrices
- $A^\mu = (V, \vec{A})$ - Lorentz tensor
- $k_p = 1.79285,$
 $k_n = -1.91315$ - Anomalous magnetic moment (AMM)
- τ_3 - Isospin projection
- q_b - Electrical charge of the particle b
- μ_N - Nuclear magneton.

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- $m_b^* = (m_b - g_{\sigma b}\sigma)$ -
Baryon effective mass
- $\Omega_{\mu\nu} = \partial_\mu\omega_\nu - \partial_\nu\omega_\mu$
- $P_{\mu\nu} = \partial_\mu\rho_\nu - \partial_\nu\rho_\mu$
- $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$
- $\sigma_{\mu\nu} = \frac{i}{2}[\gamma_\mu, \gamma_\nu]$ -
Electromagnetic tensor
- ξ - Omega meson
self-interaction parameter
- Λ_ω - Parameter that
changes the density
dependence of the
symmetry energy
- b and c - Parameters
related with the weights of
the non-linear scalar terms

Magnetic field.

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$$B(n) = B_{surf} + B_0 \left[1 - \exp \left(-\beta \left(\frac{n}{n_0} \right) \right)^\gamma \right] \quad (5)$$

- $B_{surf} = 10^{15} \text{ G}$, $n_0 = 0.153 \text{ fm}^{-3}$
- Fast: $\gamma = 4$, $\beta = 1.83$ Slow: $\gamma = 1$, $\beta = 1.92$
- The unit of the magnetic field is the critical field $B_e^c = 4.414 \times 10^{13} \text{ G}$, so that $B = B_0 / B_e^c$.

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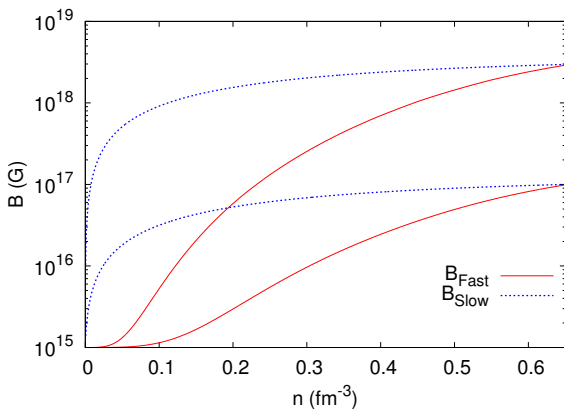


Figura: Variable density dependent magnetic field.

Magnetic field.

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Dividing the work in two types:

- Fixed magnetic field (FIX): The magnetic field is considerate constant on the equations of state.
- Variable density dependent magnetic field (VAR): The magnetic field varies on the equations of state, see eq. (??).
- Both cases consider a density dependent variable magnetic field B term on the total energy density and pressure, eq. (??).

Mesonic field equations:

Utilizing the mean field approximation and the Euler-Lagrange equations.:

$$\begin{aligned}m_{\sigma}^{*2}\sigma_0 &= -cg_{\sigma}^4\sigma_0^3 - bm_n g_{\sigma}^3\sigma_0^2 + g_{\sigma}n^s \\m_{\omega}^{*2}\omega_0 &= -\frac{\xi}{6}g_{\omega}^4\omega_0^3 + g_{\omega}n_b \\m_{\rho}^{*2}\rho_{03} &= \tau_{3b}g_{\rho}n_3,\end{aligned}\tag{6}$$

- $m_{\rho}^{*2} = m_{\rho}^2 + 2\Lambda_{\omega}g_{\rho}^2g_{\omega}^2\omega_0^2$
- $m_{\omega}^{*2} = m_{\omega}^2 + 2\Lambda_{\omega}g_{\rho}^2g_{\omega}^2\rho_{03}^2$
- $g_{iH} = X_{iH}g_{iN}$
- $X_{\sigma H} = 0.700, X_{\omega H} = 0.738,$
 $X_{\rho H} = 0.600$
- $n^s = n_p^s + n_n^s$
- $n_b = n_p + n_n$
- $n_3 = \frac{(n_p - n_n)}{2}$

Densities:

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Scalar¹:

$$n_b^s = \frac{q_p B m_p^*}{2\pi^2} \sum_{\nu}^{\nu_{\max}} \sum_s \frac{\sqrt{m_p^{*2} + 2\nu q_p B - s\mu_N k_p B}}{\sqrt{m_p^{*2} + 2\nu q_p B}} \quad (7)$$
$$\times \ln \left| \frac{p_{F,\nu,s}^p + E_F^p}{\sqrt{m_p^{*2} + 2\nu q_p B - s\mu_N k_p B}} \right|$$
$$n_b^s = \frac{m_n^*}{4\pi^2} \sum_s \left[E_{F,s}^n p_{F,s}^n - \bar{m}_n^2 \ln \left| \frac{p_{F,s}^n + E_F^n}{\bar{m}_n} \right| \right]$$

¹J.Phys.G: Nucl. Part. Phys. **35** (2008) 125201 (17pp)

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Vectorial:

$$n_b = \frac{q_p B}{2\pi^2} \sum_{\nu}^{\nu_{max}} \sum_s p_{F,\nu,s}^p \quad (8)$$

$$n_b = \frac{1}{2\pi^2} \sum_s \left[\frac{1}{3} (p_{F,s}^n)^3 - \frac{1}{2} s \mu_N k_n B \right. \\ \left. \times \left(\bar{m}_n p_{F,s}^n + E_F^{n2} \left(\arcsin \left(\frac{\bar{m}_n}{E_F^n} \right) - \frac{\pi}{2} \right) \right) \right]$$

$$n_l = \frac{|q_l| B}{2\pi^2} \sum_{\nu}^{\nu_{max}} \sum_s p_{F,\nu,s}^l$$

$$n_B = \sum_{b=n,p,hyp} n_b + n_l \quad (9)$$

Baryons:

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b	M_b (MeV)	q_b	μ_b/μ_N	k_b
p	938.3	1	2.79	1.79
n	939.6	0	-1.91	-1.91
Λ^0	1115.7	0	-0.61	-0.61
Σ^+	1189.4	1	2.46	1.67
Σ^0	1189.4	0	-1.61	-1.61
Σ^-	1189.4	-1	-1.16	-0.38
Ξ^0	1314.9	0	-1.25	-1.25
Ξ^-	1314.9	-1	-0.65	0.06

Tabela: Masses, charges, magnetic moments e anomalous magnetic moments for baryons.

$$\text{Where } k_b = \frac{\mu_b}{\mu_N} - q_b \frac{m_p}{m_b}.$$

Parametrizations:

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Modelo	m_σ	m_ω	m_ρ	g_σ	g_ω	g_ρ
GM1	512.0	783.0	770.0	8.910	10.610	8.196

Tabela: Masses and coupling constants for mesons.

Modelo	c	b	ξ	Λ_ω
GM1	-0.001070	0.002947	0.00	0.00

Tabela: Constants.

Energy spectra:

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$$E_{\nu,s}^b = \sqrt{p_z^2 + (\sqrt{m_p^* + 2\nu q_p B} - s\mu_N k_p B)^2} + g_\omega \omega^0 + \frac{1}{2} g_\rho \rho^0 \quad (10)$$

$$E_s^b = \sqrt{p_z^2 + (\sqrt{m_n^* + p_\perp^2} - s\mu_N k_n B)^2} + g_\omega \omega^0 - \frac{1}{2} g_\rho \rho^0,$$

- $\nu = n + \frac{1}{2} - \text{sgn}(q_b) \frac{s}{2} = 0, 1, 2, \dots$ - Landau levels for fermions with electrical charge q
- s is the spin, which assumes values $+1$ for spin *up* and -1 for spin *down*
- $p_\perp^2 = p_x^2 + p_y^2$

Momenta:

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$$\begin{aligned} p_{F,\nu,s}^b &= E_F^{b2} - [\sqrt{m_b^{*2} + 2\nu q_b B} - s\mu_N k_b B]^2 & (11) \\ p_{F,s}^b &= E_F^b - \bar{m}_b^2 \\ p_{F,\nu,s}^l &= E_F^l - \bar{m}_l^2 \end{aligned}$$

- Where $p_{F,\nu,s}^b$ is the Fermi momentum for the charged baryons, $p_{F,s}^b$ for the neutron and $p_{F,\nu,s}^l$ for the leptons.
- $\bar{m}_b = m^* - s\mu_N k_b B$, $\bar{m}_l = m_l^2 - 2\nu|q_l B|$.

Landau levels:

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$$\nu_{max} = \left[\frac{(E_F^b + s\mu_N k_b B)^2 - m_b^*}{2|q_b|B} \right] \quad (12)$$
$$\nu_{max} = \left[\frac{(E_F^l)^2 - m_l^2}{2|q_l|B} \right]$$

- Landau levels are summed up to its maximum value, for which the square of the Fermi moment of the particle is still positive.

Chemical Potentials:

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$$\mu_p = E_F^p + g_\omega \omega^0 + \frac{1}{2} g_\rho \rho^0 \quad (13)$$

$$\mu_n = E_F^n + g_\omega \omega^0 - \frac{1}{2} g_\rho \rho^0$$

$$\mu_p = \mu_{\Sigma^+} = \mu_n - \mu_e$$

$$\mu_\Lambda = \mu_{\Sigma^0} = \mu_{\Xi^0} = \mu_n$$

$$\mu_{\Sigma^-} = \mu_{\Xi^-} = \mu_n + \mu_e$$

$$\mu_\mu = \mu_e$$

$$\sum_{b=n,p,hyp} q_b n_b + \sum_l q_l n_l = 0$$

Total energy density:

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$$\begin{aligned}\varepsilon_m = & \sum_{b=p,n,hyp} \varepsilon_b + \sum_{l=e,\mu} \varepsilon_l + \frac{1}{2} m_\sigma \sigma_0^2 + \frac{1}{2} m_\omega \omega_0^2 + \frac{1}{2} m_\rho \rho_0^2 \quad (14) \\ & + \frac{1}{8} \xi (g_\omega \omega_0)^4 + \frac{1}{3} b m_n (g_\sigma \sigma_0)^3 + \frac{1}{4} c (g_\sigma \sigma_0)^4 \\ & + 3 \Lambda_\omega (g_\omega \omega_0)^2 (g_\rho \rho_0)^2,\end{aligned}$$

Energy density:

$$\begin{aligned} \varepsilon_b &= \sum_{n=0}^{n_{max}} \sum_s \frac{|q_b|B}{4\pi^2} \left[p_{F,n,s}^b E_F^b + (\sqrt{m_b^2 + 2q_b B n} - s\mu_N k_b B)^2 \right. \\ &\quad \left. \times \ln \left| \frac{p_{F,n,s}^b + E_F^b}{(\sqrt{m_b^2 + 2q_b B n} - s\mu_N k_b B)} \right| \right], \\ \varepsilon_b &= \frac{1}{4\pi^2} \sum_s \frac{1}{2} E_F^{b3} p_{F,n,s}^b - \frac{2}{3} s\mu_N k_b B E_F^{b3} \left(\arcsin \frac{\bar{m}}{E_F^b} - \frac{\pi}{2} \right) \\ &\quad - \left(\frac{1}{3} s\mu_N k_b B + \frac{1}{4} \bar{m} \right) \left[\bar{m} p_{F,n,s}^b E_F^b + \bar{m}^3 \ln \left(\left| \frac{E_F^b + p_{F,n,s}^b}{\bar{m}} \right| \right) \right] \\ \varepsilon_l &= \sum_{n=0}^{n_{max}} \sum_s \frac{|q_l|B}{4\pi^2} \left[p_{F,n,s}^l E_F^l + \bar{m}_l^2 \times \ln \left| \frac{p_{F,n,s}^l + E_F^l}{\bar{m}_l} \right| \right] \end{aligned} \quad (15)$$

Pressure and total energy densities:

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$$P_m = \sum_{i=e,\mu,n,p,hyp} \mu_i n_i - \varepsilon_m \quad (16)$$

Attributing the contributions of magnetic fields, we find the total energy density and pressure:

$$\varepsilon_T = \varepsilon_m + \frac{B^2}{2} \quad P_T = P_m + \frac{B^2}{2} \quad (17)$$

Symmetry energy and slope:

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$$\begin{aligned} a_{sym} &= \frac{1}{2n_B} \frac{\partial^2 E}{\partial t^2} \\ L &= 3n_0 \left. \frac{\partial a_{sym}}{\partial n_B} \right|_{n_B \rightarrow n_0} \end{aligned} \quad (18)$$

- $n_0 = 0.153 \text{ fm}^{-3}$ - Nuclear saturation density
- $t = \frac{n_n - n_p}{n_p + n_n}$ - Asymmetry parameter.

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● THANK YOU!